

# Neutrino Lorentz invariance violation from cosmic fields

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# CPT Violation and Lorentz Invariance

*CPT Symmetry: Relativistic quantum field theory (local).*

- *CPT* symmetry assumes Minkowski spacetime; hermiticity of the Hamiltonian  $\hat{H} = \hat{H}^\dagger$  and Lorentz invariance.
- *CPT*  $\rightarrow m_{e^-} = m_{e^+}$ , particle-antiparticle (mass) equality.
- Is *CPT* symmetry conserved at all energies  $E$ ?
- Quantum gravity ( $E_P \sim 10^{19}$  GeV), examples: string theory, loop quantum gravity, etc.
- String theory: Lorentz invariance violation (LIV) and *CPT* violation.
- LIV and *CPT* violation effects ( $\varepsilon$ ) suppressed by  $\varepsilon \lesssim 1/E_P$ .

V. A. Kostelecky and S. Samuel, Phys. Rev. D 39 (1989); V. A. Kostelecky and R. Potting, Nucl. Phys. B 359 (1991); V. A. Kostelecky and R. Potting, Phys. Lett. B 381 (1996). G. Barenboim, Front. in Phys. 10 (2022); C. A. Moura and F. Rossi-Torres, Universe 8 (2022) no.1, 42.

# Neutrino LIV from scalar fields

- Neutrino-scalar field interaction ( $\phi$ )

$$-\mathcal{L}_{\text{eff}} = \frac{g_{\alpha\beta}}{\Lambda} \partial_\mu \phi \bar{\nu}_{\alpha L} \gamma^\mu \nu_{\beta L}.$$

P. H. Gu, X. J. Bi and X. m. Zhang, Eur. Phys. J. C 50 (2007), 655-659; F. Simpson, R. Jimenez, C. Pena-Garay and L. Verde, Phys. Dark Univ. 20 (2018), 72-77; N. Klop and S. Ando, Phys. Rev. D 97 (2018) no.6; T. Gherghetta and A. Shkerin, Phys. Rev. D 108 (2023) no.9, 9; G. Lambiase and T. K. Poddar, JCAP 01 (2024), 069.

- Neutrino-pseudoscalar interaction ( $\phi$ )

$$-\mathcal{L}_{\text{eff}} = \frac{g_{\alpha\beta}}{\Lambda} \partial_\mu \phi \bar{\nu}_{\alpha L} \gamma^\mu \gamma_5 \nu_{\beta L}.$$

G. Y. Huang and N. Nath, Eur. Phys. J. C 78 (2018) no.11, 922.  
R. Cordero and L. A. Delgadillo, Phys. Rev. D 110 (2024) 11.  
<https://arxiv.org/abs/2407.02729>.

# Effective Interaction Model: Neutrino-Scalar Field

$$-\mathcal{L}_{\text{eff}} = \lambda_{\alpha\beta} \frac{\partial^\mu \phi}{\Lambda} \bar{\nu}_{\alpha L} \gamma_\mu \nu_{\beta L},$$

- $a_{\alpha\beta}^\mu = \lambda_{\alpha\beta} \partial^\mu \phi \Lambda^{-1}$ , couplings  $\lambda_{\alpha\beta} \sim \mathcal{O}(1)$ .
- $\phi$  as dark matter or dark energy (DE) candidate.
- Neutrino *CPT*-odd LIV induced by  $\phi$ .
- Klop, Ando:  $\phi$  is DE, signals in neutrino oscillations.
- Cordero, Delgadillo:  $\phi$  as dark matter (DM) candidate.

S. Ando, M. Kamionkowski and I. Mocioiu, Phys. Rev. D 80 (2009), 123522; N. Klop and S. Ando, Phys. Rev. D 97 (2018) no.6., R. Cordero, L. A. Delgadillo, PLB 853 (2024) 138687.

