

# Dark sector at BESIII



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## **Outline**

- ✓ Motivation
- ✓ BESIII experiment
- ✓ Status of Dark sector searches at BESIII
  - Search for a light Higgs boson  $A^0$  in radiative  $J/\psi$  decay
  - Search for an Axion-like particle
  - Dark photon (massive and massless)
  - Search for invisible decays of  $K_S^0$  meson
  - Dark baryon
  - Search for invisible muon philic scalar  $X_0$  or vector  $X_1$  via  $J/\psi \to \mu^+\mu^-+$  invisible
- ✓ Summary

## **Motivation**

- **Standard Model (SM) is incredibly successful, it is tested by experiments**
- However, it can't be quantified as a theory of everything

#### facing some tensions:

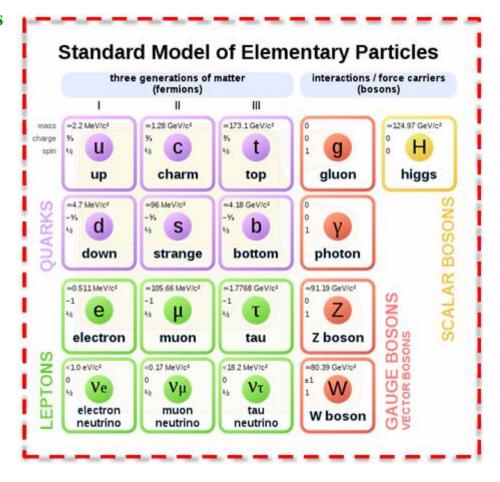
Naturalness and stability, g-2, W mass, R<sub>K</sub>, R<sub>D</sub>, R<sub>D\*</sub>, ....

#### Can not explain:

Existence & mechanism of dark matter and dark energy

Baryon asymmetry of the universe

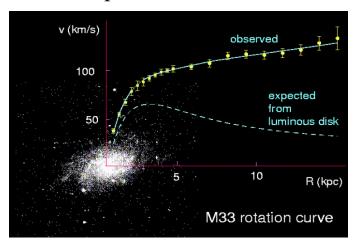
Neutrino masses and oscillations, hierarchy



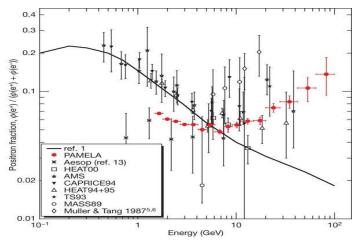
Real opportunity to search for new physics beyond the Standard Model

## **Motivation**

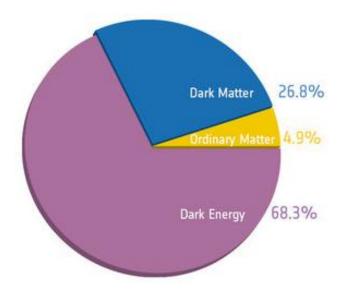
- **Standard Model (SM) is incredibly successful but not complete!** 
  - ☐ Extensions of the SM needed to solve several outstanding issues, including the missing description of Dark Matter (DM)
  - Why DM?
    - Amounts 27% of the total matter density of the universe
    - Not interact with strong and electromagnetic interactions, it presence so far can be inferred via the gravitational effects only.
    - Explain the features of recent astrophysical observations



