

# Effects of inhomogeneities on the propagation of gravitational waves from binaries of compact objects

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*A. Ali and ASM, JCAP **01**, 054 (2017)*

*S. S. Pandey, A. Sarkar, A. Ali and ASM, JCAP **06**, 021 (2022)*

*A. Halder, S. S. Pandey, ASM, JCAP **08**, 064 (2023)*

*S. S. Pandey, A. Halder, ASM, Phys. Rev. D **110**, 043531 (2024)*

# Outlook:

- ❖ Observations tell us that the present Universe is inhomogeneous up to scales ( $< 500 h^{-1} \text{ Mpc}$ ) [Features: Spatial volume is dominated by voids; peculiar structures at very large scales] [*Sloan Digital Sky Surveys; Giant arc  $\sim 1 \text{ Gpc}$* ]
- ❖ Cosmology is very well described by spatially homogeneous and isotropic FLRW model (*modulo recent tensions: Hubble, S-8.. ?*)
- ❖ Observational concordance comes with a price: more that 90% of the energy budget of the present universe comes in forms that have never been directly observed (DM & DE); DE not even theoretically understood
- ❖ Scope for alternative thinking without modifying GR or extending SM; application of GR needs to be more precisely specified on large scales
- ❖ Backreaction from inhomogeneities could modify the evolution of the Universe; Gravitational wave propagation  compact object parameters in GW astronomy

# Propagation of Gravitational Waves from binaries

- Several detection events of compact binary mergers since LIGO and VIRGO
- Observed GW parameters are crucial for inferring source parameters, viz., mass, merger rate
- Multi-messenger astronomy opening up new observational window to physics of BH formation & many aspects of early universe physics
- Present observation of GWs comes from sources that are well within (much smaller than) the scale of observed global homogeneity
- *Backreaction induced changes in observed GW parameters – corresponding modification in inferred source parameters*

## Problem of course-graining or averaging

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa T_{\mu\nu}$$

$$G_{\mu\nu} = \langle T_{\mu\nu} \rangle$$

*Einstein's equations: **nonlinear***

$$\langle G_{\mu\nu}(g_{\mu\nu}) \rangle = \langle T_{\mu\nu} \rangle \neq G_{\mu\nu}(\langle g_{\mu\nu} \rangle)$$

*Einstein tensor constructed from average metric tensor will not be same in general as the average of the Einstein tensor of the actual metrics*

# Different approaches of averaging

*Macroscopic gravity: (Zalaletdinov, GRG '92;'93)*

$$\langle g^{\mu\lambda} \rangle \langle R_{\lambda\nu} \rangle - \frac{1}{2} \delta_\nu^\mu \langle g^{\lambda\rho} \rangle \langle R_{\lambda\rho} \rangle + C_\nu^\mu = \kappa \langle T_\nu^\mu \rangle$$

(additional mathematical structure for covariant averaging scheme)

*Perturbative schemes: (Clarkson et al, RPP '11; Kolb, CQG '11)*

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + \delta g_{\mu\nu} \quad \quad \bar{G}_\nu^\mu + \delta G_\nu^\mu = \kappa \langle T_\nu^\mu \rangle$$

***Spatial averages : (Buchert, GRG '00; '01)\*\*\****

*Lightcone averages: (Gasperini et al., JCAP '09;'11)*

*Bottom-up approach [discrete cosmological models]: (Tavakol, PRD'12; JCAP'13)*

Using the Einstein equations:

$$3 \frac{\ddot{a}_D}{a_D} = -4\pi G \langle \rho \rangle_D + \mathbf{Q}_D + \Lambda$$

Q: Backreaction  
due to averaging

$$3H_D^2 = 8\pi G \langle \rho \rangle_D - \frac{1}{2} \langle \mathbf{R} \rangle_D - \frac{1}{2} \mathbf{Q}_D + \Lambda$$

$$0 = \partial_t \langle \rho \rangle_D + 3H_D \langle \rho \rangle_D$$

where the average of the scalar quantities on the domain  $D$  is

$$\langle f \rangle_D(t) = \frac{\int_D f(t, X^1, X^2, X^3) d\mu_g}{\int_D d\mu_g}$$

Integrability condition:

$$\frac{1}{a_D^6} \partial_t (a_D^6 \mathcal{Q}_D) + \frac{1}{a_D^2} \partial_t (a_D^2 \langle R_D \rangle) = 0$$

$\rho$  = local matter density

$R$  = Ricci-scalar

$H_D = \frac{\dot{a}_D}{a_D}$  = domain dependent Hubble rate

The kinematical backreaction  $Q_D$   
is defined as

$$Q_D = \frac{2}{3} \left( \langle \theta^2 \rangle_D - \langle \theta \rangle_D^2 \right) - 2\sigma_D^2$$

where  $\theta$  is the local expansion rate,

$\sigma^2 = 1/2 \sigma_{ij} \sigma^{ij}$  is the squared rate of shear

## Acceleration equation for the global domain $D$ :

$$\frac{\ddot{a}_D}{a_D} = \sum_{\ell} \lambda_{\ell} \frac{\ddot{a}_{\ell}(t)}{a_{\ell}(t)} + \sum_{\ell \neq m} \lambda_{\ell} \lambda_m (H_{\ell} - H_m)^2$$

