

# The low-scale seesaw solution to the $M_W$ and $(g - 2)_\mu$ anomalies

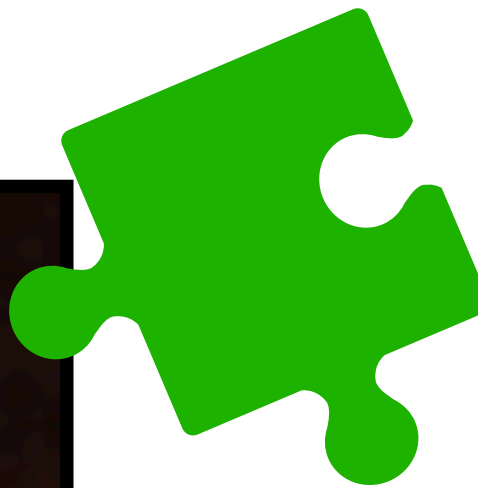
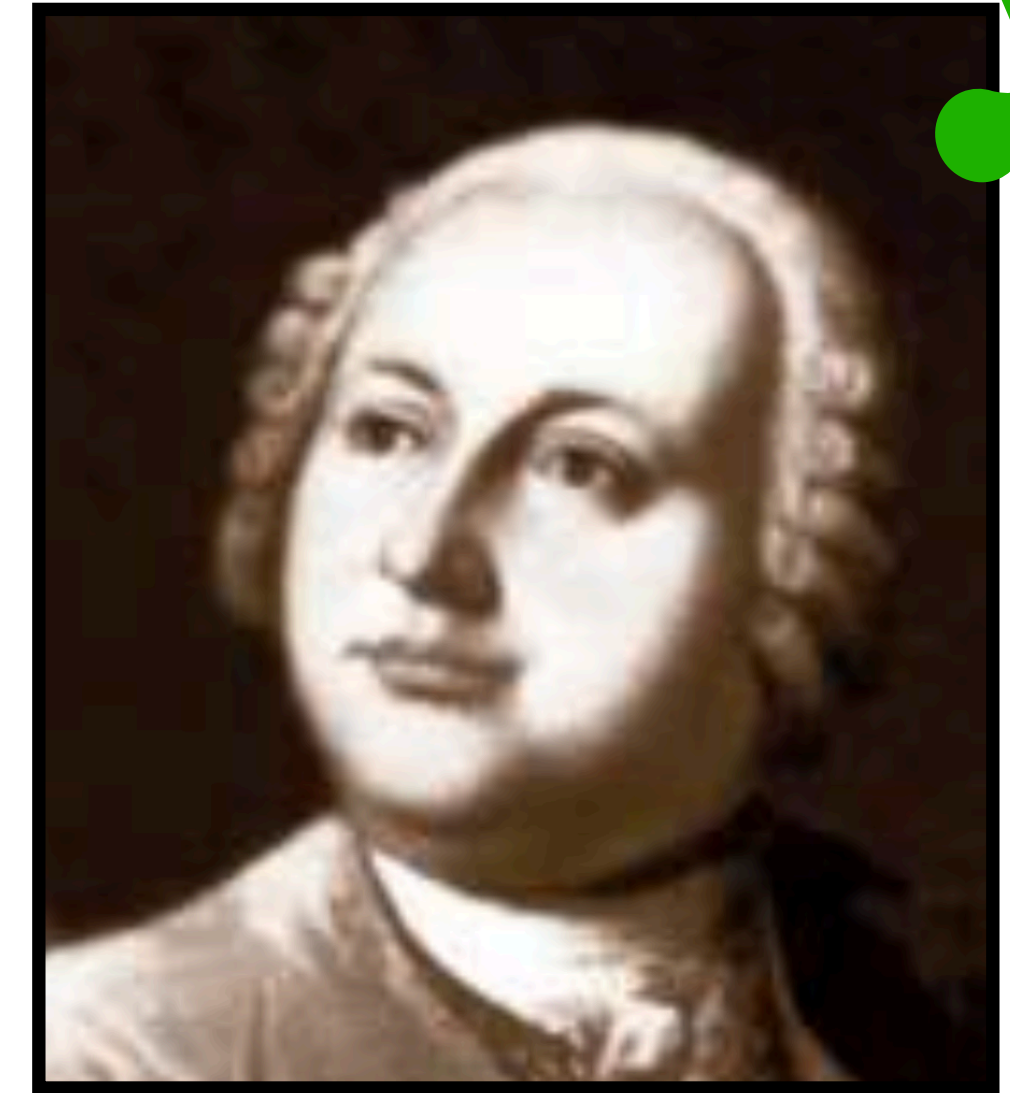
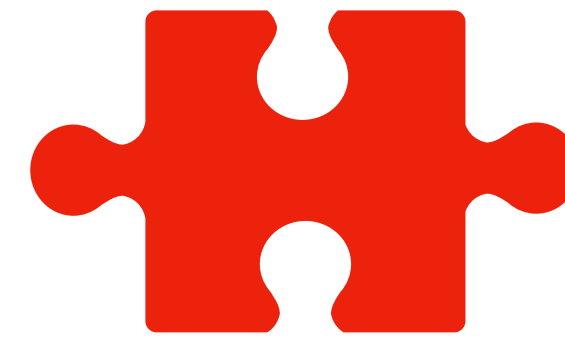
Based on:

arXiv: **2211.03797**, Fortsch.Phys. 71 (2023)

realised in collaboration with:

L. Merlo (UAM/IFT) and Stefan Pokorski (University of Warsaw)

**Arturo de Giorgi (Madrid, UAM/IFT)**

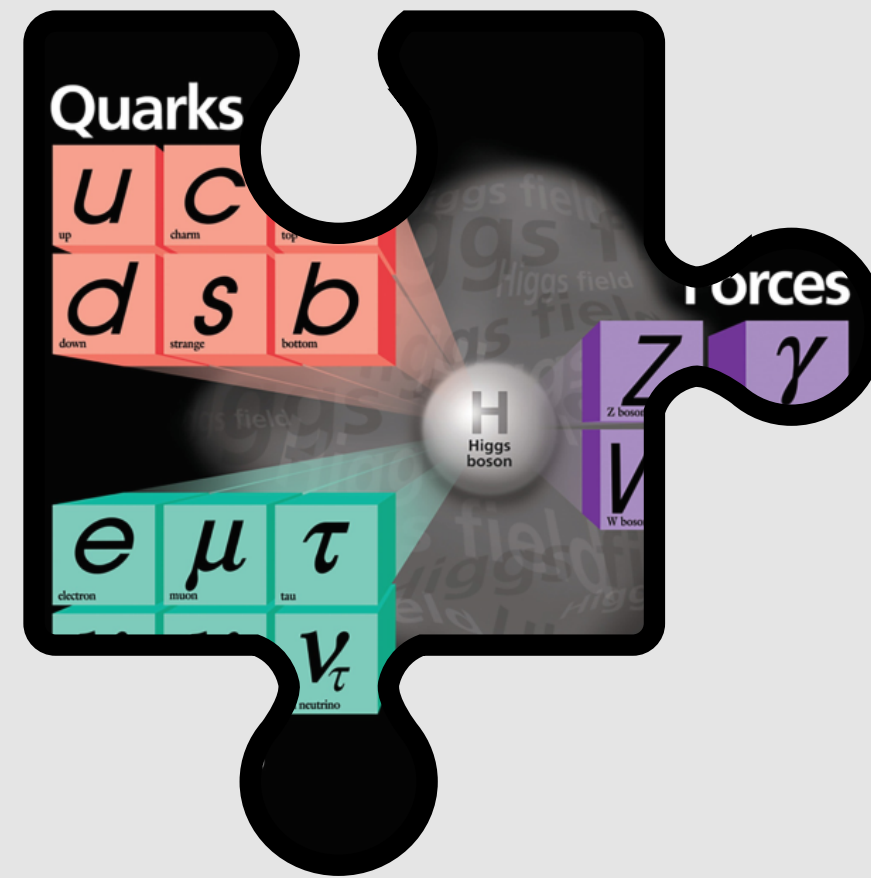


**21st Lomonosov Conference**

August 25, 2023

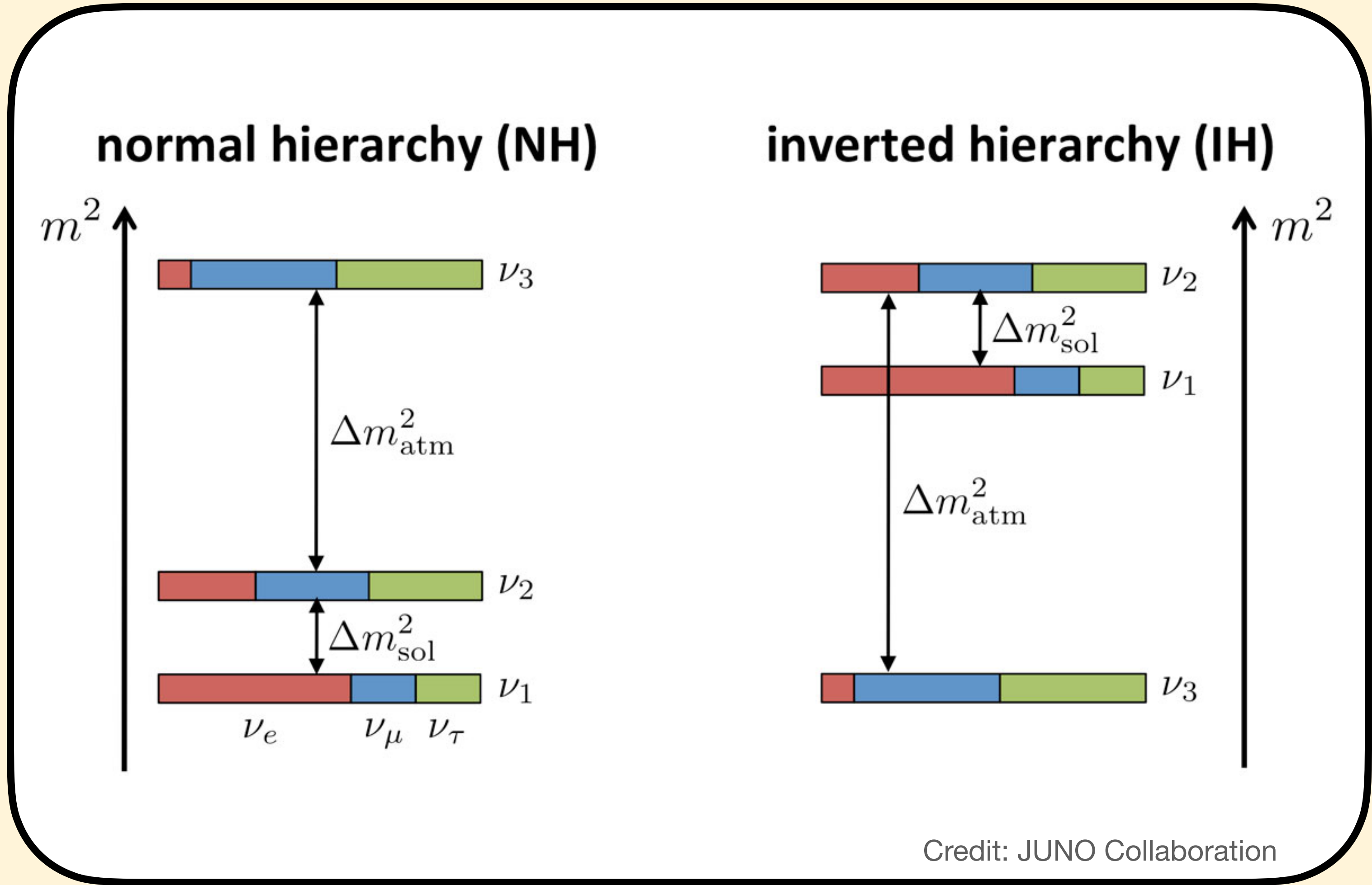


# (Some) SM's Problems/Puzzles

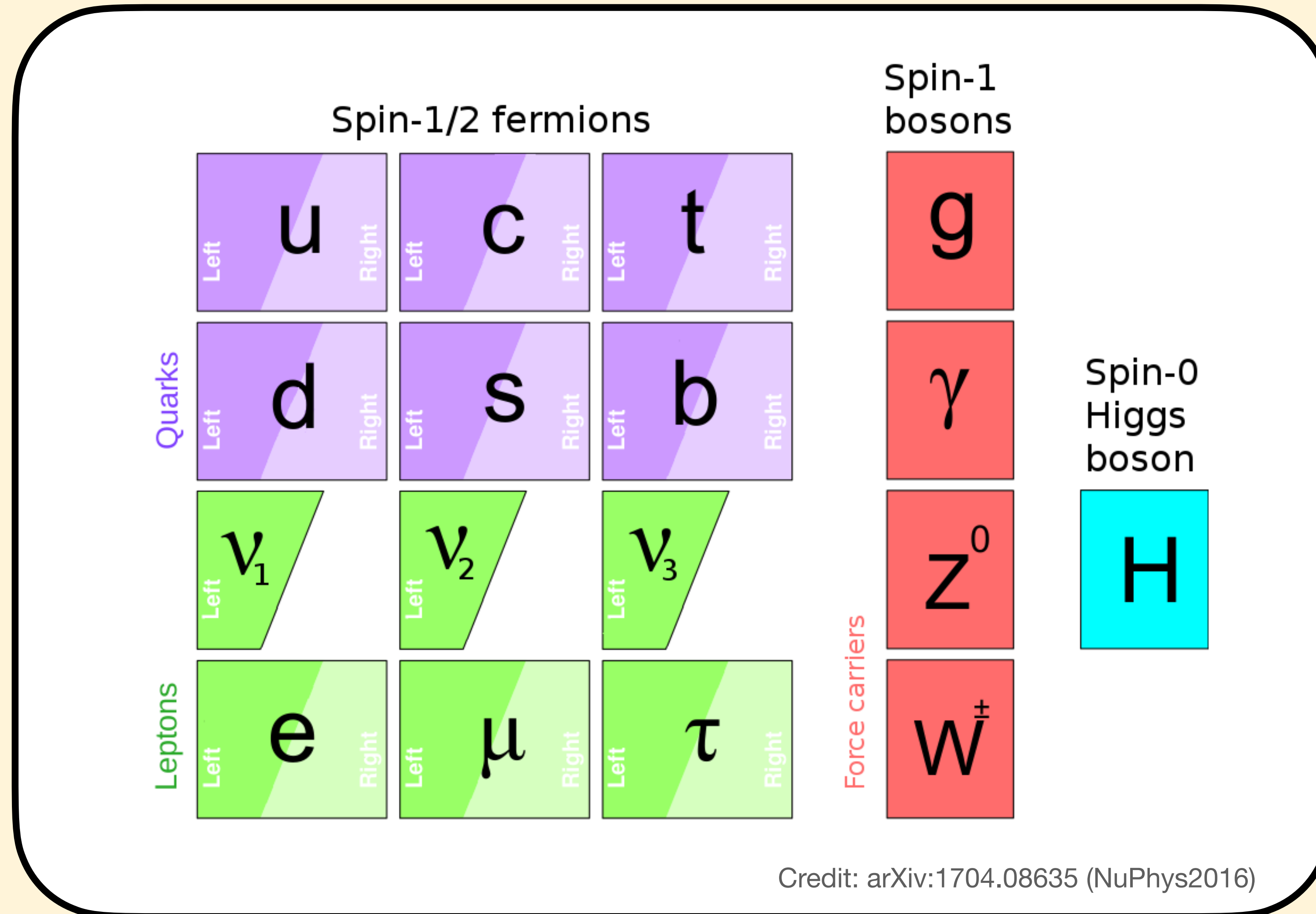


# $\nu$ -Masses

	NO	IO
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.41^{+0.21}_{-0.20}$	$7.41^{+0.21}_{-0.20}$
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.507^{+0.026}_{-0.027}$	$-2.486^{+0.025}_{-0.028}$
<div style="border: 1px solid black; padding: 2px; display: inline-block;"> <math>\nu</math>Fit Collaboration [2007.14792]         </div>		



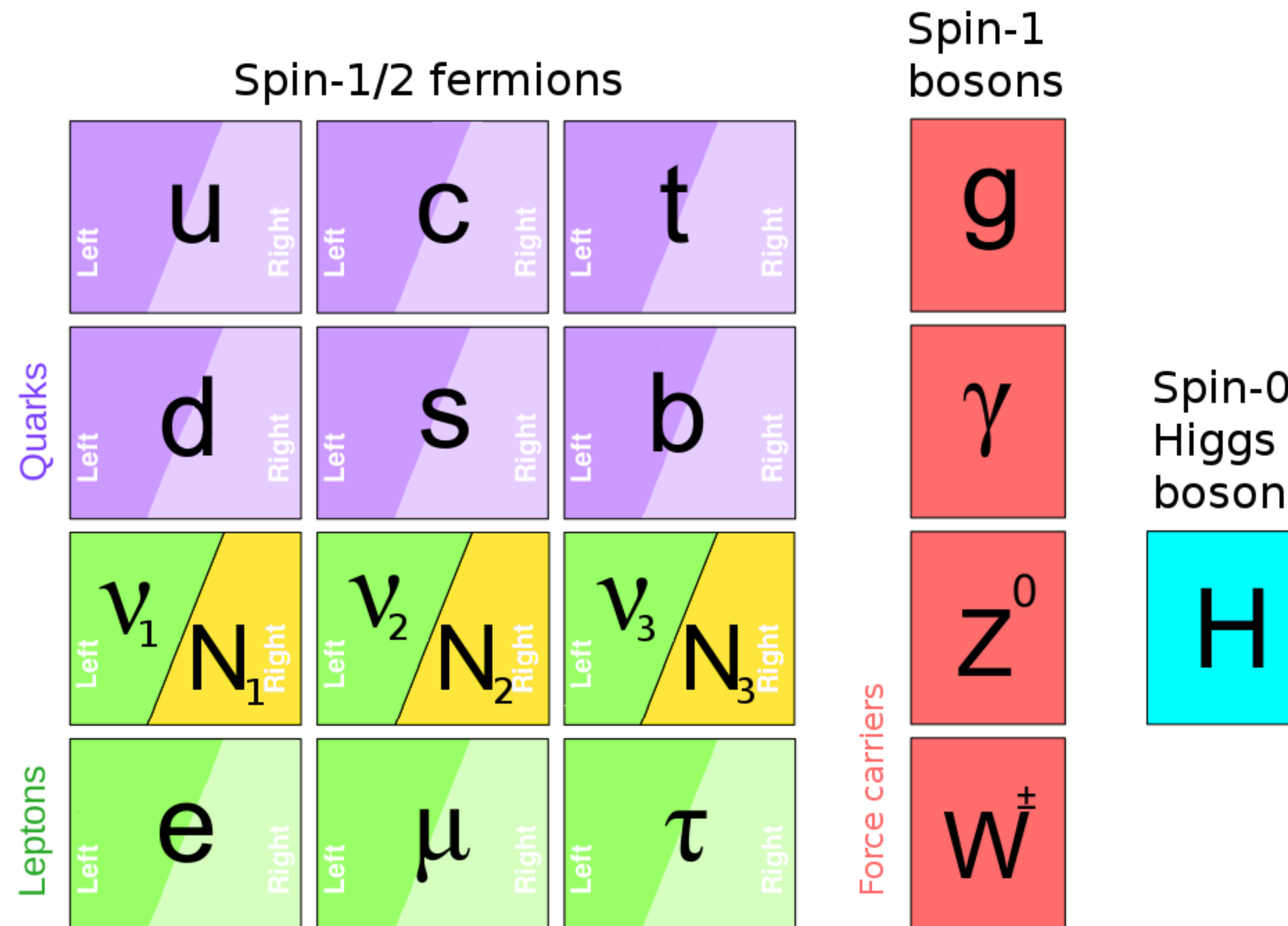
# Sterile Neutrinos?



Credit: arXiv:1704.08635 (NuPhys2016)



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# Seesaw Mechanism?