# Neutrino-Scalar Field Interaction: ( $\phi$ as DM)

Bosonic Current

$$\mathcal{L}_{\phi, Z'} \supset M_{Z'} \partial^\mu \phi Z'_\mu, \quad (1)$$

Neutrino Current

$$\mathcal{L}_{\nu, Z'} \supset g_{\alpha\beta} \bar{\nu}_{\alpha L} \gamma_\mu \nu_{\beta L} Z'^\mu. \quad (2)$$

- Integrating out  $Z'_\mu$  Eq. (1) and Eq. (2)

$$\mathcal{L}_{\phi, \nu} = g_{\alpha\beta} \frac{\partial^\mu \phi}{M_{Z'}} \bar{\nu}_{\alpha L} \gamma_\mu \nu_{\beta L},$$

- $\phi$  as DM,  $Z'_\mu$  new vector boson.

R. Cordero, L. A. Delgadillo, PLB 853 (2024) 138687

<https://arxiv.org/abs/2407.18513>

# Induced $CPT$ -odd LIV in Neutrinos

$$\mathcal{L}_{\phi,\nu}^{CPT} = \textcolor{blue}{a}_{\alpha\beta}^{\mu} \bar{\nu}_{\alpha L} \gamma_{\mu} \nu_{\beta L},$$

$CPT$ -odd LIV coefficients  $\textcolor{blue}{a}_{\alpha\beta}^{\mu}$

$$\textcolor{blue}{a}_{\alpha\beta}^{\mu} \rightarrow g_{\alpha\beta} \frac{\partial^{\mu} \phi}{M_{Z'}},$$

- $g_{\alpha\beta}$ : neutrino– $Z'_\mu$  couplings,  $M_{Z'} \sim [10 \text{ GeV}–\text{TeV}]$ .
- $\phi$  as ULDM:  $m_\phi \sim 10^{-15} \text{ eV}$  with  $V(\phi) \simeq m_\phi^2 \phi^2$ ,

$$\phi(x^\mu) \simeq \frac{\sqrt{2\rho_\phi}}{m_\phi} \sin((m_\phi(t - \langle v_\phi \rangle x)) \sim \phi(t).$$

E. G. M. Ferreira, Astron. Astrophys. Rev. 29 (2021) no.1, 7; A. Suárez, V. H. Robles and T. Matos, Astrophys. Space Sci. Proc. 38 (2014), 107-142. R. Cordero, L. A. Delgadillo, PLB 853 (2024) 138687

# Radiative Corrections

Lagrangian  $\mathcal{L}_{\phi,\nu}^{CPT}$  induces radiative corrections to  $m_\phi$

$$\delta m_\phi \simeq \frac{g_{\alpha\beta} m_\phi}{\pi M_{Z'}} \tilde{\Lambda} \lesssim m_\phi \left( g_{\alpha\beta} / \lesssim \pi \right) \left( \tilde{\Lambda} / M_{Z'} \right).$$

Neutrino mass corrections  $m_\nu$

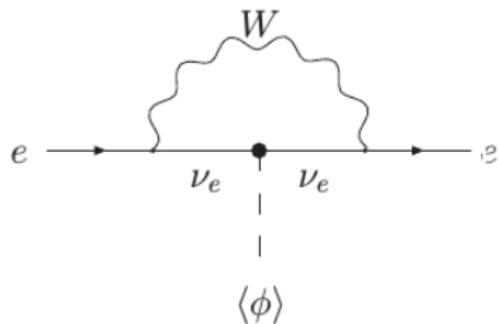
$$\delta m_\nu \simeq \frac{g_{\alpha\beta}^2 m_\nu}{\pi^2 M_{Z'}^2} \tilde{\Lambda}^2 \lesssim 10^{-3} m_\nu \left( g_{\alpha\beta} / \lesssim 0.1 \right) \left( \tilde{\Lambda} / M_{Z'} \right).$$

- Required ratio:  $g_{\alpha\beta} / M_{Z'} \lesssim 10^{-13} \text{ eV}^{-1}$ .
- Search for  $(Z'_\mu)$  at colliders and  $(a_{\alpha\beta}^\mu)$  in neutrino oscillation experiments.