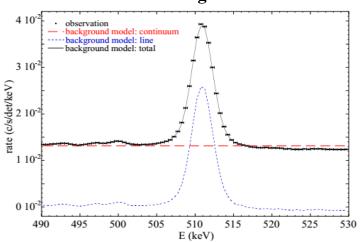
arXiv:astro-ph/0403324



PAMELA: Positron fraction O Adriani et al., Nature 458 (2009) 607

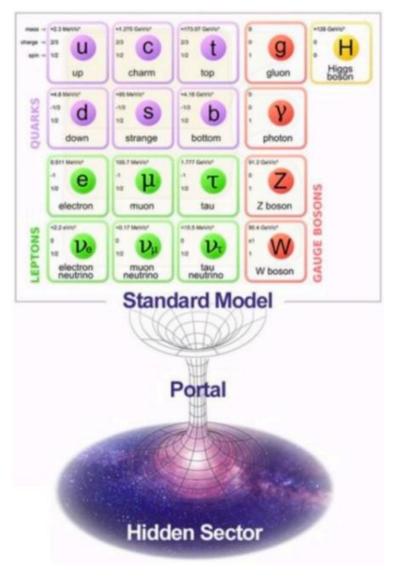


#### **SPI/Integral**



P. Jean et al., A&A 407, L-55-L58 (2003)

## **Coupling of DM with Standard Model**



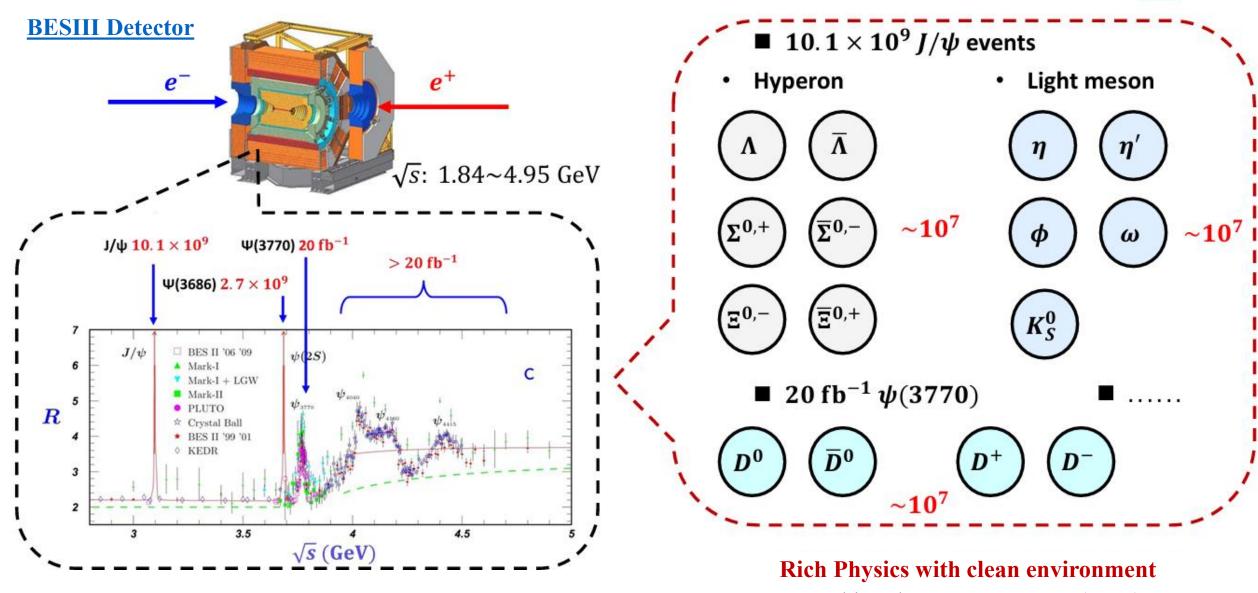
- Dark matter has not seen yet in particle physics experiments.
  - SM can't explain DM  $\Longrightarrow$  Extend to the SM to include Dark matter
  - One of the simplest models is "DM hidden sector" that allows the coupling between DM and SM particles via the so called "portals"

#### R. Essig et al., arXiv:1311.0029 (2013)

$$\mathcal{L} \supset \left\{ \begin{array}{ll} -\frac{\varepsilon}{2\cos\theta_W} B_{\mu\nu} F'^{\mu\nu} \,, & \text{vector portal} \\ (\mu\phi + \lambda\phi^2) H^\dagger H \,, & \text{Higgs portal} \\ y_n L H N, & \text{neutrino portal} \\ \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, & \text{axion portal} \\ \end{array} \right. \quad \begin{array}{ll} \textbf{A' kinetic mixing with } \gamma, \textbf{Z} \\ \textbf{Dark Higgs (mixes with SM Higgs)} \\ \textbf{Sterile neutrino} \\ \textbf{Axion, coupling to DM} \end{array}$$

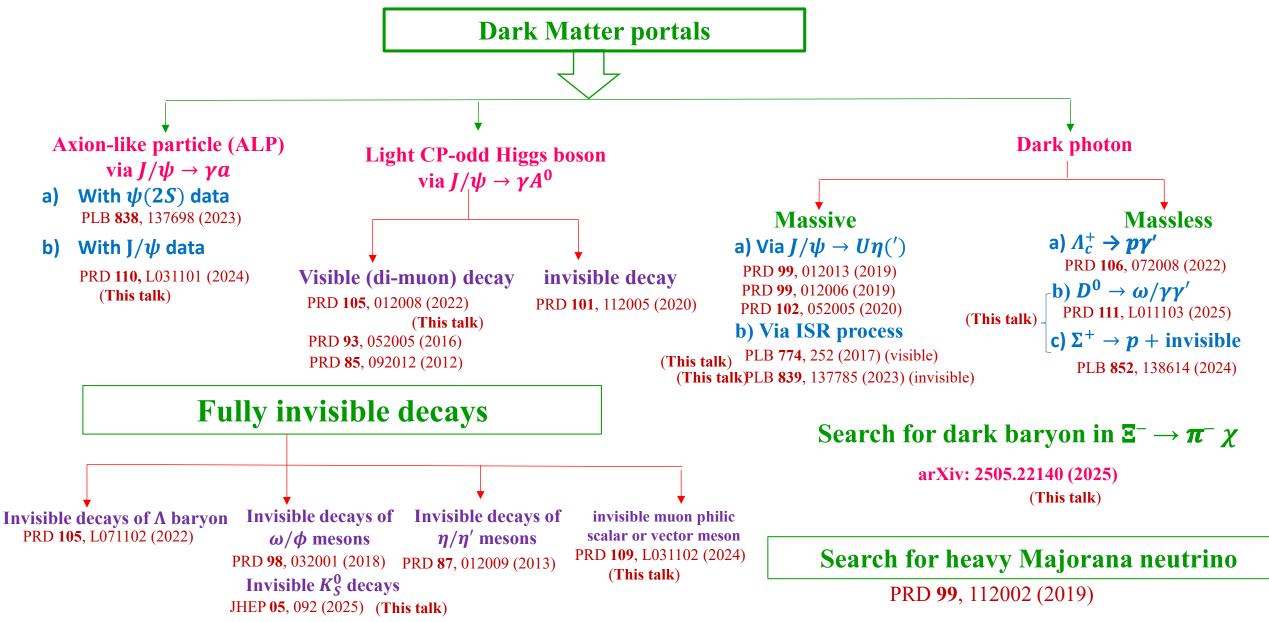
• Can be accessible by high intensity e<sup>+</sup>e<sup>-</sup> collider experiments, such as BESIII experiment, if their masses are a few GeV

## **BESIII Experiment**



Chin. Phys. C 44, 040001 (2020)

## Status of Dark sector searches at BESIII



## Light Higgs boson A<sup>0</sup> search

#### PRD 105, 012008 (2022)

- A light Higgs boson is predicted by many extensions of Standard Model, such as Next-to-Minimal Supersymmetric Standard Model (NMSSM).
  - ➤ NMSSM contains a total of three CP-even, two CP-odd and two charged Higgs bosons.
    - $\triangleright$  The lighter state of the  $A^0$  is defined as:

➤ Coupling of fermions and the CP-odd Higgs A<sup>0</sup>

$$L_{\text{int}}^{f\overline{f}} = -\cos\theta_A \tan\beta \frac{m_f}{v} A^0 \overline{d}(i\gamma_5) d, \quad d = d, \, s, \, b, \quad e, \, \mu, \, \tau$$

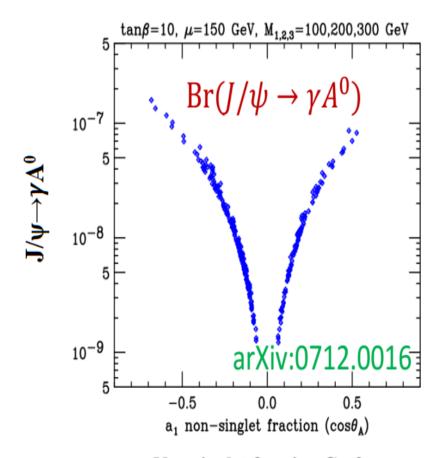
$$L_{\text{int}}^{f\bar{f}} = -\cos\theta_A \cot\beta \frac{m_f}{v} A^0 \bar{u}(i\gamma_5) u, \quad u = u, \quad \boldsymbol{c}, \quad t, \quad v_e, \quad v_\mu, \quad v_\tau$$

$$\tan \beta = \frac{v_u}{v_v}$$
 Ratio of the VEVs of the up and down-types of Higgs doublets

E. Fullana et. al, Phys. Lett. B 653, 67 (2007)

 $\triangleright$  Can be detectable via radiative decays of J/ $\psi$  and  $\Upsilon(1S)$ 

[Phys. Rev. Lett. **39**, 1304 (1977)]

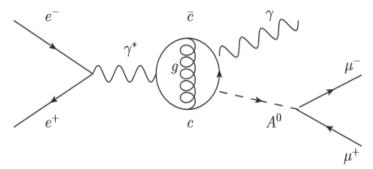


Non singlet fraction  $Cos\theta_A$ 

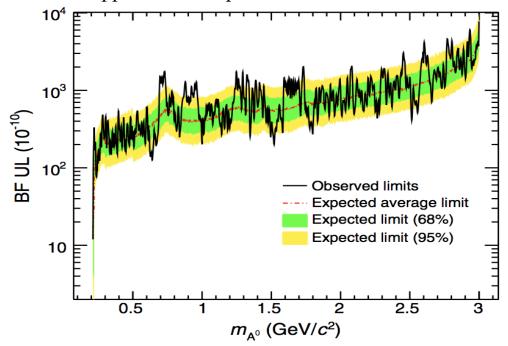
## Light Higgs boson $A^0$ search in radiative $J/\psi$ decay

Expected  $B(J/\psi \to \gamma A^0) \sim 10^{-9} - 10^{-7}$  [PRD 76, 051105 (2007)]

PRD 105, 012008 (2022)



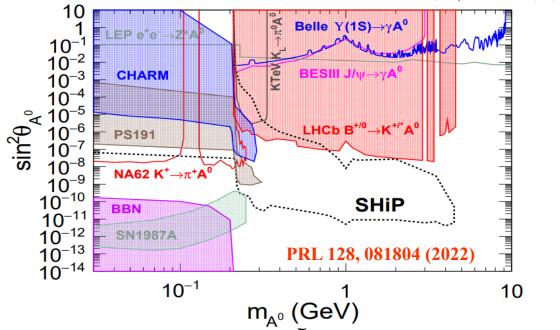
No evidence of A<sup>0</sup> production is found and set 90% confidence level upper limits on product BFs.



 $\triangleright$  Use 9 billion J/ψ events collected by BESIII experiment to perform this study.

### Mixing angle $(sin\theta_{A^0})$

$$\frac{\mathcal{B}(\Upsilon(1S) \to \gamma A^0)\mathcal{B}(A^0 \to \text{hadrons})}{\mathcal{B}(\Upsilon(1S) \to \ell^+\ell^-)} = \sin^2\theta_{A^0} \frac{G_F m_b^2}{\sqrt{2}\pi\alpha} \sqrt{(1 - \frac{m_{A^0}^2}{m_{\Upsilon(1S)}^2})}$$



Our result in the low-mass region is better than recent **BELLE** 

measurement

## Search for an Axion-like particle

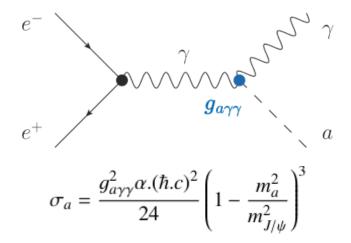
Phys. Lett. B 753, 482 (2016)

An Axion-like particle (ALP), a

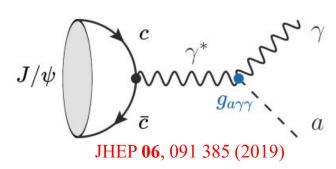
Phys. Rev. D 110, L031101 (2024)

- is a pseudo-scalar particle
- introduced by the spontaneous breaking of Peccei-Quinn symmetry to solve the strong CP problem of the QCD Phys. Rev. Lett. 40, 223 (1978); Phys. Rev. Lett. 40, 279 (1978) Phys. Rev. Lett. 38, 1440 (1977); Phys. Rev. D 16, 1791 (1977)
- Predicted by many models beyond the SM and proposed to be a cold DM candidate.
- ALP production at e<sup>+</sup>e<sup>-</sup> colliders
  - ALP-Strahlung process

Phys. Rev. D 52, 1755 (1995)