## 2-scale interaction-free model (Weigand & Buchert, PRD '10):

$M$  – those parts that have initial overdensity (“Wall”)

$E$  – those parts that have initial underdensity (“Void”)

$$D = M \cup E$$

$$H_D = \lambda_M H_M + \lambda_E H_E$$

Void fraction:  $\lambda_E = \frac{|E|}{|D|}$       Wall fraction:  $\lambda_M = \frac{|M|}{|D|}$

$$\lambda_M + \lambda_E = 1$$

Acceleration equation:

$$\frac{\ddot{a}_D}{a_D} = \lambda_M \frac{\ddot{a}_M}{a_M} + \lambda_E \frac{\ddot{a}_E}{a_E} + 2\lambda_M \lambda_E (H_M - H_E)^2$$



Future evolution assuming present acceleration ★

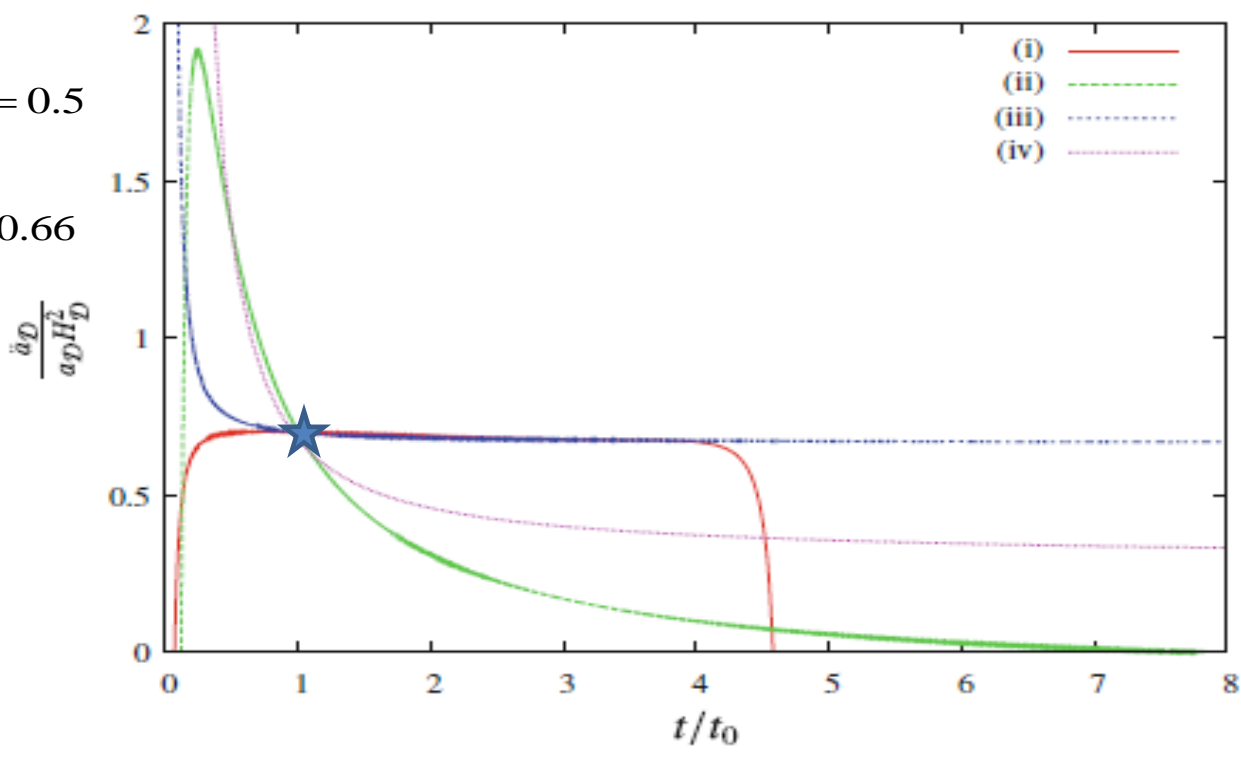
$$a_M \propto c_M t^\beta$$

$$a_E \propto c_E t^\alpha$$

Present wall fraction,  $\lambda_{M_0} = 0.09$  [Weigand & Buchert, PRD '10]