W. Abdallah, A. K. Barik, S. K. Rai and T. Samui, Phys. Rev. D 104 (2021) no.9, 095031; M. G. Aartsen et al. [IceCube], Nature Phys. 14 (2018) no.9; K. Abe et al. [Super-Kamiokande], Phys. Rev. D. 91 (2015) no.5.

# Radiative Corrections

Lagrangian  $\mathcal{L}_{\phi,\nu}^{CPT}$  induces  $CPT$  violation in electrons  $e$



P. H. Gu, X. J. Bi and X. m. Zhang, Eur. Phys. J. C 50 (2007), 655-659.

$$\mathcal{L}_{\text{1-loop}} = \tilde{a}_e^\mu \bar{e} \gamma_\mu (1 - \gamma_5) e, \quad \tilde{a}_e^\mu \simeq g_{ee} \frac{\partial^\mu \phi}{M_{Z'}} \frac{\alpha M_{Z'}^2}{8\pi \sin^2 \theta_W M_W^2},$$

$\alpha$  fine structure constant,  $\theta_W$  Weinberg angle. Experimental bounds:  $\tilde{a}_e^T < 5 \times 10^{-25}$  GeV,  $\tilde{a}_e^X < 10^{-25}$  GeV.

$$\tilde{a}_e^T \lesssim 10^{-26} \text{ GeV} \quad (g_{ee} \lesssim 10^{-1}) (M_{Z'}/3M_W), \quad \tilde{a}_e^X \lesssim 10^{-3} \tilde{a}_e^T.$$

R. Cordero, L. A. Delgado, PLB 853 (2024) 138687.

# Neutrino-Scalar Field Interaction Potential

In early epochs,  $\mathcal{L}_{\phi,\nu}^{CPT}$  induces potential

$$V(\dot{\phi}) = \frac{g_{\alpha\beta} n_\nu}{M_{Z'}} \dot{\phi}.$$

- $n_\nu$  neutrino number density,  $\phi(x^\mu) \simeq \phi(t)$ .

$\phi$  evolution with interaction ( $Q$ )

$$\ddot{\phi} + 3H\dot{\phi} + \frac{dV(\phi)}{d\phi} = -\frac{Q}{\dot{\phi}},$$

$$Q = \tilde{g}_{\alpha\beta} n_\nu \ddot{\phi} \quad (\text{Simpson et al.})$$

$$\ddot{\phi} \left( 1 + \frac{g_{\alpha\beta} n_\nu}{M_{Z'} \dot{\phi}} \right) + 3H\dot{\phi} + \frac{dV(\phi, \dot{\phi})}{d\phi} = 0.$$

F. Simpson, R. Jimenez, C. Pena-Garay and L. Verde, Phys. Dark Univ. 20 (2018), 72-77.

# Scalar Field Evolution: Energy Exchange ( $Q$ )

- $Q = \tilde{g}_{\alpha\beta} n_\nu \ddot{\phi}$  (Simpson et al.),  $\phi$  behaves as DE,  $\phi = \text{constant}$ .
- Proposal:  $Q = m_\nu \tilde{g}_{\alpha\beta} n_\nu \dot{\phi}$  ( $Q \propto \dot{\phi}$ , Boehmer et al.).
- Is  $\phi$  a ULDM candidate?

$$\ddot{\phi} + 3H\dot{\phi} + \frac{dV(\phi, \dot{\phi})}{d\phi} + \frac{g_{\alpha\beta} n_\nu m_\nu}{M_{Z'}} = 0.$$