#### Radiative decay process

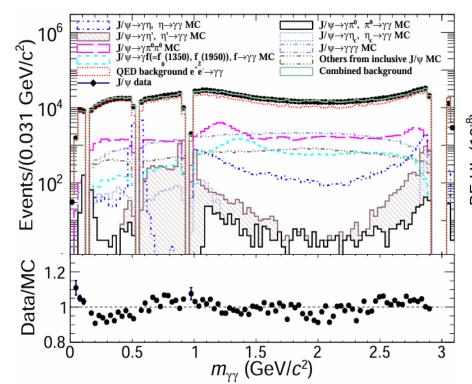


$$\mathcal{B}(J/\psi \to \gamma a) = \frac{m_{J/\psi}^2}{32\pi\alpha} g_{a\gamma\gamma}^2 \left(1 - \frac{m_a^2}{m_{J/\psi}^2}\right)^3 \mathcal{B}(J/\psi \to e^+e^-)$$

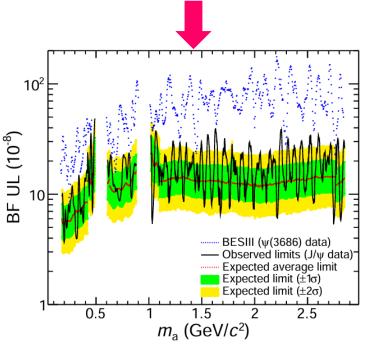
## Search for ALP via radiative $J/\psi$ decays at BESIII

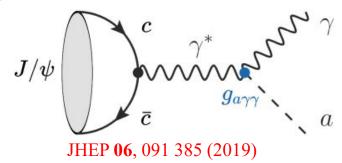
Phys. Rev. D 110, L031101 (2024

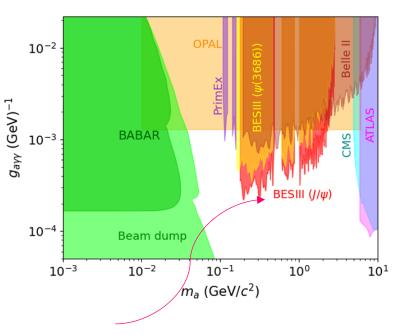
- $10^{10} J/\psi$  events
- Extract signal from  $M_{\gamma\gamma}$  distribution
- Maximum signal significance:  $< 3\sigma$



• UL on the BF of  $\mathcal{B}(J/\psi \to \gamma a) \times \mathcal{B}(a \to \gamma \gamma)$ (3. 6~53. 1) × 10<sup>-8</sup>@ .95% CL.



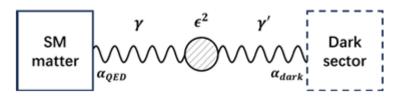




New stringent constraints on ALP-photon coupling for  $0.18 \le m_a \le 2.85 \text{ GeV}$ 

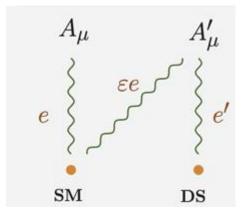
## Dark photon

Simplest extension of the SM  $\Rightarrow$  An extra Abelian gauge group,  $U(1)_D \Rightarrow$  dark photon



### Massive dark photon

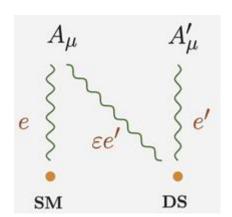
arises when the symmetry of the additional Abelian gauge group is spontaneously broken



- Massive dark photon
- Coupling with SM fermion
- Strong constraint

## Massless dark photon

Symmetry remains unbroken



- B. Batell, et al, PRD **79**, 115008 (2009);
- R. Essig, et al, PRD **80**, 015003 (2009)

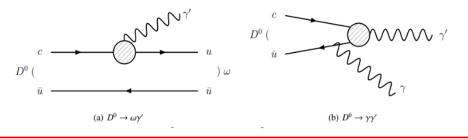
- Massless dark photon
- No direct coupling with SM fermion
- Less constraint
- Also, important role in dark sector

A dimension-six operator has been proposed to provide a connection between SM fermions and the massless dark photon

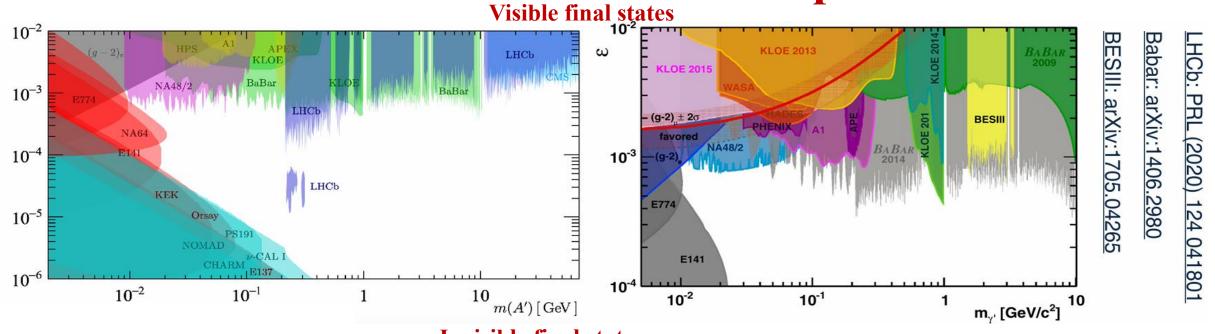
$$\mathcal{L}_{ ext{NP}} = rac{1}{\Lambda_{ ext{NP}}^2} (C_{jk}^U ar{q}_j \sigma^{\mu
u} u_k ilde{H} + C_{jk}^D ar{q}_j \sigma^{\mu
u} d_k H + C_{jk}^L ar{l}_j \sigma^{\mu
u} e_k H + ext{H.c.}) F'_{\mu
u}$$