- (i)  $\alpha = 0.995, \beta = 0.5$
- (ii)  $\alpha = 1.02, \beta = 0.66$

★  
N.Bose & ASM,  
GRG (2013)



$$q = 0.6, \beta = 0.7$$

$$a_E = c_E t^\beta$$

$$a_M = \frac{q}{2q-1} (1 - \cos \theta)$$

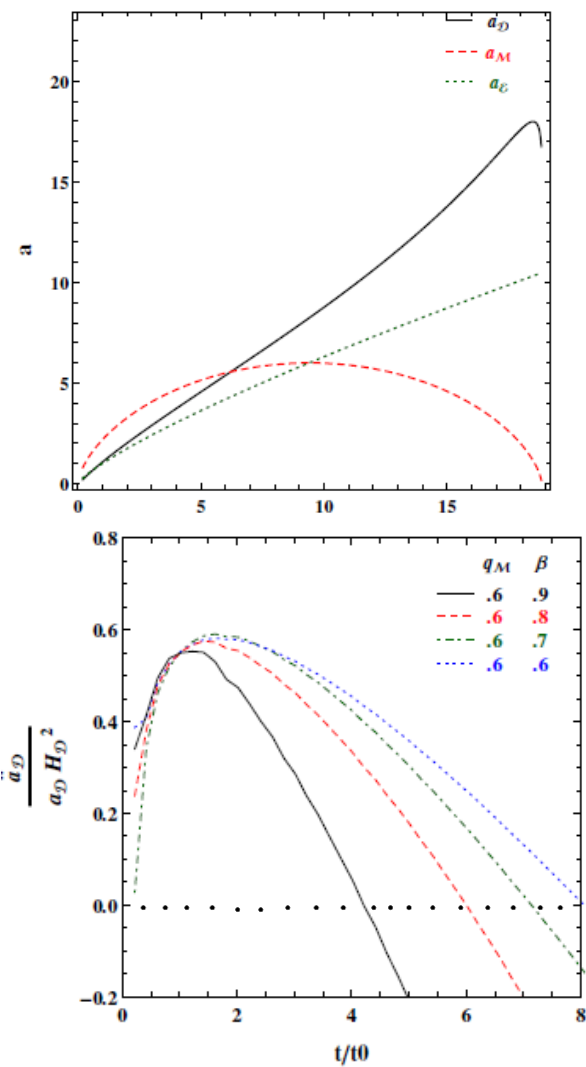
$$t = \frac{q}{2q-1} (\theta - \sin \theta)$$

$$H_D = \lambda_M H_M + \lambda_E H_E$$

*As time evolves,  $H_E$  falls off more rapidly compared to  $H_M$*

*Even though the wall occupies a tiny fraction of the total volume, the decrease of  $\lambda_M$  is more than compensated by the comparative evolution of  $H_E$  and  $H_M$*

*Decelerating future evolution !*



*Effective perfect fluid E-M tensor in the Backreaction formalism:*

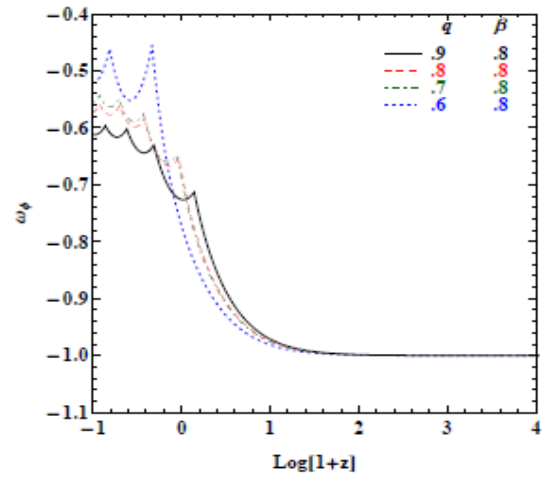
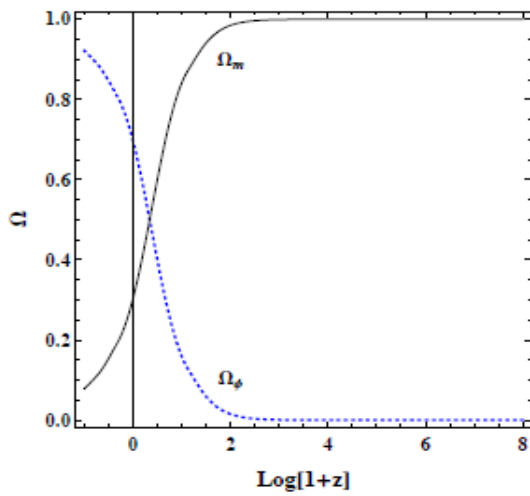
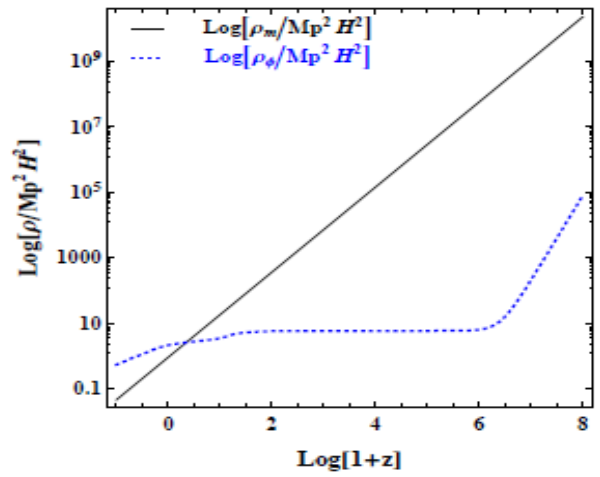
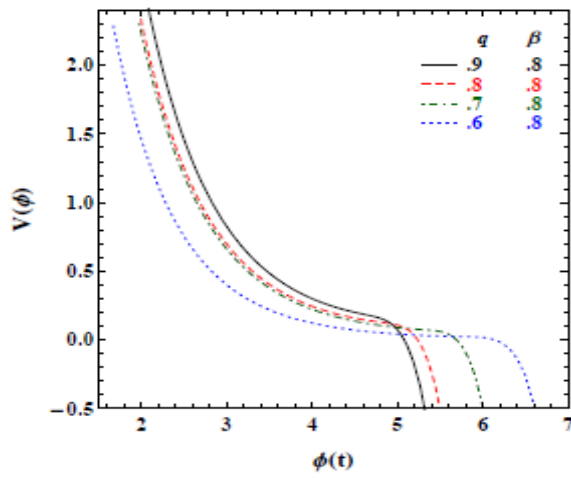
$$\rho_{eff}^D = \langle \rho \rangle_D - \frac{1}{16\pi G} Q_D - \frac{1}{16\pi G} \langle R \rangle_D$$
$$P_{eff}^D = -\frac{1}{16\pi G} Q_D + \frac{1}{16\pi G} \langle R \rangle_D$$