In the **SMEFT**, only 1-operator at dimension 5 that can describe  $\nu$ -masses: “**Weinberg-Operator**”

$$\delta\mathcal{L}^{d=5} = \left(\frac{c_{\alpha\beta}}{\Lambda}\right) \left(\overline{\ell}_{L\alpha}^c \tilde{H}^*\right) \left(\tilde{H}^\dagger \ell_{L\beta}\right)$$

Steven Weinberg, PRL **43**, 1566 (1979)

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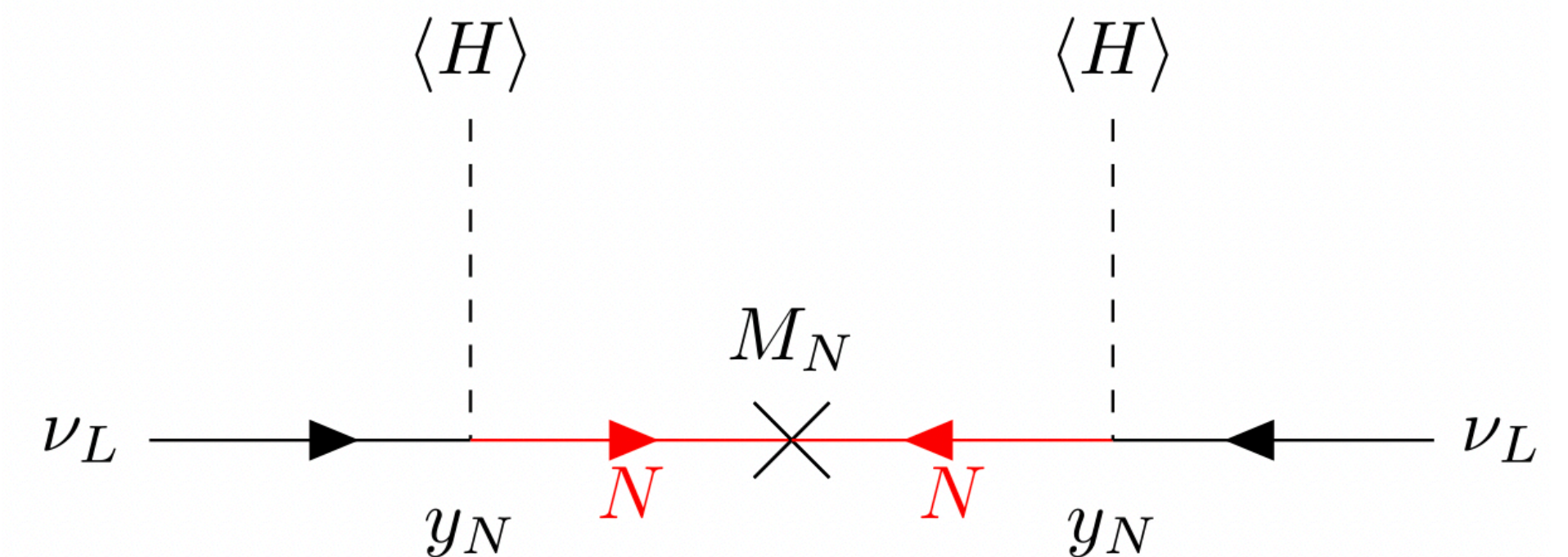
Could we use a very **large scale** to justify the **smallness** of  $\nu$ -masses?

$$\mathcal{L} = \mathcal{L}_{SM} + i\bar{N}_\alpha \not{\partial} N_\alpha - \left[ \lambda_N^{\alpha b} \bar{N}^\alpha \tilde{H}^\dagger \ell_L^b + \frac{M_{\alpha\beta}}{2} \bar{N}^\alpha N^{\beta c} + h.c. \right]$$

$$M_\nu = \begin{pmatrix} 0 & \lambda_N^T v / \sqrt{2} \\ \lambda_N v / \sqrt{2} & M \end{pmatrix} \quad m_\nu = \frac{\lambda_N^2 \langle H \rangle^2}{2M_N}$$

$$\lambda_N \sim \mathcal{O}(1) \longrightarrow M_N \sim 10^{14} \text{ GeV}$$

Diagrammatically:



P. Minkowski, P.L. B 67 (1977).  
 M. Gell-Mann, P. Ramond, and R. Slansky, Conf. Proc. C 790927 (1979).  
 T. Yanagida, C.P.C 7902131 (1979).  
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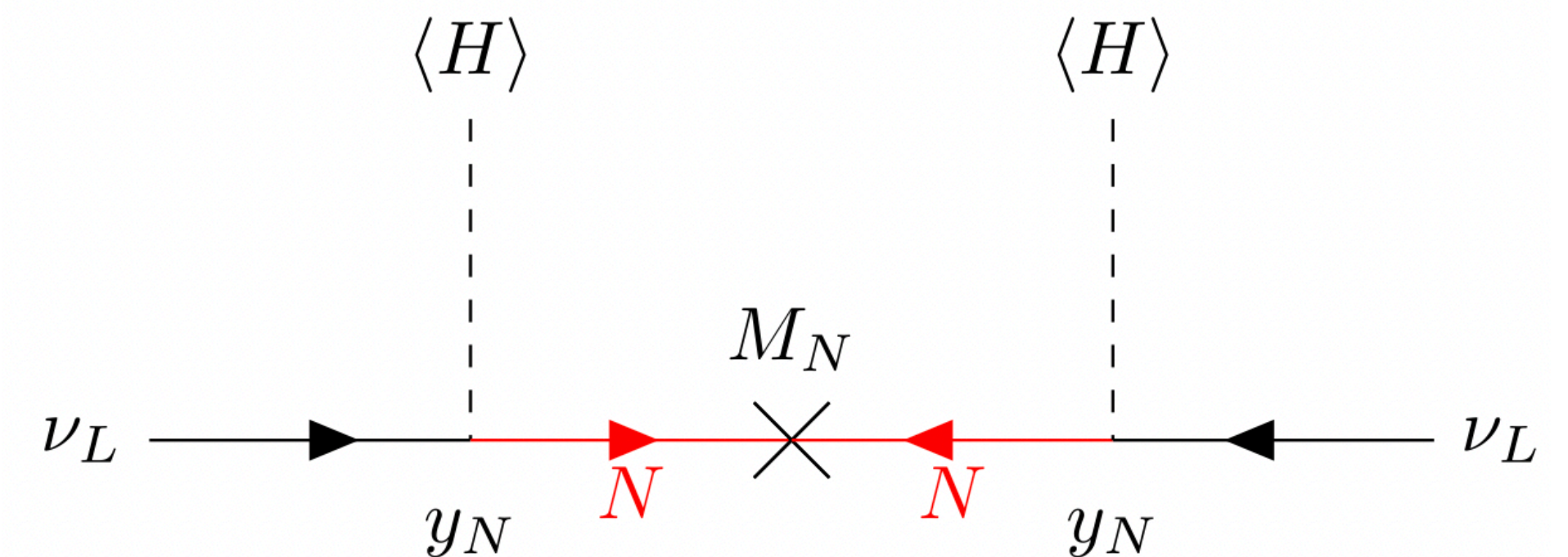
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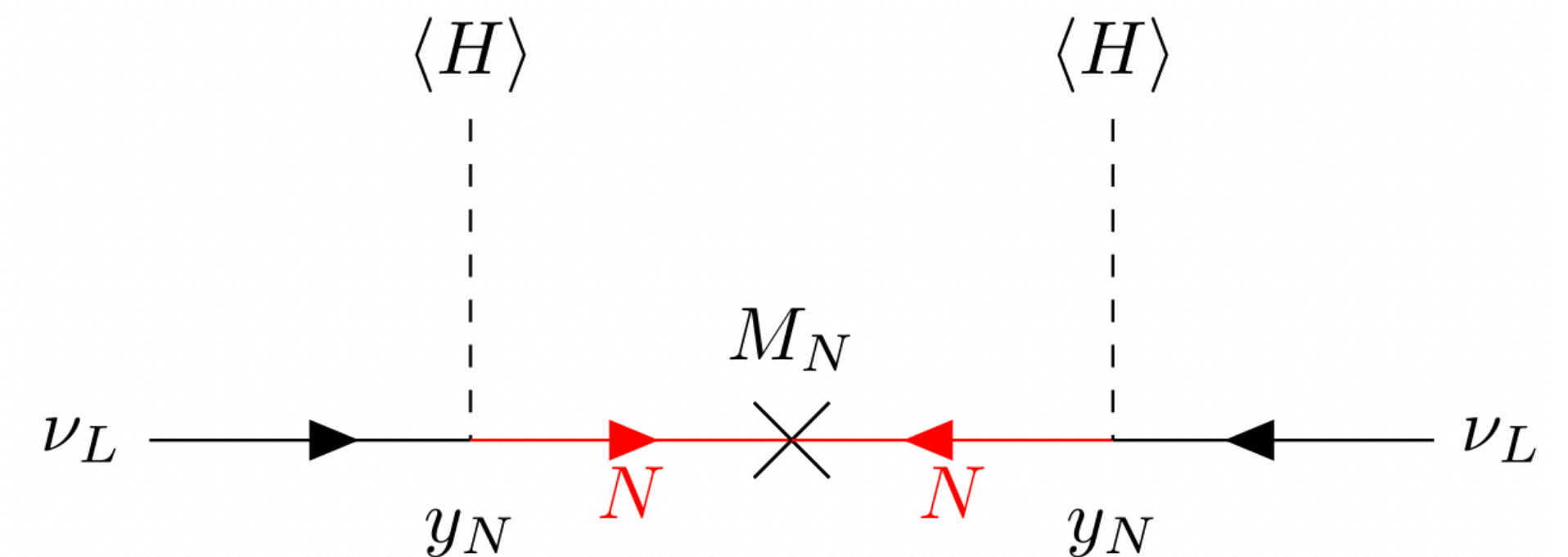
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HUGE! Something more “testable”?