- Naturally allow the FCNC coupling
- Less background and higher sensitivity

PRL **94**, 151802 (2005)

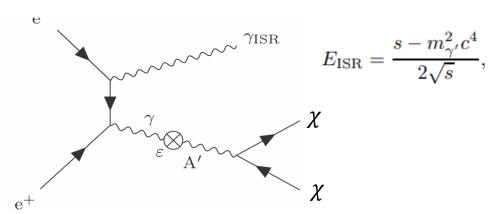


## **Current status of massive dark photon searches**

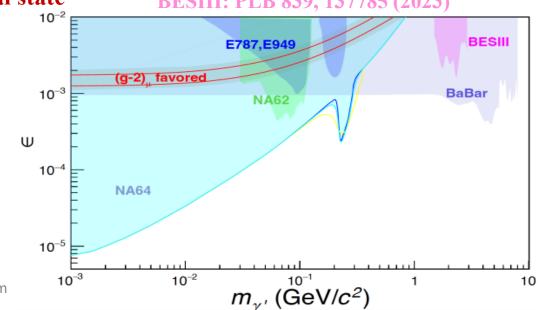


#### **Invisible final state**

[: PLB 839, 137785 (2023)



BaBar: Phys. Rev. Lett. 119, 131804 (2017)

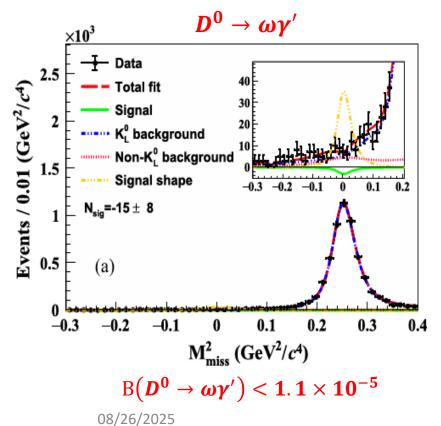


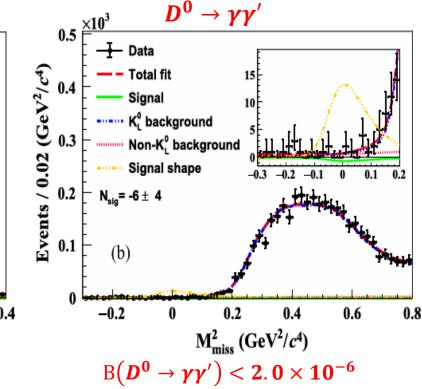
# Search for a massless dark photon in $D^0 \to \omega/\gamma\gamma'$

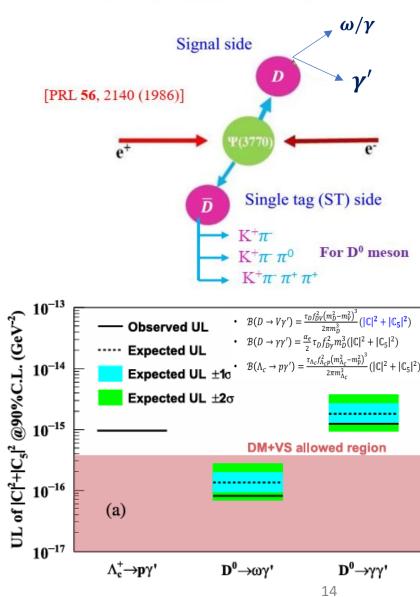
PRD 111, L011103 (2025)

- Search is based on 7.9 fb<sup>-1</sup>  $\psi(3770)$  data using a double tag (DT) technique
- The signals of the massless dark photon are extracted from a fit on the distribution of missing mass square

$$M_{\rm miss}^2 = |p_{\rm c.m.s.} - p_{\bar{D}^0} - p_{\omega(\gamma)}|^2/c^4$$







# $\Sigma^+ \rightarrow p$ + invisible and QCD axion

PLB 852, 138614 (2024)

•  $s \rightarrow d\nu\bar{\nu}$  is a FCNC process and highly suppressed by GIM mechanism

PRD 2, 1285 (1970)

• BF  $< 10^{-11}$ 

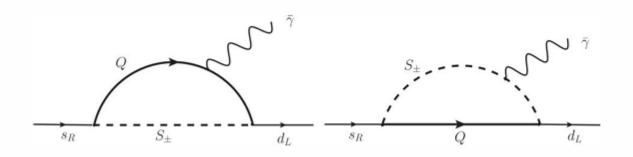
Rep. Prog. Phys. 86, 016201 (2023)

 $\Sigma^+ \rightarrow p +$  Decay to BSM particle invisible

$$m_a \propto 1/f_a \Rightarrow m_a \ll 1 \text{eV}$$

$$\frac{\Gamma(\Sigma^{+} \to pa)}{16\pi} = \frac{M_{\Sigma^{+}}^{3}}{16\pi} \left(1 - \frac{M_{p}^{2}}{M_{\Sigma^{+}}^{2}}\right)^{3} \left(\frac{(-1)^{2}}{|F_{sd}^{V}|^{2}} + \frac{0.34^{2}}{|F_{sd}^{A}|_{-}^{2}}\right)$$

Our measurement can set limits on axion-fermion effective decay constants

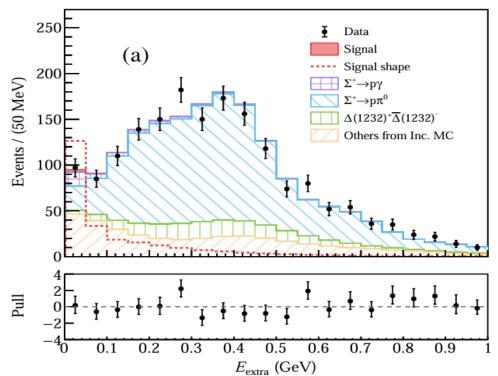


PRD 102 015023 (2020)