*Buchert equations recast in standard Friedman form:*

$$3 \frac{\dot{a}_D}{a_D} = -4\pi G (\rho_{eff}^D + 3P_{eff}^D) + \Lambda$$
$$3H_D^2 = 8\pi G \rho_{eff}^D + \Lambda$$
$$\dot{\rho}_{eff}^D + 3H_D (\rho_{eff}^D + P_{eff}^D) = 0$$

*corresponding to energy density and pressure of effective global scalar field at scales much larger than the scale of inhomogeneities*

# Scalar field dynamics



*Thawing quintessence [Recent DESI BAO Observations !]*

*Observational Impact of  
Inhomogeneities:*

*Multi-messenger Astronomy*

# Backreaction model

(2-scale void-wall) [Pandey, Ali, Sarkar, ASM '22]

$$t=t_0\left(\frac{\phi-\sin\phi}{\phi_0-\sin\phi_0}\right),$$

$$a_o=\frac{fo^{1/3}}{2}(1-\cos\phi),$$

$$a_u=\frac{fu^{1/3}(\phi_0-\sin\phi_0)}{\pi t_0}t^\beta$$

$$a_{\mathcal{D}}=\left(\frac{a_u^3+a_o^3}{a_{u,0}^3+a_{o,0}^3}\right)^{1/3}$$

$$H_{\mathcal{D}}=H_u\frac{a_u^3}{a_u^3+a_o^3}+H_o\frac{a_o^3}{a_u^3+a_o^3}$$

$$\mathcal{Q}_{\mathcal{D}}=\mathcal{Q}_o+\mathcal{Q}_u+6fo(1-fo)(H_o-H_u)^2$$

$$\mathcal{Q}_o=0\qquad\mathcal{Q}_u=0.$$

# Light propagation in backreaction model

Angular diameter distance

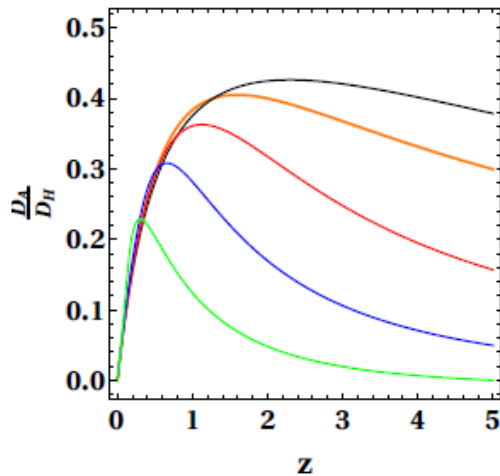
$$D_A = \frac{c}{1+z_1} \int_0^{z_1} \frac{dz}{H(z)}$$

