

TWENTY-FIRST LOMONOSOV CONFERENCE August, 24-30, 2023 ON ELEMENTARY PARTICLE PHYSICS MOSCOW STATE UNIVERSITY

Light QCD exotics at BESIII

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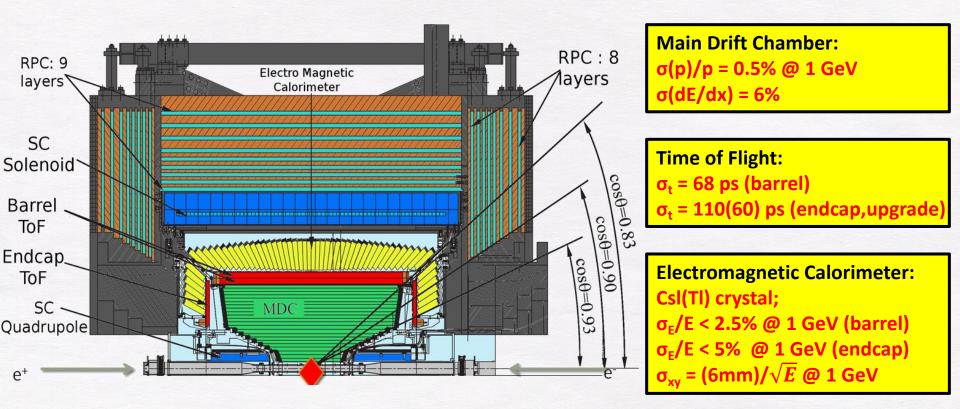
OUTLINE

The BESIII experiment @ BEPCII Hadron spectroscopy Recent results @ BESIII Summary

Beijing Electron Positron Collider II



BESIII Detector



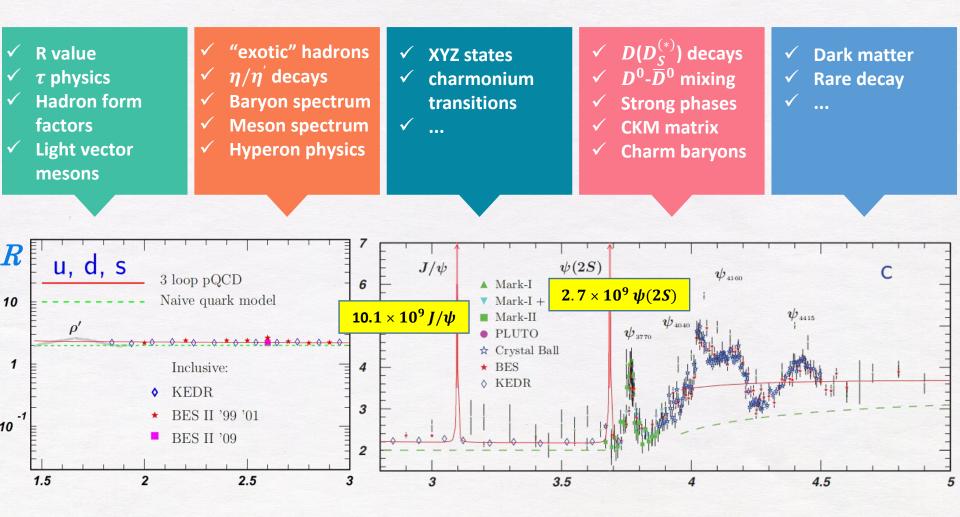
Total weight 730 tons, ~40,000 read-out channels, Data rate: 5kHz, 50Mb/s

In full operation since 2008, still in very good status now!