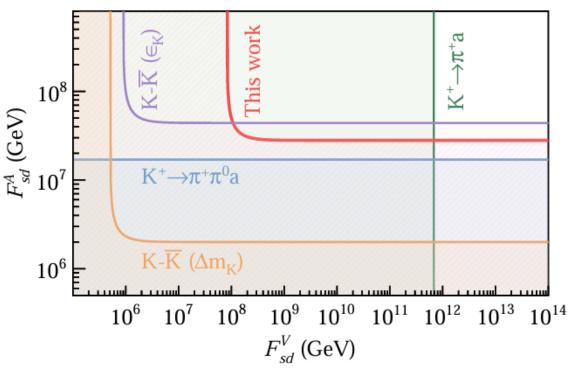
Massless dark photon

# Search for $\Sigma^+ \rightarrow p$ + invisible at BESIII

- PLB 852, 138614 (2024)
- Search is performed via  $J/\psi \to \Sigma^+ \bar{\Sigma}^-$  using 10 billion  $J/\psi$  events using double tag technique
- Invisible particle with mass hypothesis of zero
- Total energy deposited in EMC by extra photons is used to extra the signal



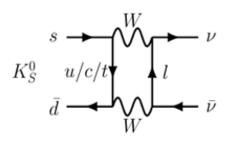
 $B(\Sigma^+ \to p + invisible) < 3.2 \times 10^{-5}$  @ 90% C.L.



Competitive limits on the axial-vectorial part of axion-fermion effective decay constant  $F_{sd}^{A}$ 

# Invisible $K_S^0$ decays

✓ SM decays

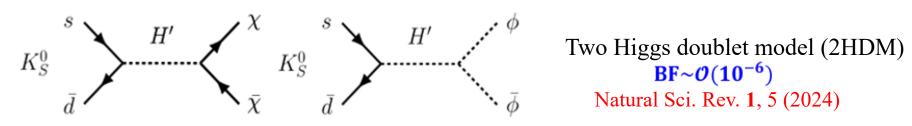


Very rare (BF <  $10^{-16}$ )  $K_S^0$  u/c/t  $\bar{\nu}$ Very rare (BF <  $10^{-16}$ )

FCNC and Helicity suppression

Phys. Rev. D **91**, 015004 (2015)

✓ Decay to DM particles



**✓** Ordinary matter particle oscillation

$$K_S^0 \leadsto K_S^{0\prime}$$

Mirror matter model BF~ $\mathcal{O}(10^{-6})$ 

arXiv: 2006.10746

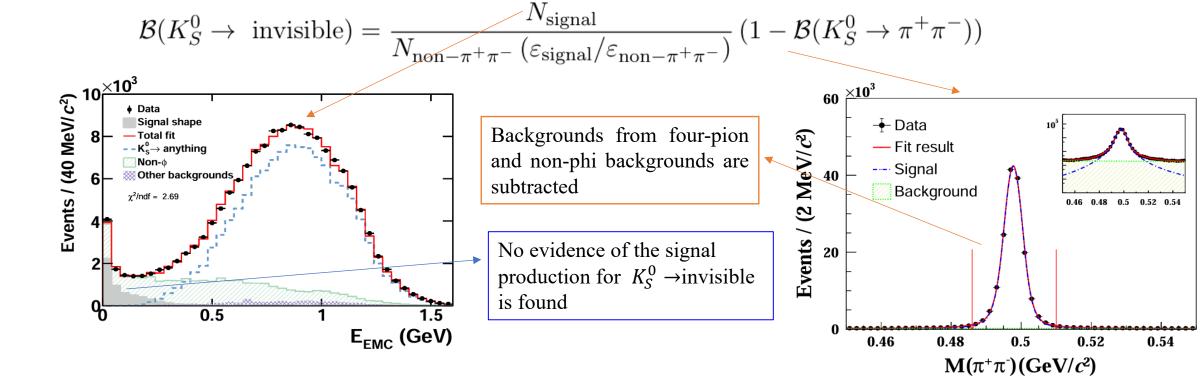
✓ Input for CP test

Bell-Steinberger relation **connects CPTV** to the amplitudes of all decay channels of neutral kaons. BUT currently assumes no invisible modes

# Invisible $K_S^0$ decays

JHEP 05, 092 (2025)

- $\triangleright$  Search is based on 10 billion  $J/\psi$  events collected by the BESIII detector via  $J/\psi \to \phi K_S^0 K_S^0$  decay.
- $\rightarrow J/\psi \rightarrow \phi K_S^0 K_L^0$  is forbidden by C-parity conservation.
- $\succ K_S^0$   $\rightarrow$  invisible decay rate can be calculated as,



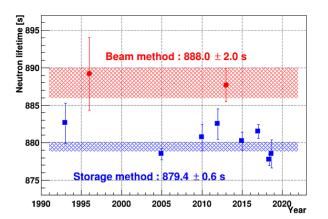
 $B(K_S^0 \rightarrow \text{invisible}) < 8.4 \times 10^{-4}$  at the 90% CL (First direct measurement).