Covariant scheme:  
[Rasanen '09]

$$1+z = \frac{1}{a_{\mathcal{D}}},$$

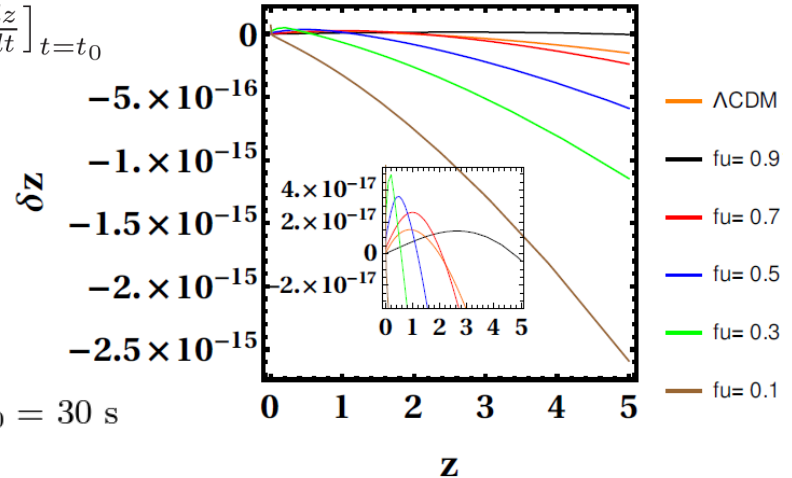
$$H_{\mathcal{D}} \frac{d}{dz} \left( (1+z)^2 H_{\mathcal{D}} \frac{dD_A}{dz} \right) = -\frac{4\pi G}{c^4} \langle \rho_{\mathcal{D}} \rangle D_A$$

Change in observed redshift during the time interval of observation: Red-shift drift



$$\delta z = \left[ \frac{dz}{dt} \right]_{t=t_0}$$

$$\delta t_0 = 30 \text{ s}$$



Angular diameter versus red-shift

Departure from  $\Lambda$ CDM in terms  
model parameters

# Gravitational wave amplitude

*Binaries in early inspiral stage  
(Keplerian approximation)*

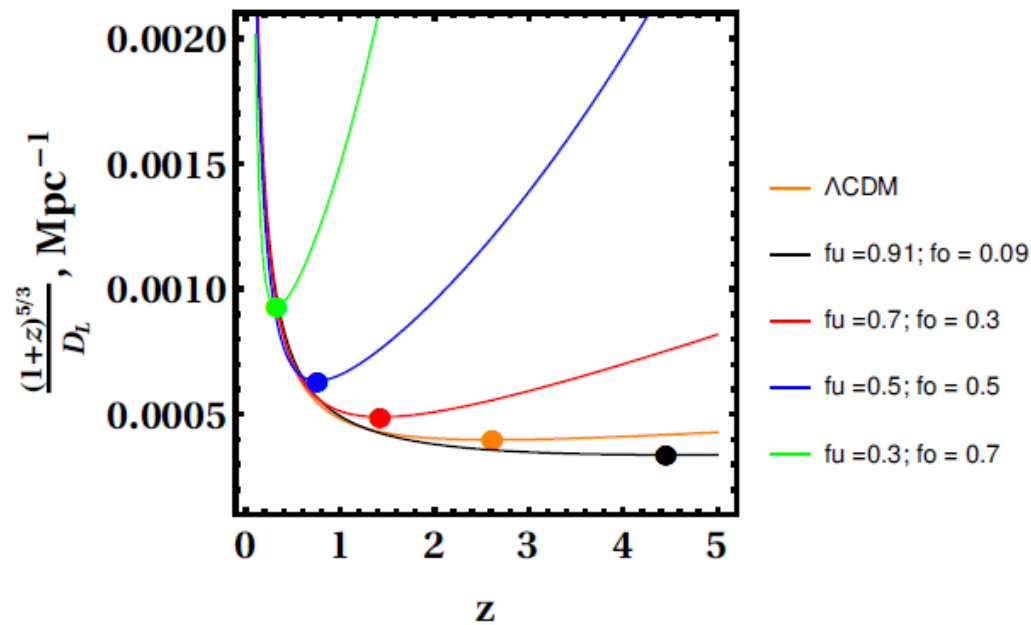
$$h_{\times} = \frac{G^{5/3}(1+z)^{5/3}}{D_L c^4} \frac{m_1 m_2}{(m_1 + m_2)^{1/3}} (-4\omega^{2/3}) \sin 2\omega t$$