- $\phi$  as ULDM, but oscillations delayed if:

$$\frac{g_{\alpha\beta} n_\nu m_\nu}{M_{Z'}} \lesssim \left| \frac{dV(\phi, \dot{\phi})}{d\phi} \right| \sim m_\phi^2 |\dot{\phi}|.$$

R. Cordero, L. A. Delgadillo, PLB 853 (2024) 138687.

C. G. Boehmer, G. Caldera-Cabral, R. Lazkoz and R. Maartens, Phys. Rev. D 78 (2008), 023505.

# Lorentz Invariance Violation (LIV) in Neutrinos

LIV/CPT Parameterization (effective theory  $E \ll E_P$ ):

$$\mathcal{L}_{\text{eff}}^{\psi} = \mathcal{L}_{\text{SM}}^{\psi} + \mathcal{L}_{\text{LIV}} + \text{h.c.},$$

$$\begin{aligned}\mathcal{L}_{\text{LIV}} = -\frac{1}{2} \Big\{ & p_{\alpha\beta}^{\mu} \bar{\psi}_{\alpha} \gamma_{\mu} \psi_{\beta} + q_{\alpha\beta}^{\mu} \bar{\psi}_{\alpha} \gamma_{\mu} \gamma_5 \psi_{\beta} - \\ & ir_{\alpha\beta}^{\mu\nu} \bar{\psi}_{\alpha} \gamma_{\mu} \partial_{\nu} \psi_{\beta} - is_{\alpha\beta}^{\mu\nu} \bar{\psi}_{\alpha} \gamma_{\mu} \gamma_5 \partial_{\nu} \psi_{\beta} \Big\},\end{aligned}$$

For neutrinos, we define:

$$(a_L)_{\alpha\beta}^{\mu} = (p + q)_{\alpha\beta}^{\mu} \text{ [Energy]}, \quad (c_L)_{\alpha\beta}^{\mu\nu} = (r + s)_{\alpha\beta}^{\mu\nu} [E^0].$$

V.A. Kostelecky and M. Mewes, Phys. Rev. D 69 (2004). G. Barenboim, Front. in Phys. 10 (2022); C. A. Moura and F. Rossi-Torres, Universe 8 (2022) no.1, 42.

# *CPT*-even SME $C_{\alpha\beta}^{\mu\nu}$ in neutrino oscillations

- Effective Hamiltonian (Neutrino Oscillations)

$$H_{\text{eff}} = H_{\text{vacuum}} + H_{\text{matter}} + H_{\text{LIV}},$$

- LIV Hamiltonian, *CPT*-even (*T*-time and *Z*-spatial) SME coefficients:

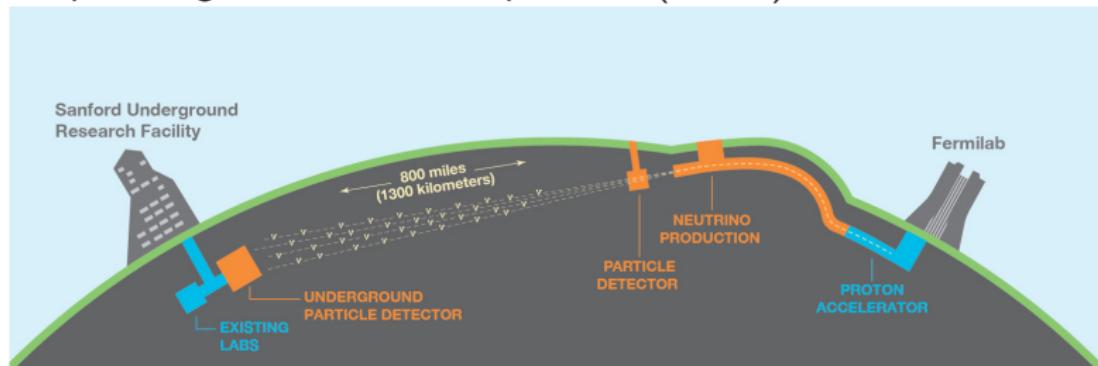
$$H_{\text{LIV}} = -\frac{E_\nu}{2} \left[ (3 - \hat{N}_Z^2)(c_L)^{TT}_{\alpha\beta} + (3\hat{N}_Z^2 - 1)(c_L)^{ZZ}_{\alpha\beta} - 2\hat{N}_Z(c_L)^{TZ}_{\alpha\beta} \right].$$

Neutrinos  $E_\nu \gtrsim$  GeV: Higher sensitivity to LIV and *CPT* effects!