## **Dark baryon**

**Coincidence issue:** 

Similarity between DM and baryon densities:  $\rho_{DM} \approx 5.4 \rho_{baryon}$ 

**❖** Neutron lifetime puzzle



arXiv: 2505.22140 (2025)

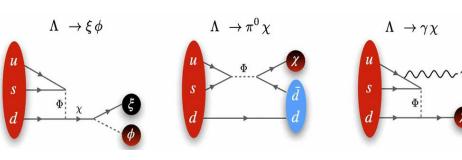
Potential connection between their origins **DM may have non-zero baryon number** 



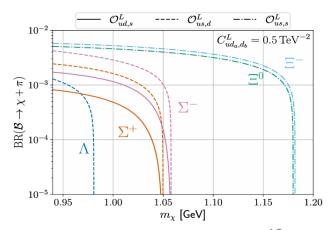
$$\tau_n^{\text{beam}} = \frac{\tau_n^{\text{bottle}}}{\text{Br}(n \to p + \text{ anything })}$$

$$\mathcal{B}(n \to \text{dark}) \sim 1\%$$

- Motivates the existence of dark baryon
- ❖ B-Mesogenesis mechanism: Can explain the symmetry between visible mater and antimatter and origin of DM.
- **❖** Hyperon dark decays



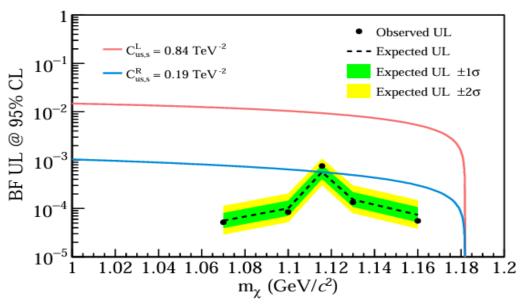
PRD **105**, 115005 (2022)



# Search for dark baryon in $\Xi^- \to \pi^- \chi$

#### **Analysis strategy:**

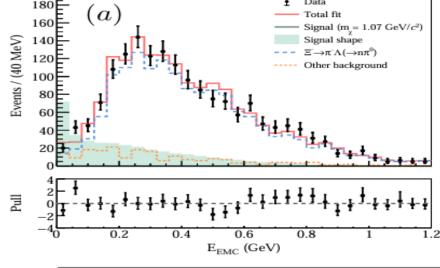
- $J/\psi \to \overline{\Xi}^+\Xi^-$  from 10 billion  $J/\psi$  events
- Double tag method:  $\overline{\Xi}^+ \to \overline{\Lambda}\pi^+, \overline{\Lambda} \to \overline{p}\pi^+, \Xi^- \to \pi^- + \chi$
- $\chi$  is the dark baryon with an invisible signature with masses of 1.07, 1.10,  $m_{\Lambda}$ , 1.13, 1.13 GeV/c<sup>2</sup>
- The invisible signal should have EMC energy deposit peaking at zero
- No evidence of significant signal events

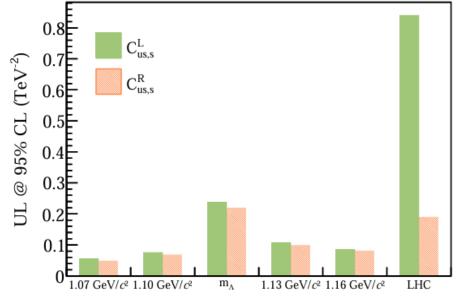


Corresponding Wilson coefficients  $C_{uss}^L$  and  $C_{uss}^R$  are more stringent than the previous limits from the LHC searches for the colored mediators

90% C.L. UL on B( $\Xi^- \to \pi^- + \chi$ ) varies from  $(4.5 - 76) \times 10^{-5}$ 

### 





# Search for an invisible muon philic scalar $X_0$ or vector $X_1$ via $J/\psi \to \mu^+\mu^-+$ invisible PRD 109, L031102 (2024)

- A new type of massive vector meson  $X_1$  or scalar boson  $X_0$  may appear in SM extension of the anomaly free gauged U(1) or  $U(1)_{L_{\mu}-L_{\tau}}$  model.
- They only couple to the second or third generations of leptons  $(\mu, \nu_{\mu}, \tau, \nu_{\tau})$  with the coupling strength  $g'_{0,1}$ .
- The  $X_{0.1}$  can contribute to the muon anomalous magnetic moment and explain the  $(g-2)_{u}$  anomaly.

$$\Delta a_{\mu}^{scalar} = \frac{g_0^2}{8\pi^2} \int_0^1 dx \frac{m_{\mu}^2 (1-x)(1-x^2)}{m_{\mu}^2 (1-x)^2 + m_{Z'}^2 x}$$

$$\Delta a_{\mu}^{vector} = \frac{g_1^2}{8\pi^2} \int_0^1 dx \frac{2m_{\mu}^2 x (1-x)^2}{m_{\mu}^2 (1-x)^2 + m_{Z'}^2 x}$$

arXiv:1610.06587 (2016)

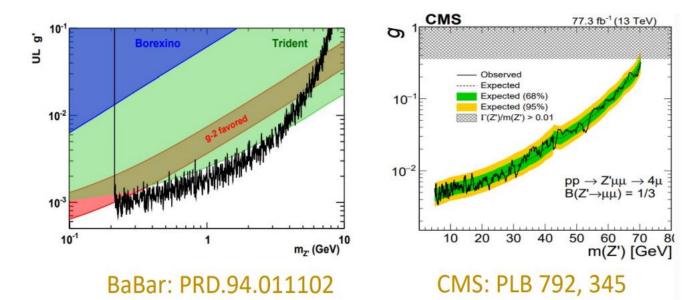
• Can be accessible via  $J/\psi \to \mu^+ \mu^- X_{0,1}$ 

JHEP 10 (2020)