$$D_L = (1+z)D$$

*Red-shift dependent part  
of GW amplitude*

$$(1+z)^{5/3}/D_L$$

*Deviations from  $\Lambda$ CDM  
get amplified at higher  $z$*





# Change in gravitational wave observables

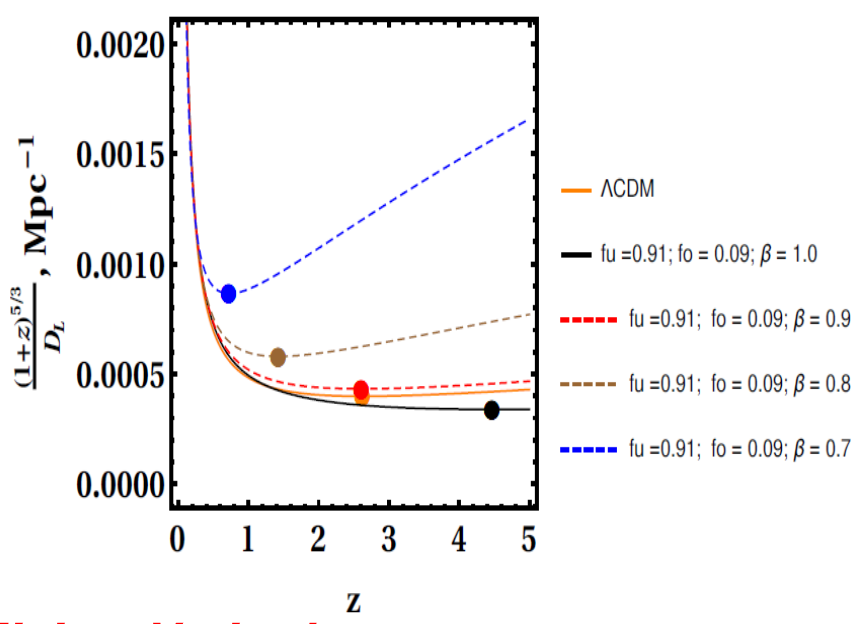
\* Amount of change depend upon effect of inhomogeneities  
in the path of propagation – model parameters  $f_u, f_o$ , and  $\beta$

## Red-shift minima (model dependence)

$(1+z)^{5/3}/D_L$  has a minimum at  $z_{\min}$  S. S. Pandey, A. Sarkar, A. Ali, ASM, JCAP 06, 021 (2022)

$$(1+z_{\min}) \left[ \frac{d}{dz} \ln[D_L] \right]_{z=z_{\min}} = \frac{5}{3}$$

Independent of binary characteristics,  
cosmological model, GW detector  
[Rosado et al., PRL 116, 101102 (2016)]



Minima Varies !

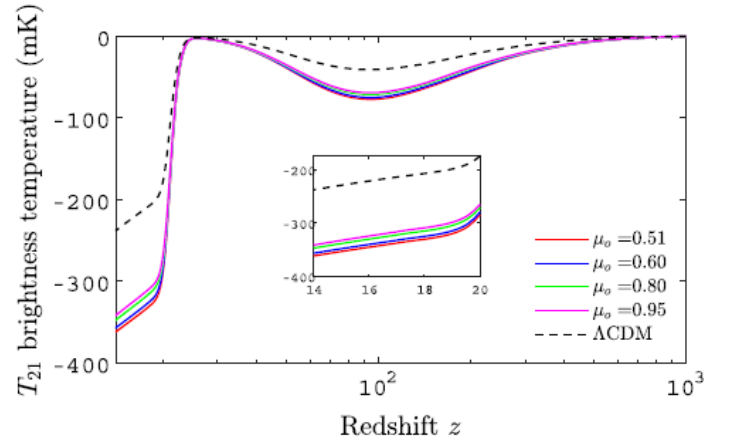
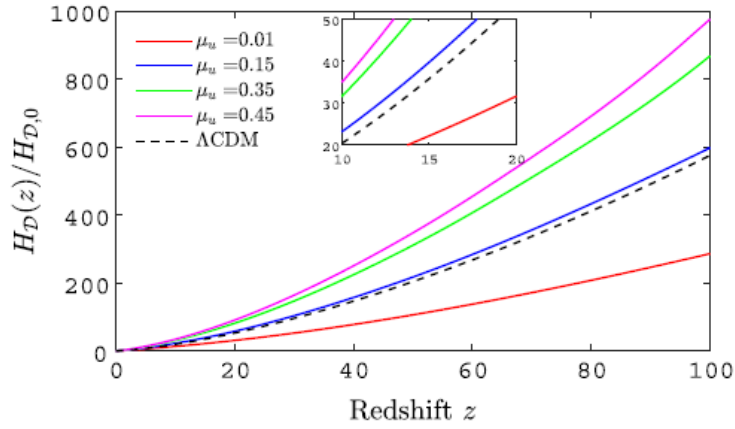
# Analyzing the 21-cm signal brightness temperature in the Universe with inhomogeneities

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**Multidomain model** [A. Halder, S. S. Pandey, ASM, JCAP **08**, 064 (2023)]