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 M. Gell-Mann, P. Ramond, and R. Slansky, Conf. Proc. C 790927 (1979).  
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# Low-scale Seesaws: a global symmetry?

$$-\mathcal{L}_Y = \bar{\ell}_L H Y_\mu \mu_R + \bar{\ell}_L \tilde{H} Y_N N_R + \epsilon \bar{\ell}_L \tilde{H} Y_S S_R + \frac{1}{2} \mu' \overline{N_R^c} N_R + \frac{1}{2} \mu \overline{S_R^c} S_R + \Lambda \overline{N_R^c} S_R$$

Small Symmetry Breaking Parameters

$$\chi \equiv (\nu_L, N_R^c, S_R^c)^T$$

$$-\mathcal{L}_Y \supset \frac{1}{2} \bar{\chi} \mathcal{M}_\chi \chi^c + \text{h.c.}$$

$$\mathcal{M}_\chi = \begin{pmatrix} 0 & \frac{v}{\sqrt{2}} Y_N & \epsilon \frac{v}{\sqrt{2}} Y_S \\ \frac{v}{\sqrt{2}} Y_N^T & \mu' & \Lambda \\ \epsilon \frac{v}{\sqrt{2}} Y_S^T & \Lambda^T & \mu \end{pmatrix}$$

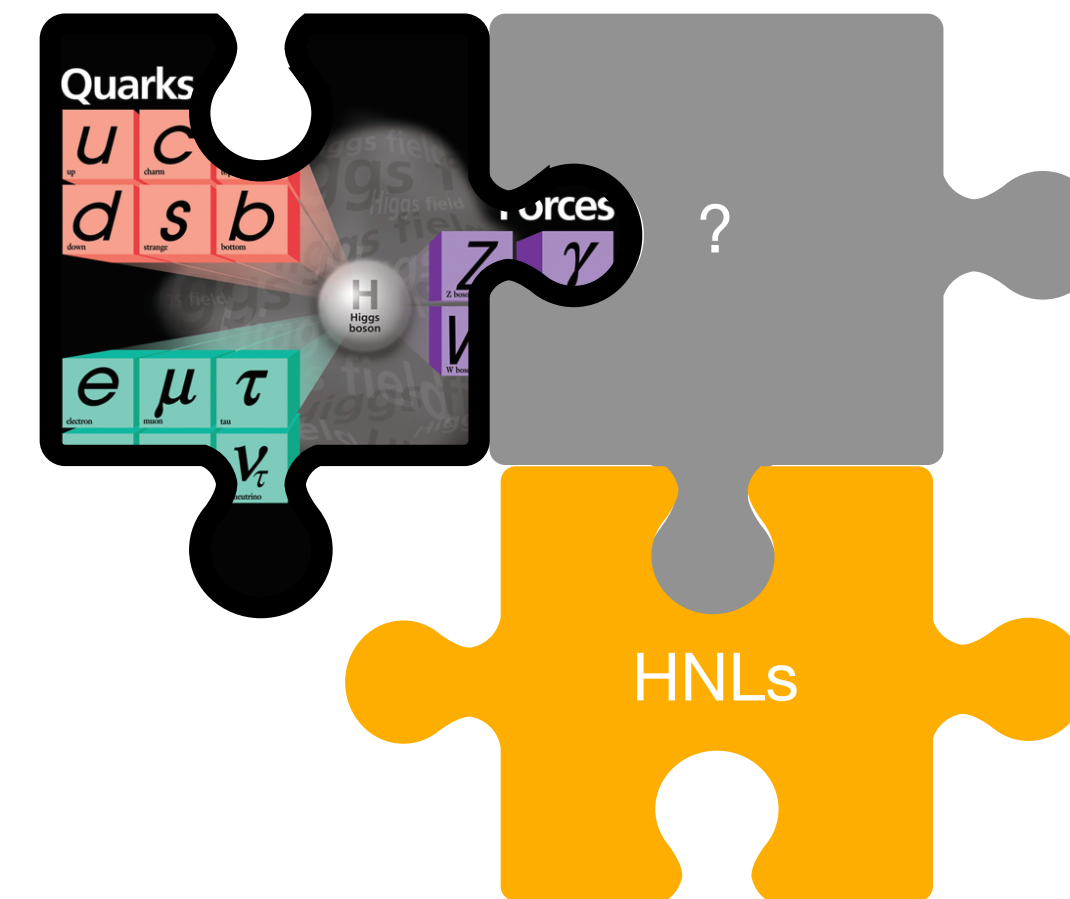
$$m_\nu \simeq \frac{v^2}{2} \left[ \left( Y_N \frac{1}{\Lambda^T} \mu \frac{1}{\Lambda} Y_N^T \right) - \epsilon \left( Y_S \frac{1}{\Lambda} Y_N^T + Y_N \frac{1}{\Lambda^T} Y_S^T \right) \right]$$

D. Wyler and L. Wolfenstein, Nucl. Phys. B 218 (1983) 205–214.  
 R. N. Mohapatra and J. W. F. Valle, Phys. Rev. D 34 (1986) 1642.  
 J. Bernabeu et al., Phys. Lett. B 187 (1987) 303–308.  
 M. Malinsky et al., Phys. Rev. Lett. 95 (2005) 161801.



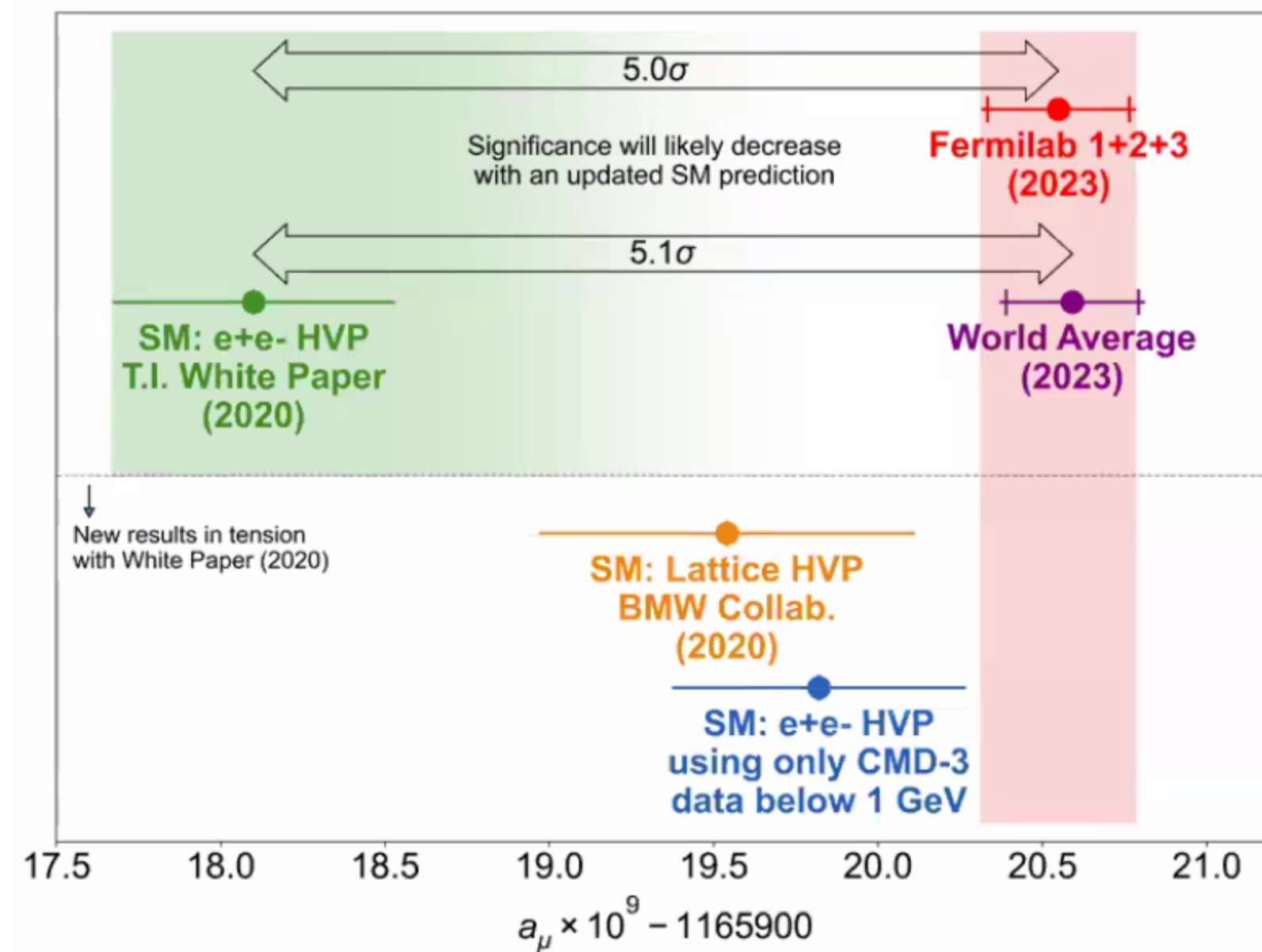
HNLs cannot account for **all** problems/puzzles of the SM

What extra **“piece”** could make a difference?



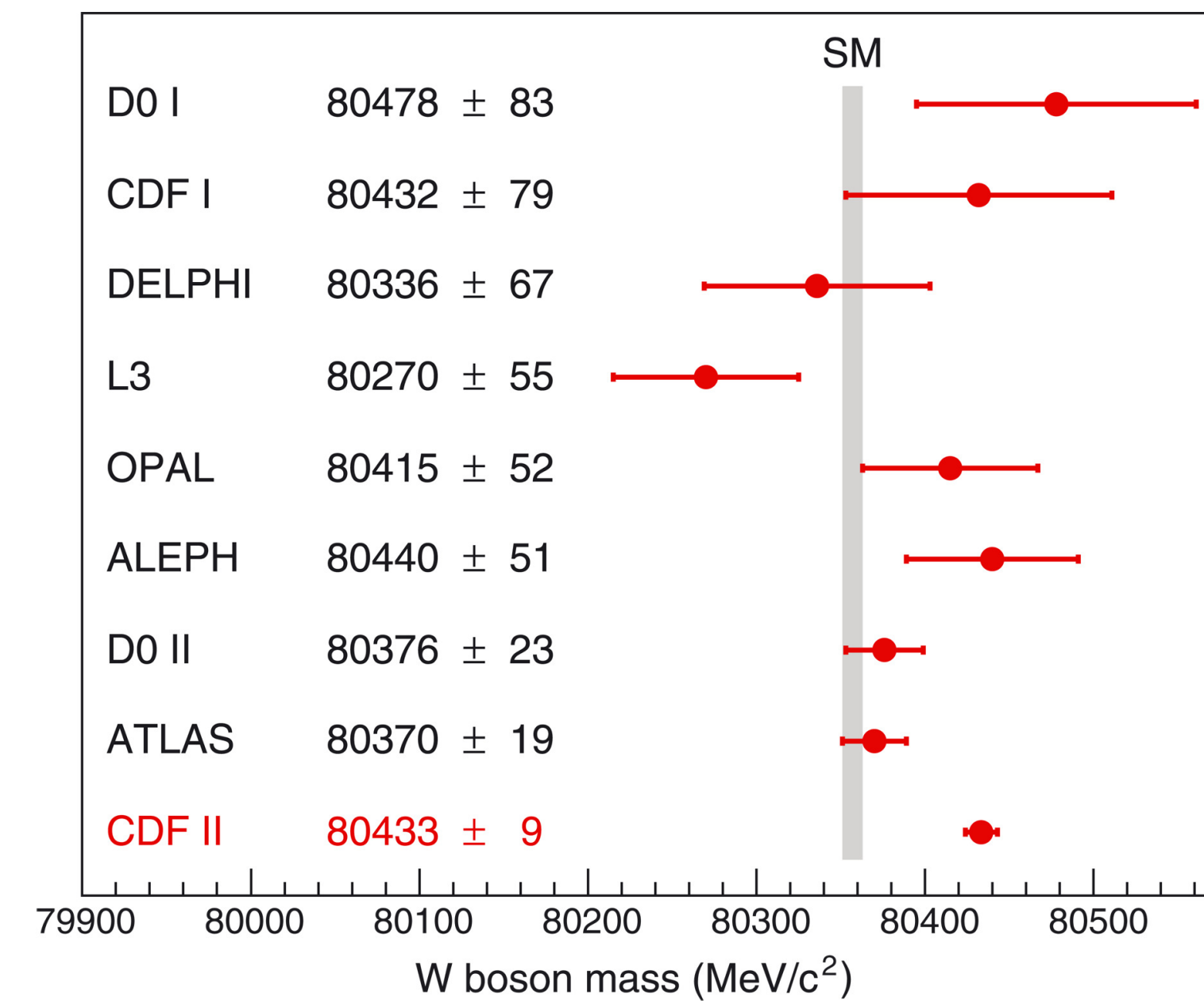
# Anomalies?