V.A. Kostelecky and M. Mewes, Phys. Rev. D 69 (2004). G. Barenboim, Front. in Phys. 10 (2022); C. A. Moura and F. Rossi-Torres, Universe 8 (2022) no.1, 42.

# Tensorial neutrino interaction with cosmic fields: Sensitivity to $CPT$ -even SME ( $C_{\alpha\beta}^{\mu\nu}$ ) in DUNE

Deep Underground Neutrino Experiment (DUNE)



$$H_{\text{LIV}} = -\frac{E_\nu}{2} \left[ (3 - \hat{N}_Z^2)(c_L)^{TT}_{\alpha\beta} + (3\hat{N}_Z^2 - 1)(c_L)^{ZZ}_{\alpha\beta} - 2\hat{N}_Z(c_L)^{TZ}_{\alpha\beta} \right].$$

DUNE:  $E_\nu \sim 3 \text{ GeV}$ ,  $\hat{N}^Z \simeq 0.16$ .

J. S. Diaz, V. A. Kostelecky and M. Mewes, Phys. Rev. D 80, 076007 (2009) R. Cordero,  
L. A. Delgado, O. G. Miranda, C. Moura, EPJC (2025) 85:6

# Tensorial neutrino interaction with cosmic fields

$$-\mathcal{L}_{\text{eff}} = -i \frac{\lambda_{\alpha\beta}}{M_*^4} T_\varphi^{\mu\nu} \bar{\nu}_{\alpha L} \gamma_\mu \partial_\nu \nu_{\beta L}$$

- $CPT$ -even LIV coefficients of the SME  $(c_L)_{\alpha\beta}^{\mu\nu} = c_{\alpha\beta}^{\mu\nu}$ :

$$c_{\alpha\beta}^{\mu\nu} \rightarrow \frac{\lambda_{\alpha\beta}}{M_*^4} T_\varphi^{\mu\nu}.$$

- $T_\varphi^{\mu\nu}(\varphi) \rightarrow \varphi$  as either dark matter or dark energy candidate.

- Isotropic LIV coefficients  $(c_L)_{\alpha\beta}^{00} = c_{\alpha\beta}^{00}$ :