*Overdense regions*

$$a_{o_i} = \frac{q_{o_i}}{2q_{o_i} - 1} (1 - \cos \phi)$$

$$t_i = \frac{q_{o_i}}{2q_{o_i} - 1} (\phi - \sin \phi)$$

*Underdense regions*

$$a_{u_i} = c_{u_i} t^{\beta_i}$$

*Global acceleration:*

$$\frac{\ddot{a}_{\mathcal{D}}}{a_{\mathcal{D}}} = \left( \sum_i -\lambda_{o_i} q_{o_i} H_{o_i}^2 \right) + \left( \sum_j \lambda_{u_j} \frac{\beta(\beta - 1)}{t^2} \right) + \left( \sum_k \sum_l \lambda_k \lambda_l (H_l - H_k)^2 \right)$$

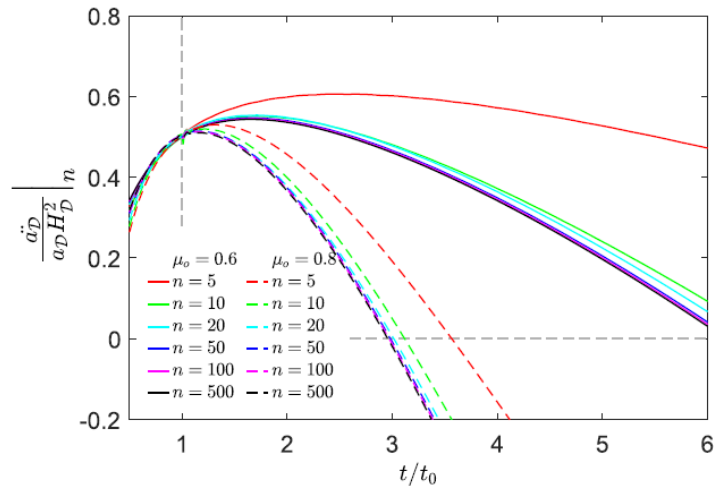
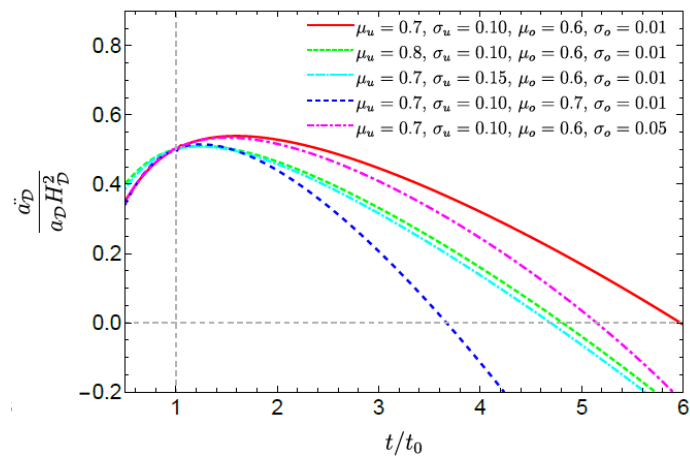
*Gaussian profile for volume fractions*

$$\lambda_{u_i,0} = \frac{N_u}{\sigma_u \sqrt{2\pi}} e^{-(\beta_i - \mu_u)^2 / 2\sigma_u^2}$$

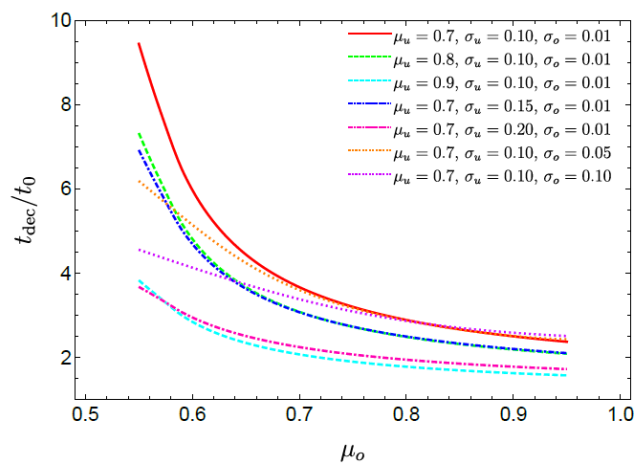
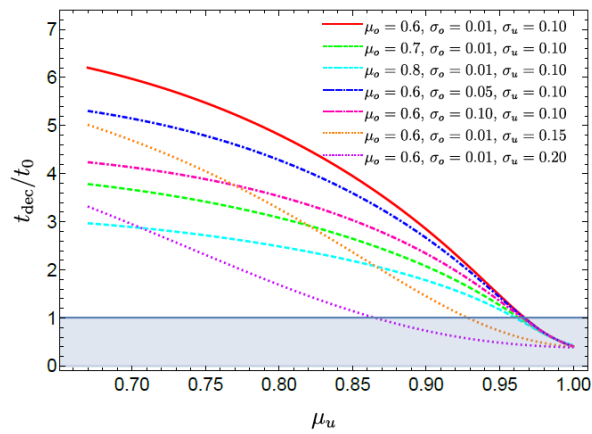
$$\lambda_{o_i,0} = \frac{N_o}{\sigma_o \sqrt{2\pi}} e^{-(q_{o_i} - \mu_o)^2 / 2\sigma_o^2}$$