$$(g - 2)_\mu$$



g-2 Collaboration, arXiv: 2308.06230

$$M_W$$

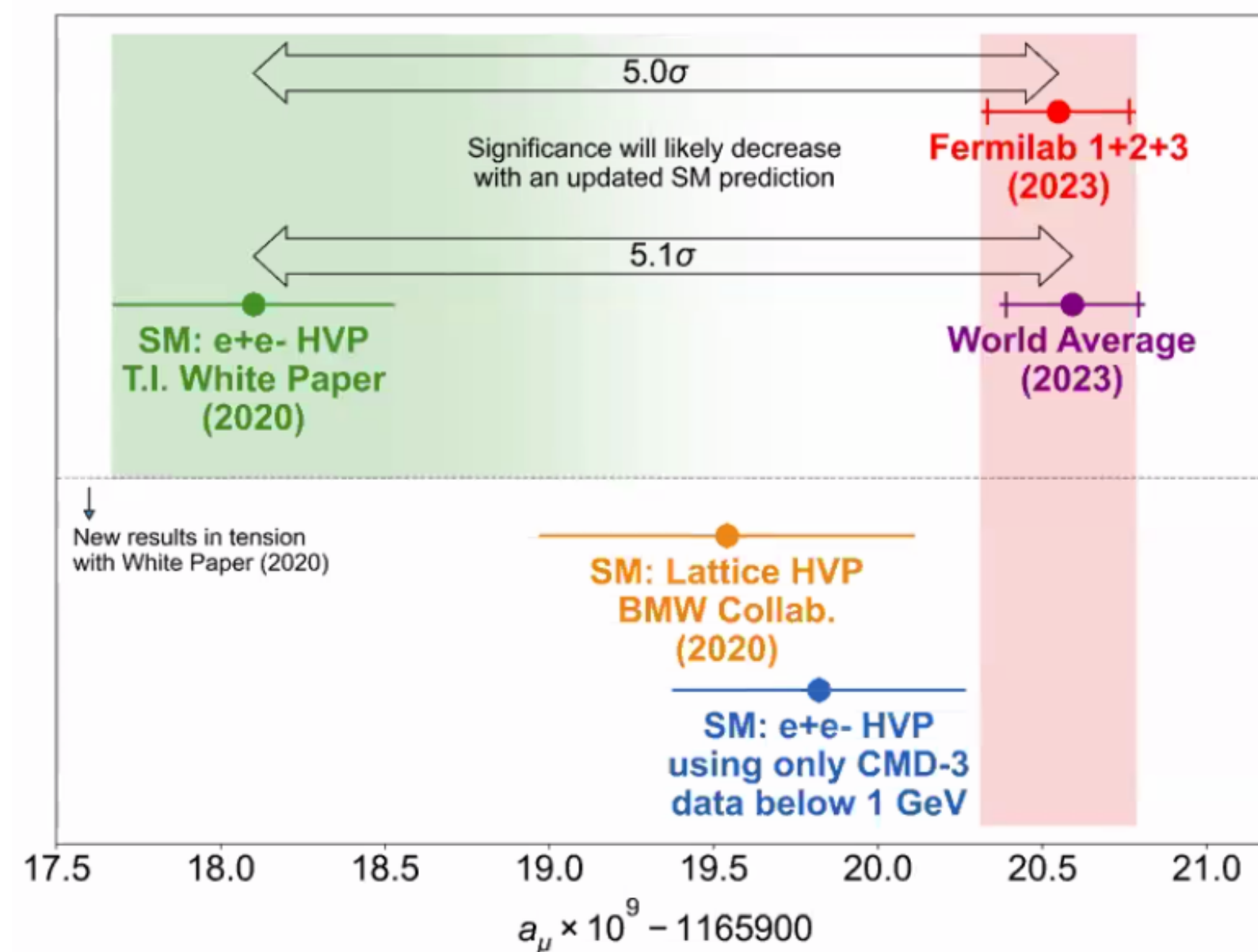


CDF Collaboration, Science 376 (2022), no. 6589 170–176.



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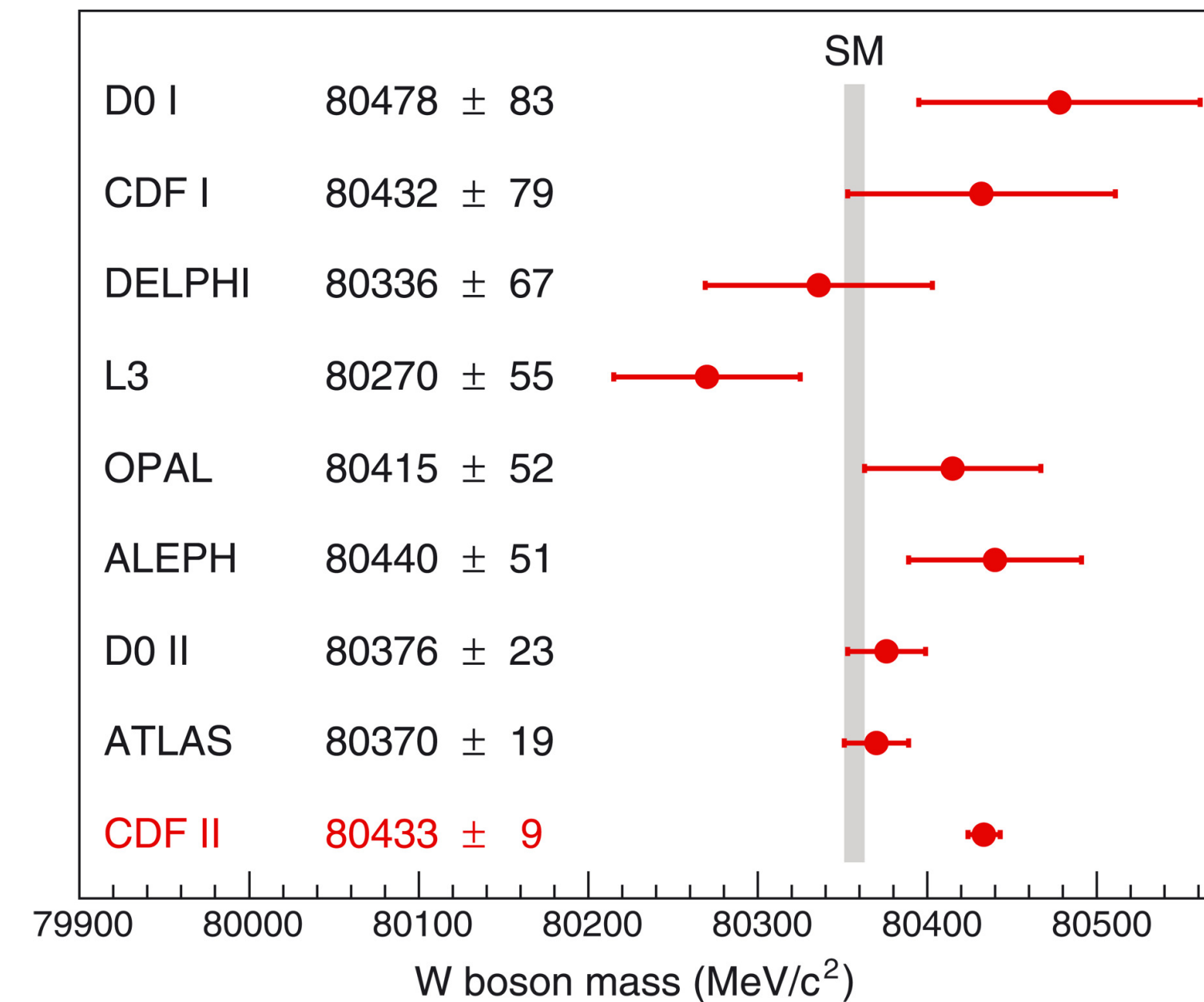
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**Внимание!  
Controversial!**

$$M_W$$



CDF Collaboration, Science 376 (2022), no. 6589 170–176.

# A Minimal Extension

- **Only** sterile neutrinos?

- $m_\nu$  ? ✓

- $M_W$  ? ✓

- $(g - 2)_\mu$  ? ✗




$$\left(\frac{\delta a_\mu}{10^{-9}}\right) \sim \frac{1}{16\pi^2} |\Theta_{\mu N}|^2 G_F m_\mu^2 \sim \mathcal{O}(1) |\Theta_{\mu N}|^2$$



K. Kannike et al., JHEP 02 (2012) 106. [Erratum: JHEP 10, 136 (2012)].  
R. Dermisek and A. Raval, Phys. Rev. D 88 (2013) 013017.  
G. Arcadi et al., Phys. Rev. Lett. 127 (2021), no. 6 061802.  
C.-T. Lu et al., JHEP 08 (2021) 073.  
G. Guedes and P. Olgoso, JHEP 09 (2022) 181.