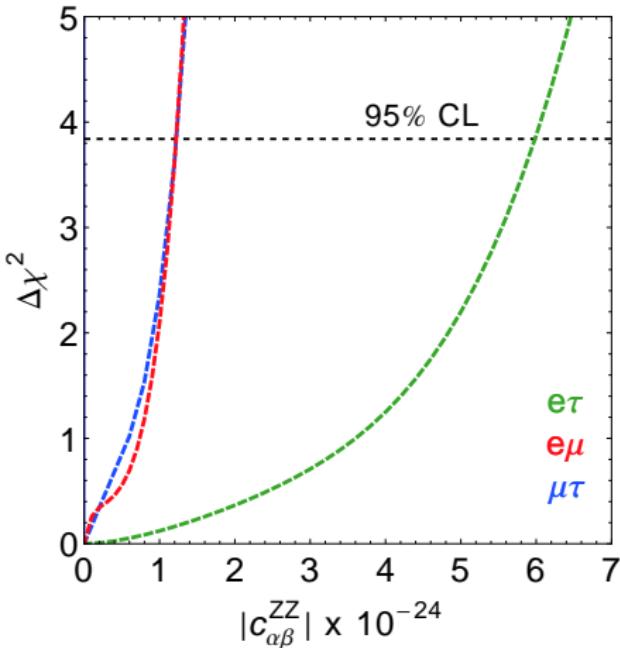
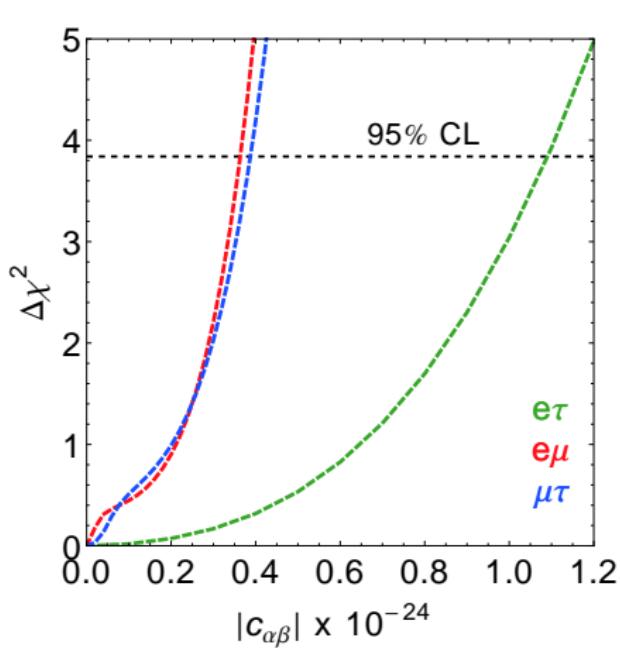
$$c_{\alpha\beta}^{00} \rightarrow \frac{\lambda_{\alpha\beta}}{M_*^4} T_\varphi^{00} = \frac{\lambda_{\alpha\beta}}{M_*^4} \rho_\varphi.$$

- $Z$ -spatial LIV coefficients  $(c_L)_{\alpha\beta}^{33} = c_{\alpha\beta}^{ZZ}$ :

$$c_{\alpha\beta}^{ZZ} \rightarrow \frac{\lambda_{\alpha\beta}}{M_*^4} (T_\varphi^{ZZ} + \delta T_\varphi^{ZZ}). \quad (3)$$

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# Future Sensitivity in DUNE ( $C_{\alpha\beta}^T = C_{\alpha\beta}$ , $C_{\alpha\beta}^{ZZ}$ )



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<https://arxiv.org/abs/2407.18513>

# Future Sensitivities at DUNE

Table: DUNE projected sensitivities, last column shows the interaction energy scale  $M_*$  associated to the SME coefficients  $c_{\alpha\beta}^{\mu\nu}$ .

Sensitivity	Interaction energy scale
$ c_{\alpha\beta}  \sim [1 - 10] \times 10^{-25}$	$M_* \sim [3 - 6] \times 10^4 \text{ eV } (\rho_{\varphi, \odot}^{\text{DM}})$
$ c_{\alpha\beta}^{ZZ}  \sim [1 - 10] \times 10^{-24}$	$M_* \sim [2 - 3] \times 10^3 \text{ eV } (\rho_{\varphi}^{\text{DE}})$
$c_{ee} - c_{\tau\tau} \simeq 1.3 \times 10^{-24}$	$M_* \sim 3 \times 10^4 \text{ eV } (\rho_{\varphi, \odot}^{\text{DM}})$
$c_{\mu\mu} - c_{\tau\tau} \simeq 2.2 \times 10^{-24}$	$M_* \sim 3 \times 10^4 \text{ eV } (\rho_{\varphi, \odot}^{\text{DM}})$
$c_{ee}^{ZZ} - c_{\tau\tau}^{ZZ} \simeq 5.0 \times 10^{-24}$	$M_* \sim 2 \times 10^3 \text{ eV } (\rho_{\varphi}^{\text{DE}})$
$c_{\mu\mu}^{ZZ} - c_{\tau\tau}^{ZZ} \simeq 7.0 \times 10^{-24}$	$M_* \sim 2 \times 10^3 \text{ eV } (\rho_{\varphi}^{\text{DE}})$