# Evolution in multidomain model *(No Big-rip !)*

## Global acceleration



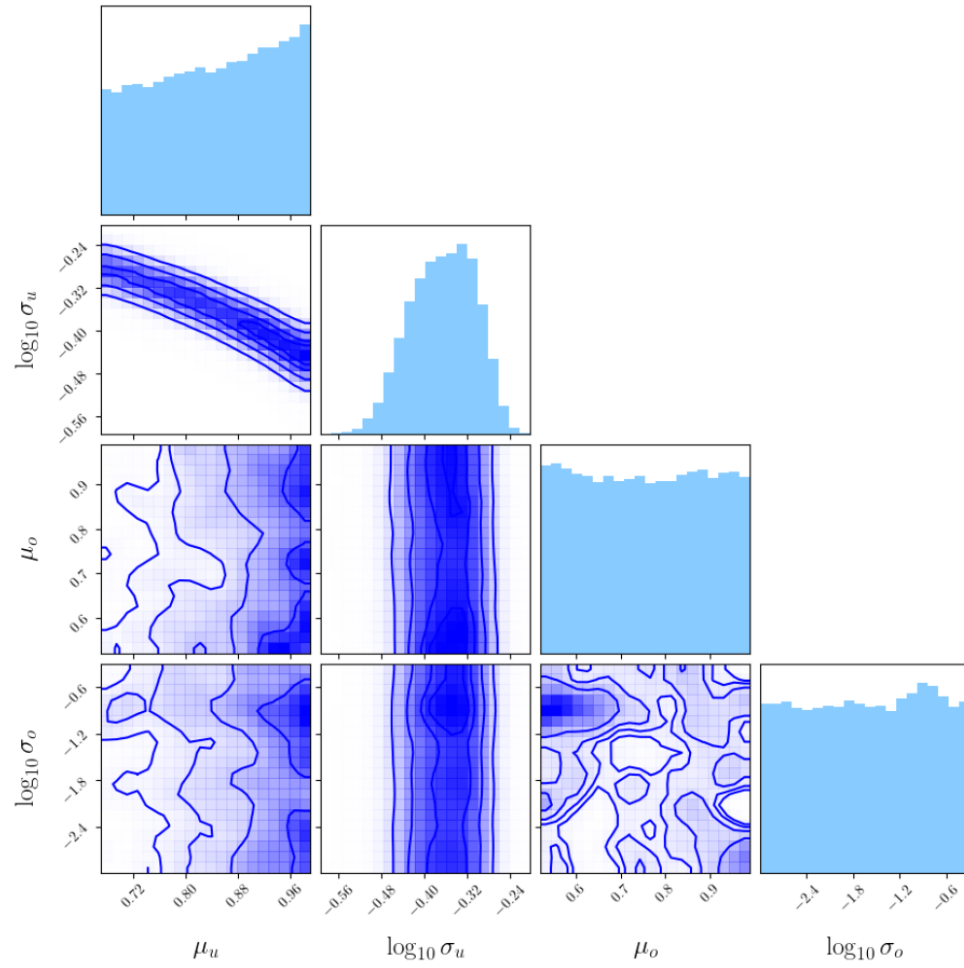
## Onset of future decelerating phase $t_{\text{dec}}$



# Observational constraints *[multi-scale Gaussian model]*

[A. Halder... ASM, JCAP '23]

*(Bayesian analysis using Union 2.1 Supernova Ia data)*



# Effect of inhomogeneities: gravitational wave propagation (Summary)

[A. Ali, ASM, JCAP 01, 054 (2017); S. S. Pandey, A. Sarkar, A. Ali, ASM, JCAP 06, 021 (2022); A. Halder, S. S. Pandey, ASM, JCAP 08, 064 (2023); S. S. P., A. H., ASM, PRD 110. 043531 (2024)]

- Effect of backreaction due to inhomogeneities on the future evolution of the accelerating universe (*Spatial averaging in the **Buchert framework***)
- The global homogeneity scale (or cosmic event horizon) impacts the role of inhomogeneities on the evolution, causing the acceleration to slow down significantly with time.
- Backreaction could be responsible for a decelerated era in the future.  
(***Avoidance of big rip !***) Possible within a small region of parameter space
- Analogous scalar field cosmology: Form of potential fixed by backreaction model; Observational constraints from data analysis
- Dip in 21-cm signal (explanation for EDGES result !)
- Modification of binary GW parameters due to backreaction from inhomogeneities. Clear signature of red-shift drift. ***Significance in Multimessenger Astronomy & EU Physics***
- Effect may be tested in more realistic models, e.g., models with no ansatz for subdomains, & other schemes of backreaction