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	$SU(2)_L$	$U(1)_Y$	$U(1)_L$
$\ell_L$	<b>2</b>	-1/2	1
$\mu_R$	1	1	1
$H$	<b>2</b>	+1/2	0
$N_R$	1	1	1
$S_R$	1	1	-1
$\psi_L$	<b>2</b>	-1/2	1
$\psi_R$	<b>2</b>	-1/2	1

**Minimal**  
extension in the  
leptonic sector?



**2 Vector-Like  
Leptons**  
coupling to  $\mu$

K. Kannike et al., JHEP 02 (2012) 106. [Erratum: JHEP 10, 136 (2012)].  
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# Extended Seesaw

Small Symmetry Breaking Parameters

$$\begin{aligned}
 -\mathcal{L}_Y = & \bar{\ell}_L H Y_\mu \mu_R + \bar{\ell}_L \tilde{H} Y_N N_R + \epsilon \bar{\ell}_L \tilde{H} Y_S S_R + \frac{1}{2} \mu' \bar{N}_R^c N_R + \frac{1}{2} \mu \bar{S}_R^c S_R + \Lambda \bar{N}_R^c S_R + \\
 & + Y_R \bar{\psi}_L H \mu_R + Y_V \bar{S}_R^c \tilde{H}^\dagger \psi_R + Y'_V \bar{\psi}_L \tilde{H} N_R + M_\psi \bar{\psi}_L \psi_R + M_L \bar{\ell}_L \psi_R + \text{h.c.},
 \end{aligned}$$

New!

$$\chi \equiv (\nu_L, N_R^c, S_R^c, \psi_L^0, \psi_R^{0c})^T$$

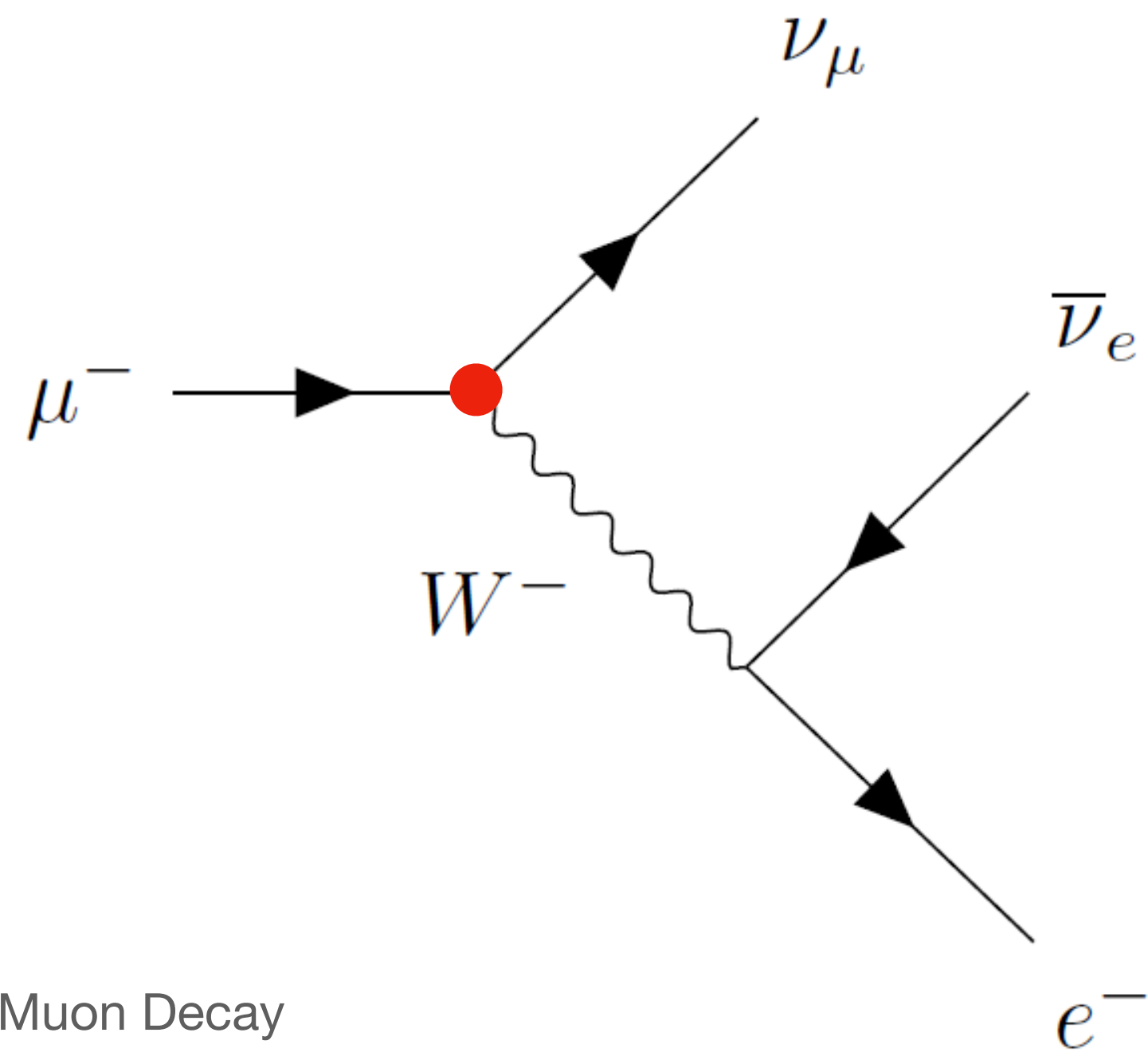
$$-\mathcal{L}_Y \supset \frac{1}{2} \bar{\chi} \mathcal{M}_\chi \chi^c$$

$$\mathcal{M}_\chi = \begin{pmatrix} 0 & m_N & \epsilon m_S & 0 & M_L \\ m_N & 0 & \Lambda & m_{V'} & 0 \\ \epsilon m_S & \Lambda & \mu & 0 & m_V \\ 0 & m_{V'} & 0 & 0 & M_\psi \\ M_L & 0 & m_V & M_\psi & 0 \end{pmatrix}$$

$$m_\nu \simeq \frac{v^2}{2} \left[ \left( Y_N \frac{1}{\Lambda^T} \mu \frac{1}{\Lambda} Y_N^T \right) - \epsilon \left( Y_S \frac{1}{\Lambda} Y_N^T + Y_N \frac{1}{\Lambda^T} Y_S^T \right) \right]$$



$M_W$  is computed from **three** input parameters:  $G_F$ ,  $\alpha_{em}$ ,  $M_Z$



Muon Decay

$$G_F = G_\mu(1 + \Delta_G)$$

$$M_W = M_Z \sqrt{\frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\pi \alpha (1 - \Delta_G)}{\sqrt{2} G_\mu M_Z^2 (1 - \Delta r)}}}$$

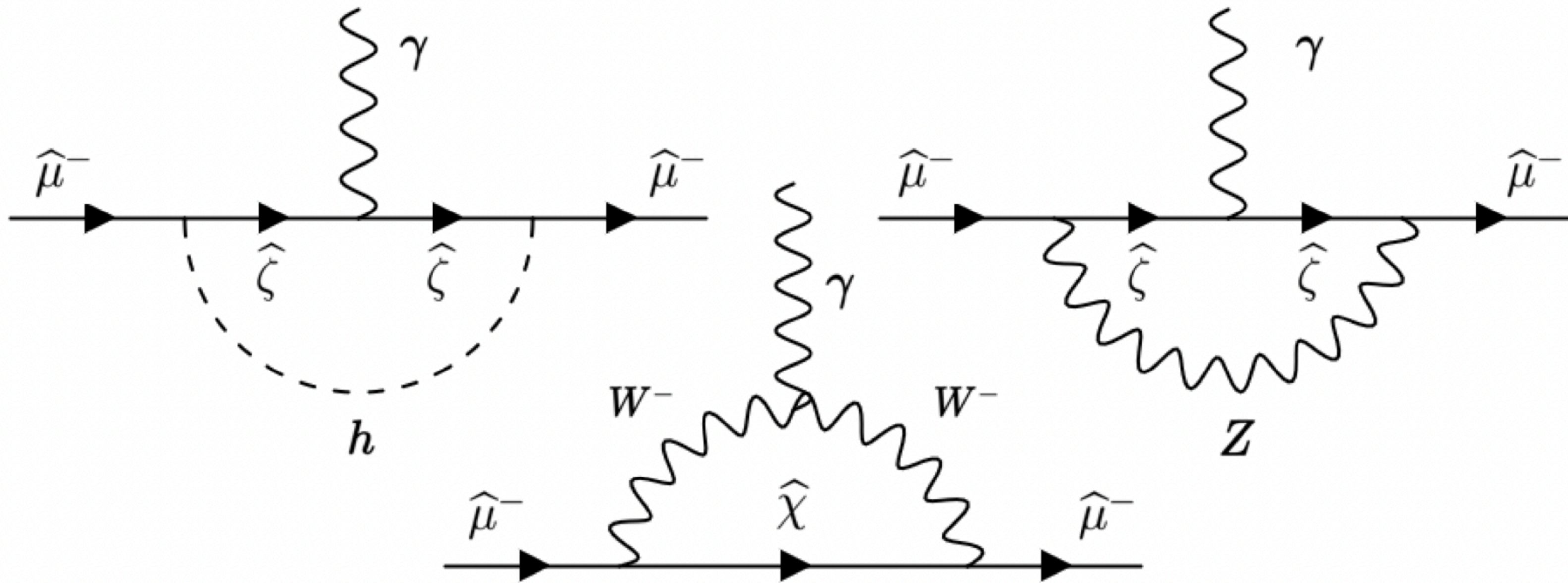
$$\Delta_G \sim 5 \times 10^{-3}$$

$(g - 2)_\mu$

$$\frac{1}{2m_\ell} \bar{\ell} \sigma^{\mu\nu} \ell F_{\mu\nu}$$

Dimension-5 LR Operator  
→ **mass insertion**

$$\chi \equiv (\nu_L, N_R^c, S_R^c, \psi_L^0, \psi_R^{0c})^T, \quad \zeta \equiv (\mu, \psi^-)^T$$