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# Summary

- Neutrino–ultralight DM interaction ( $m_\phi \sim 10^{-15}$  eV):

$$\mathcal{L}_{\phi,\nu} = g_{\alpha\beta} \frac{\partial^\mu \phi}{M_{Z'}} \bar{\nu}_{\alpha L} \gamma_\mu \nu_{\beta L}.$$

- Signatures via  $Z'_\mu$  searches and neutrino oscillations:

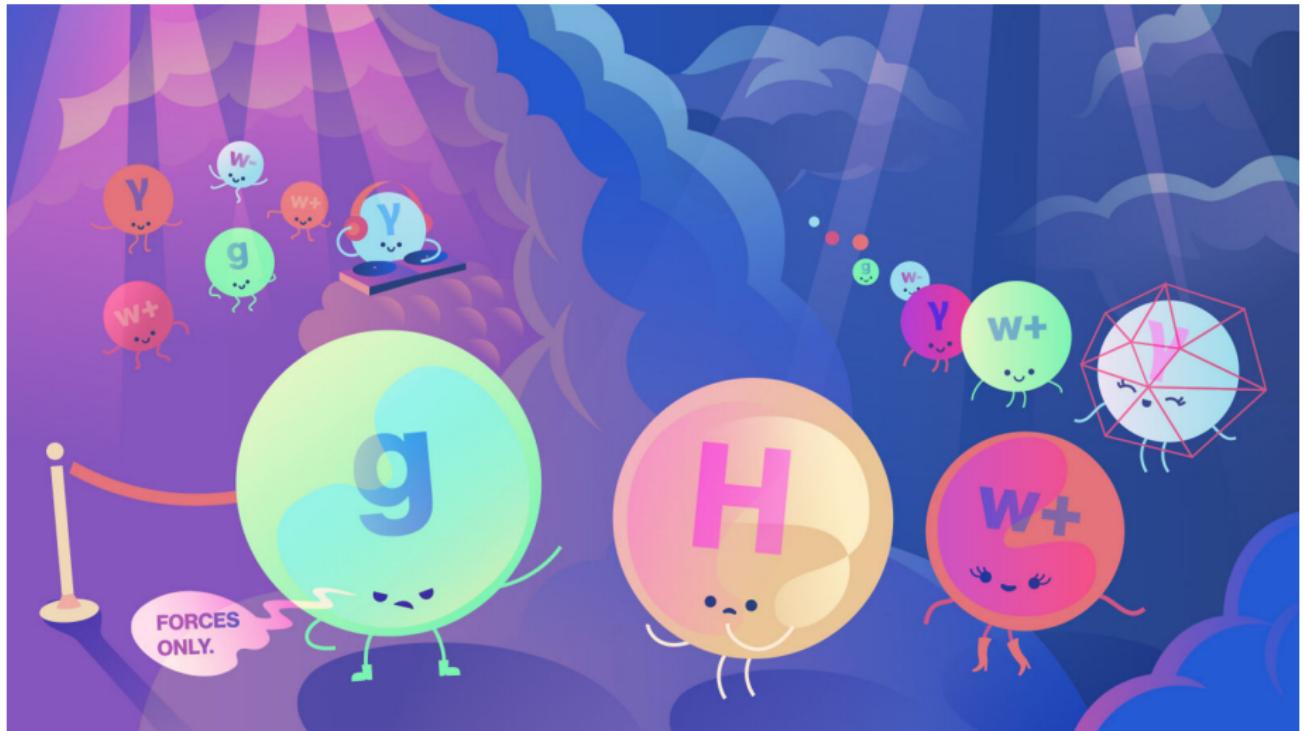
$$a_{\alpha\beta}^\mu \rightarrow g_{\alpha\beta} \frac{\partial^\mu \phi}{M_{Z'}}, \quad a_{\alpha\beta}^T \lesssim 10^{-25} \text{ GeV}, \quad a_{\alpha\beta}^X \lesssim 10^{-28} \text{ GeV}.$$