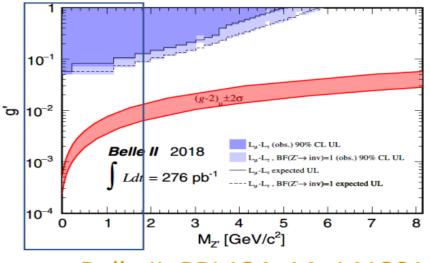
08/26/2025

# Search for an invisible muon philic scalar $X_0$ or vector $X_1$ via $I/\psi \to \mu^+\mu^-+$ invisible PRD 109, L031102 (2024)

- Current experimental constraints:
  - The g' space with  $Z' \rightarrow \mu^+ \mu^-$



• The g' space with  $Z' \rightarrow invisible$ 



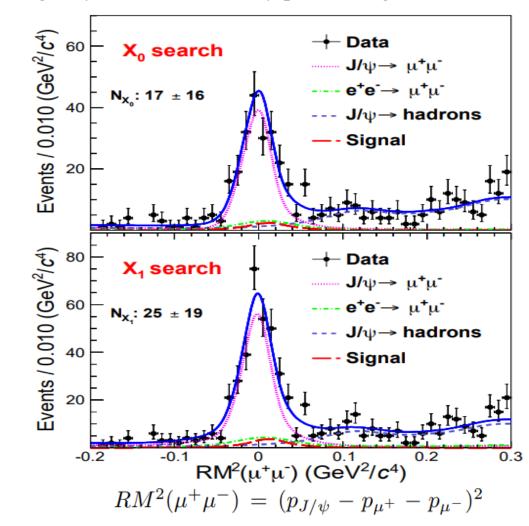
Belle II: PRL124 14, 141801

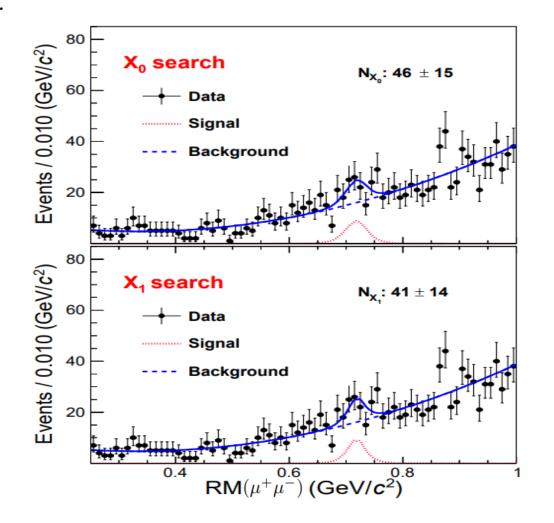
✓ BESIII can contribute to the low mass region

Search for a light muon philic scalar  $X_0$  or vector  $X_1$  is performed via  $J/\psi \to \mu^+\mu^- X_{0,1}$  with  $X_{0,1}$  invisible decays using  $(8.998 \pm 0.039) \times 10^9$  J/ $\psi$  events collected by the BESIII experiment.

# Search for an invisible muon philic scalar $X_0$ or vector $X_1$ via $J/\psi \to \mu^+\mu^-+$ invisible PRD 109, L031102 (2024)

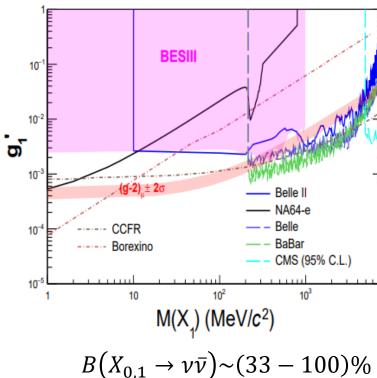
Signal yield is extracted by performing a series of ML fits.





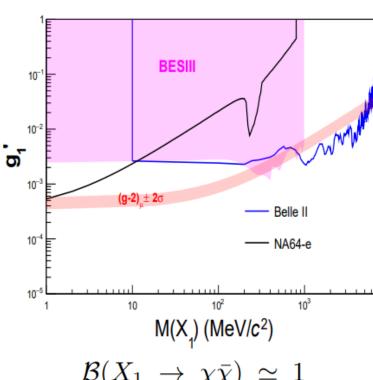
## Search for an invisible muon philic scalar $X_0$ or vector $X_1$ via $J/\psi \rightarrow \mu^{+}\mu^{-}$ +invisible PRD 109, L031102 (2024)

"vanilla"  $L_{\mu} - L_{\tau}$  model



$$B(X_{0.1} \to \nu \bar{\nu}) \sim (33 - 100)\%$$

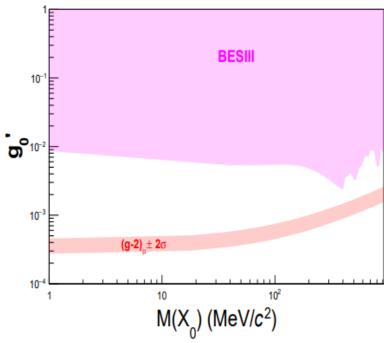
"invisible"  $L_{\mu} - L_{\tau}$  model



 $\mathcal{B}(X_1 \to \chi \bar{\chi}) \simeq 1$ 

JHEP10(2020)207

"scalar" U(1) model



 $X_0$  is long-lived with displaced decay or predominately decays to invisible particles

## **Summary**

- > Top priority is to search for new physics beyond the SM.
- ➤ BESIII plays a unique role to search for dark sector from e<sup>+</sup>e<sup>-</sup> collisions at the tau-charm region.
- ➤ A series of searches for invisible decays, dark photon, ALP and muon-philic particles have been performed at BESIII.
- > Only null results are available so far.
- ➤ BESIII limits exclude a large fraction of the parameter space of the new physics models beyond SM.
- $\triangleright$  More results is expected to come in the near future, especially with recently collected 20 fb<sup>-1</sup> of  $\psi(3770)$  data.

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# Thanks!