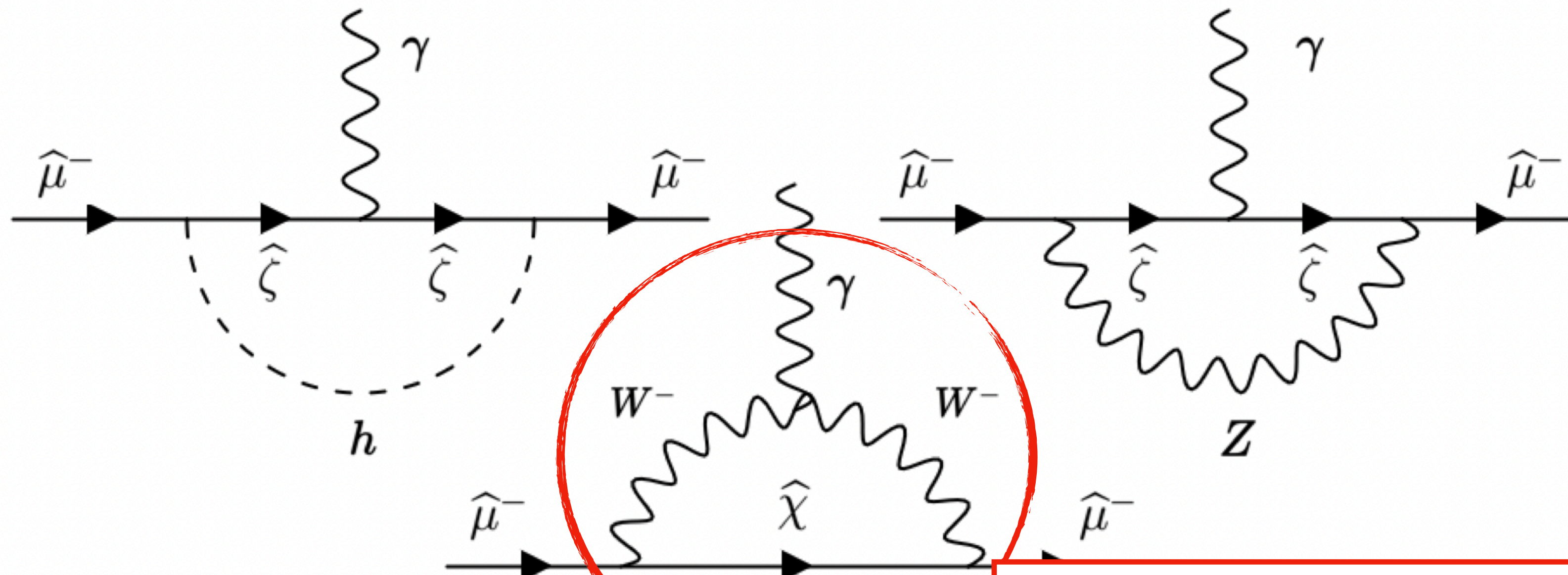


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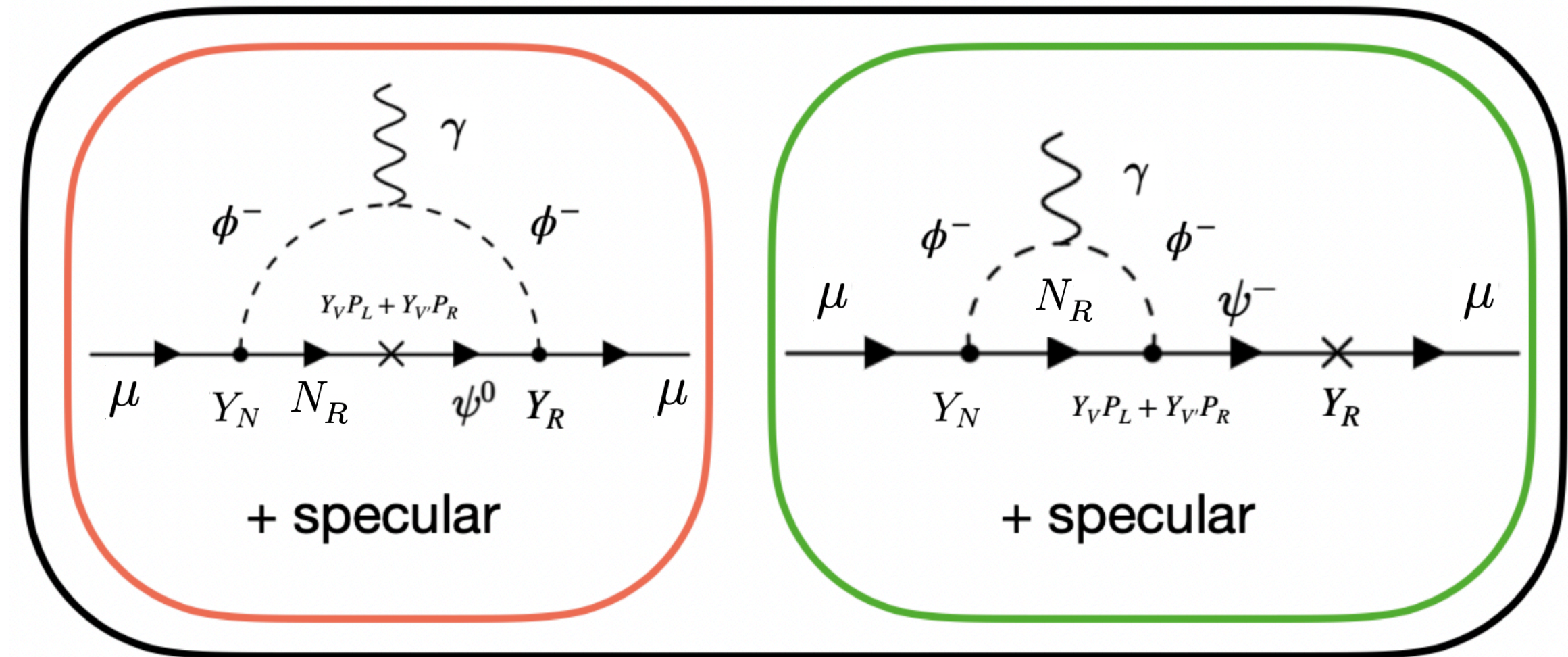


Chirally enhanced contribution

$(g - 2)_\mu$

What to expect?

$$\delta a_\mu \sim \frac{Y_i^3}{16\pi^2} \frac{m_\mu}{v} \left[ \mathcal{O} \left( \frac{v^2}{M_\psi \Lambda} \right) + \mathcal{O} \left( \frac{v^2}{M_\psi \Lambda} \right)^2 \right]$$



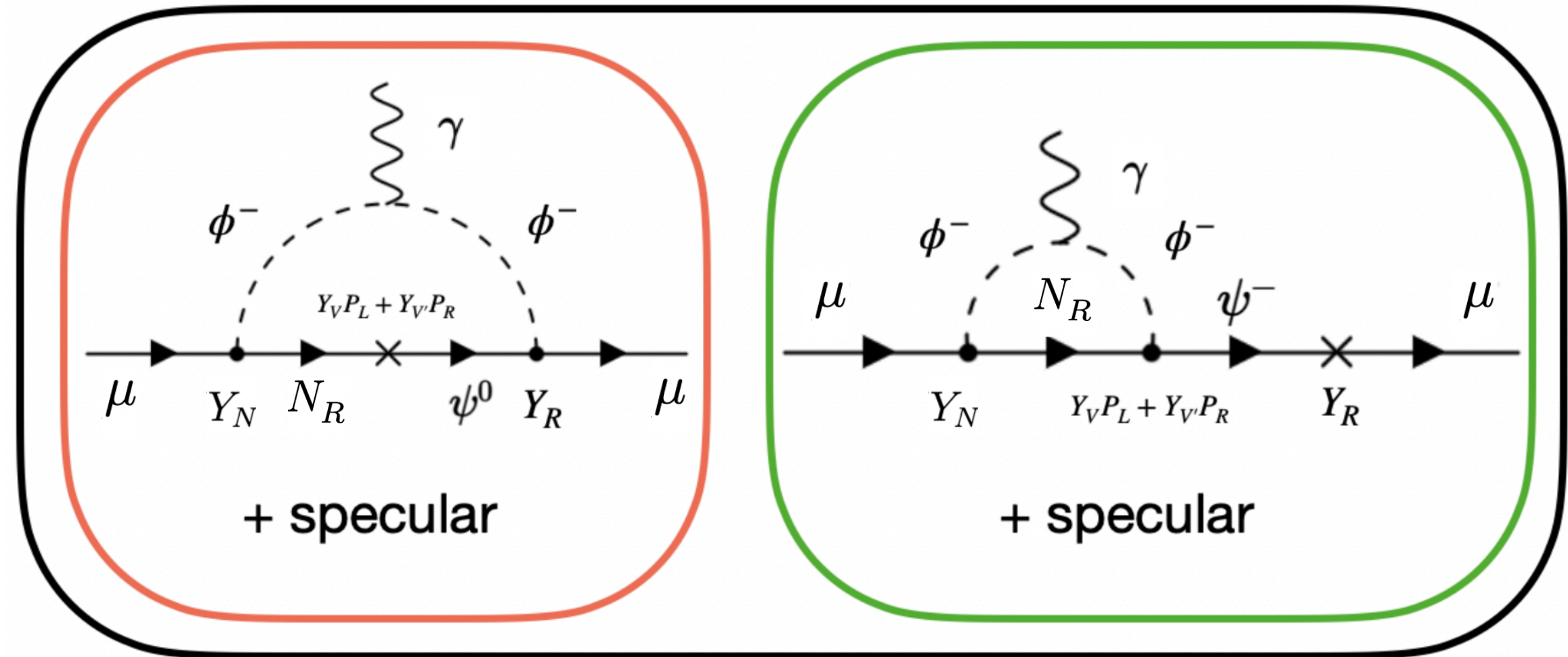


$(g - 2)_\mu$

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$$M_\psi, \Lambda \sim \mathcal{O}(10) \text{ TeV}$$