- Depending on  $\nu - \phi$  energy exchange ( $Q$ ), scalar field behaves as DE or DM.
- Tensorial–neutrino interaction:  $T_\varphi^{\mu\nu}(\varphi) \rightarrow \varphi$  as either DM or DE candidate.
- *CPT*-even SME coefficients related to the energy-momentum tensor:

$$c_{\alpha\beta}^{\mu\nu} \rightarrow \frac{\lambda_{\alpha\beta}}{M_*^4} T_\varphi^{\mu\nu}.$$

- Competitive sensitivities to the *CPT*-even SME coefficients at DUNE.

**Thank you for your attention!**



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# BACK UP: Scalar Field Evolution ( $T \gtrsim 1$ eV, $n_\nu \sim 0.1 T^3$ )

Equation of Motion:  $\phi(t) \rightarrow \phi(T)$

$$H(t) = \frac{1}{2t} \simeq 1.66 \sqrt{g_*} \frac{T^2}{M_{\text{Pl}}},$$

$g_*$ : relativistic degrees of freedom.

$$\begin{aligned}\phi(T) &\rightarrow \left( \frac{d\phi}{dT} - \frac{0.1 g_{\alpha\beta} M_{\text{Pl}}}{(1.66) \sqrt{g_*} M_{Z'}} \right) \frac{d^2\phi}{dT^2} = \\ &\left( \frac{(0.3) g_{\alpha\beta} M_{\text{Pl}}}{M_{Z'} (1.66) \sqrt{g_*} T} - \frac{m_\phi^2 M_{\text{Pl}}^2 \phi}{(1.66)^2 g_* T^6} \right) \frac{d\phi}{dT}.\end{aligned}$$

Thus:  $\phi = \text{constant}$  is a solution!

# LIV and CPT Violation in Neutrino sector

CPT-odd LIV coefficients:  $CPT(a_{\alpha\beta}^{\mu}) = (-)a_{\alpha\beta}^{\mu}$ ,

$$CPT(\bar{\psi}_{\alpha}\gamma_{\mu}\psi_{\beta}) = (-)\bar{\psi}_{\alpha}\gamma_{\mu}\psi_{\beta}, \quad CPT(\bar{\psi}_{\alpha}\gamma_{\mu}\gamma_5\psi_{\beta}) = (-)\bar{\psi}_{\alpha}\gamma_{\mu}\gamma_5\psi_{\beta},$$

$$CPT(\partial_{\nu}) = (-)\partial_{\nu}, \quad CPT(\bar{\psi}_{\alpha}\partial^{\mu}\gamma_{\mu}\psi_{\beta}) = (-)^2\bar{\psi}_{\alpha}\partial^{\mu}\gamma_{\mu}\psi_{\beta}.$$

CPT-even LIV coefficients:  $CPT(c_{\alpha\beta}^{\mu\nu}) = (-)^2c_{\alpha\beta}^{\mu\nu}$ ,

$$CPT(\bar{\psi}_{\alpha}\gamma_{\mu}\partial_{\nu}\psi_{\beta}) = (-)^2\bar{\psi}_{\alpha}\gamma_{\mu}\partial_{\nu}\psi_{\beta},$$

$$CPT(\bar{\psi}_{\alpha}\gamma_{\mu}\gamma_5\partial_{\nu}\psi_{\beta}) = (-)^2\bar{\psi}_{\alpha}\gamma_{\mu}\gamma_5\partial_{\nu}\psi_{\beta}.$$

G. Barenboim, Front. in Phys. 10 (2022).