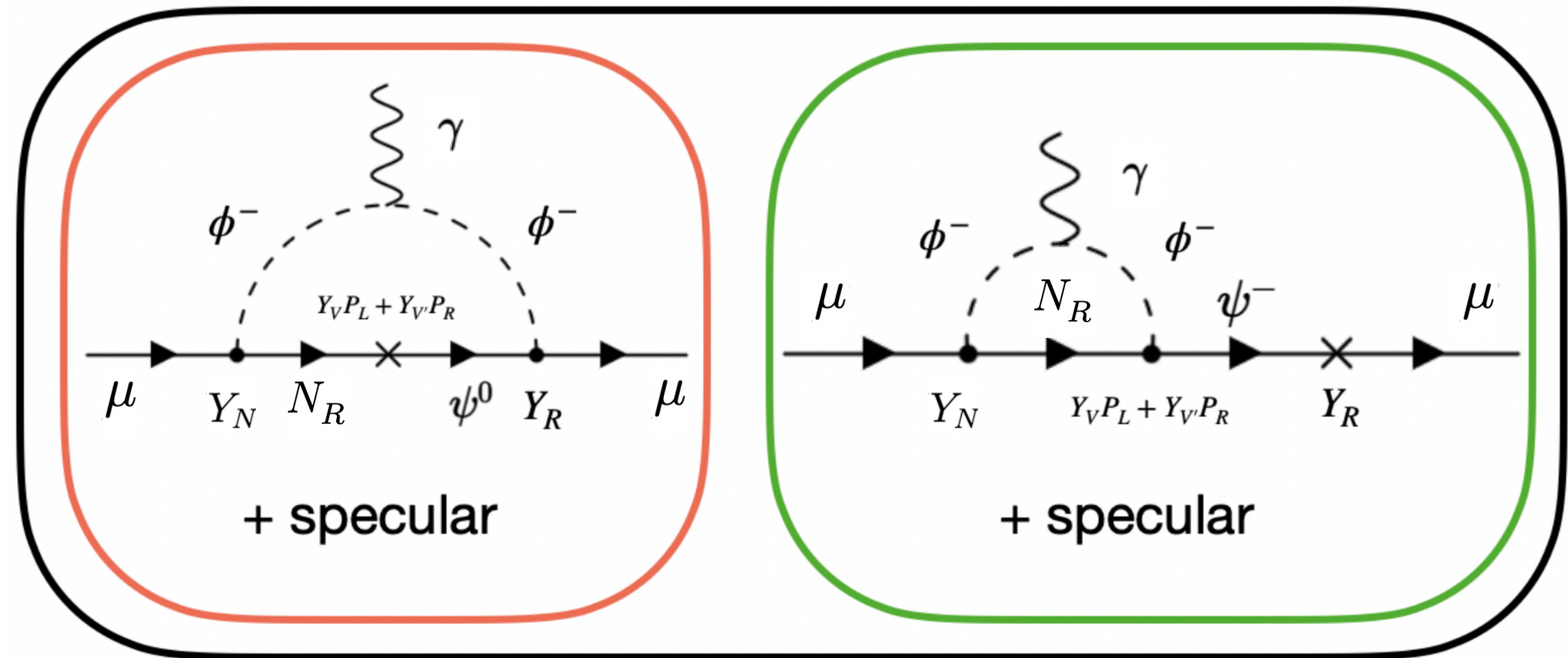
$$(g - 2)_\mu$$

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$$M_\psi, \Lambda \sim \mathcal{O}(1) \text{ TeV}$$



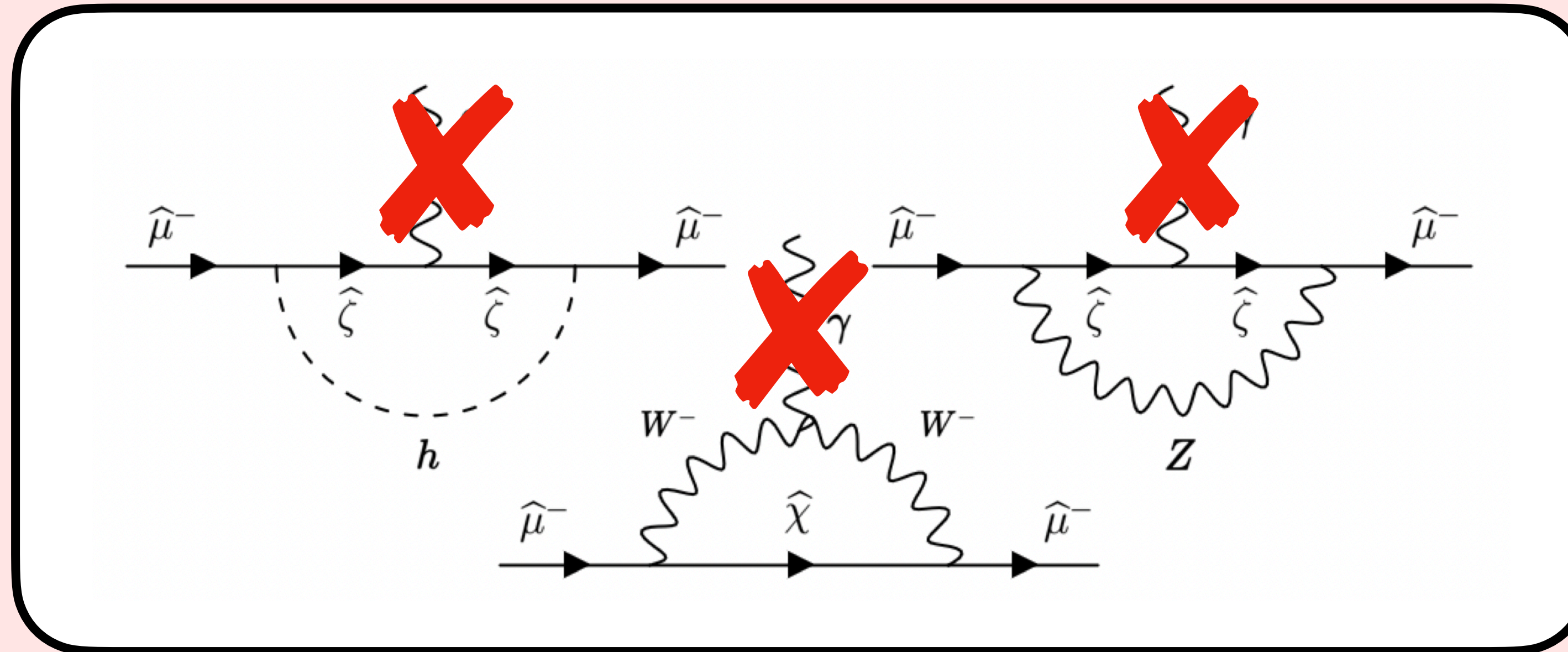
**Accidental** cancellation of the LO

→ violation of Wilsonian naturalness

N. Arkani-Hamed and K. Harigaya, JHEP 09 (2021) 025.  
 N. Craig et al., JHEP 05 (2022) 079.  
 L. Delle Rose et al., JHEP 05 (2022) 120.



# Lepton masses radiatively?

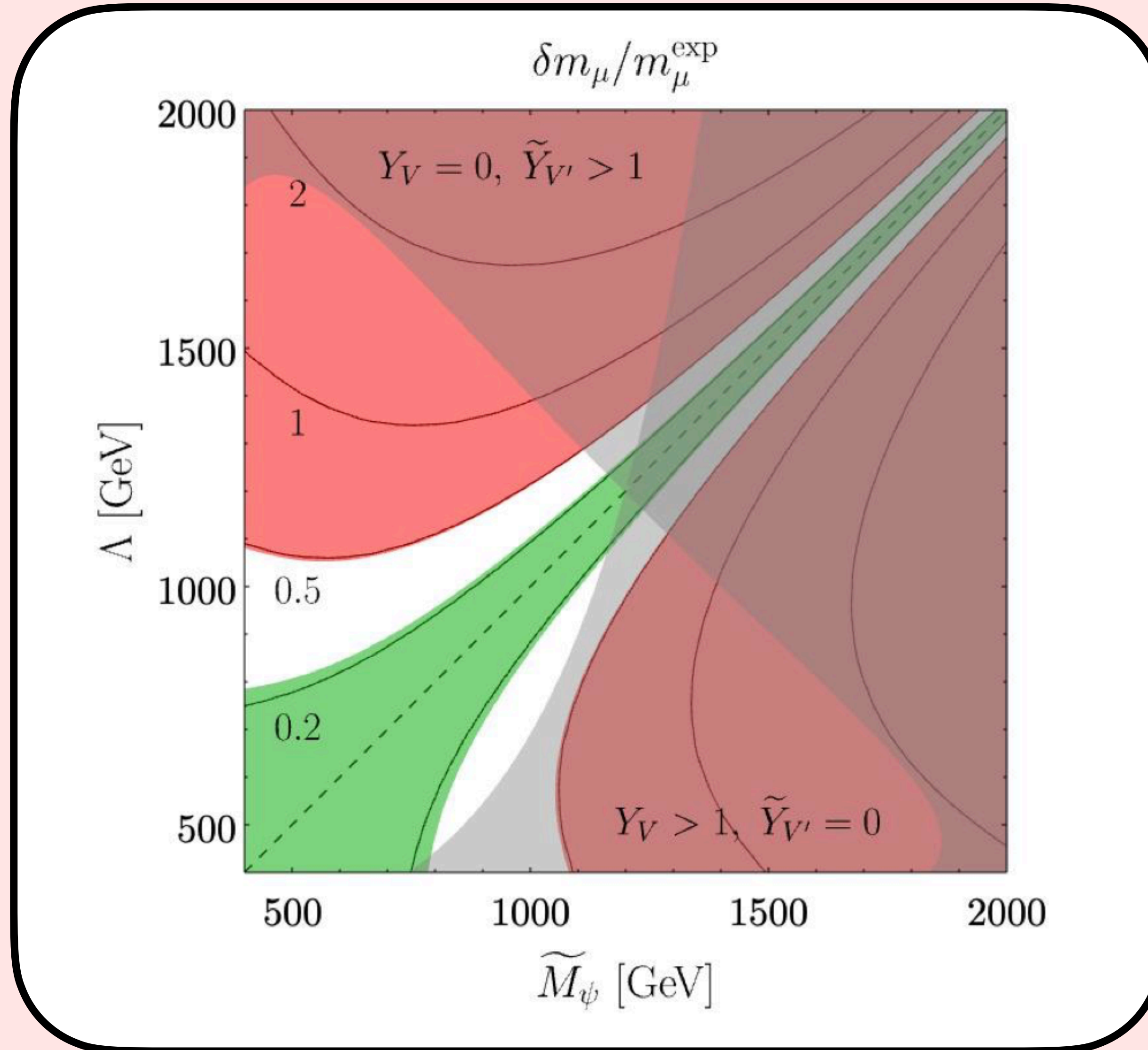


$$\delta m_\mu \propto \delta a_\mu$$

Could this type of models be used for “**charged seesaws**” and generate the **lepton masses radiatively**?



# Parameter Space and Summary



$$m_\mu^{\text{exp}} = m_\mu^{\text{TL}} + \delta m_\mu$$

$$|\delta m_\mu| < 30\% m_\mu^{\text{TL}}$$

$$|\delta m_\mu| > |m_\mu^{\text{TL}}|$$

# Summary

- New particles:

**2 HNLs +  
Vector-Like  
EW-Doublet**

	$SU(2)_L$	$U(1)_Y$	$U(1)_L$
$N_R$	1	1	1
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- Possible **simultaneous** explanation of  $m_\nu$ ,  $(g - 2)_\mu$  and  $M_W$
- NP at  $\sim \mathcal{O}(1)$  TeV, or even lower!
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**Stay tuned!**



**Thanks!**

**Arturo de Giorgi - 2211.03797**