Magnet: 1T superconducting

Muon Counter: Resistive Plate Chamber σ(xy) < 2 cm

BESIII physics program



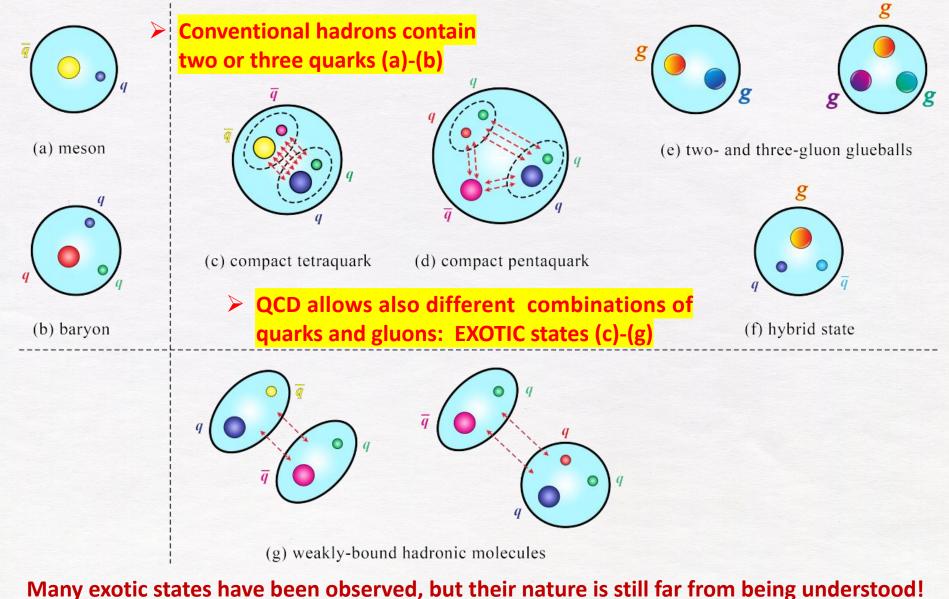
World's largest J/ψ and $\psi(2S)$ data samples directly collected,

~40fb⁻¹ data in $E_{\rm cm}$ = 2-4.95 GeV.

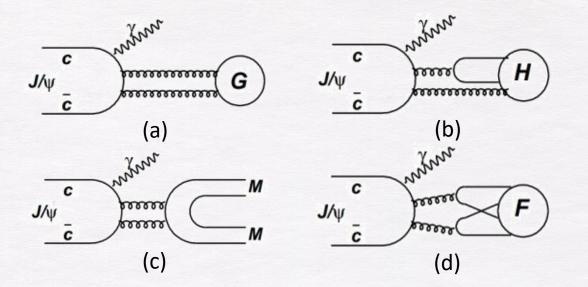
An excellent opportunity for studying light hadron physics!

Hadron spectroscopy

Hadron spectroscopy: Establish the spectrum and study the exotic hadrons properties.



2023-8-30



 $\Gamma(J/\psi o \gamma g g) \cong 6\% \Gamma_{J/\psi}$

- $\succ \Gamma(J/\psi \to \gamma G) \sim \mathcal{O}(\alpha \alpha_s^2)$
- $\succ \Gamma(J/\psi \to \gamma H) \sim O(\alpha \alpha_s^3)$
- $\succ \Gamma(J/\psi \to \gamma M) \sim O(\alpha \alpha_s^4)$
- $\succ \Gamma(J/\psi \to \gamma F) \sim \mathcal{O}(\alpha \alpha_s^4)$

Charmonium radiative decays provide an ideal laboratory for exotic states studies

- High statistics data samples.
- > Well defined initial and final states (clean).
 - ✓ Kinematic constraints
 - ✓ $I(J^{PC})$ filter

 $\succ \text{``Gluon-rich''} environment. \ \Gamma(J/\psi \to \gamma G) > \Gamma(J/\psi \to \gamma H) > \Gamma(J/\psi \to \gamma M) \ge \Gamma(J/\psi \to \gamma F)$

Manifest exotic

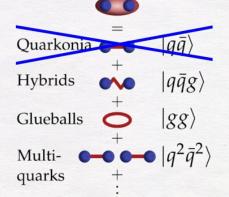
- > Quark contents more than $q\overline{q}$ or qqq.
- Quantum number not reachable for ordinary mesons or baryons.

Ambiguous exotic

- Overpopulation of states.
- Mass/width not fit in spectra.
- Production incompatible with standard mesons/baryons.

- Light-flavor exotic states are hard to establish.
- > Exotic quantum numbers are forbidden by $q\overline{q}$ configuration.
- > Focus on spin-exotic states $J^{PC} = 0^{--}$, even⁺⁻, odd⁻⁺.

J ^{PC}	0++	0+-	0-+	0	1++	1+-	1-+	1	2++	2+-	2-+	2	3++	3+-	3-+	3
$q\overline{q}$	1	×	1	×	1	1	×	1	1	×	1	1	1	1	×	1
$q\overline{q}q\overline{q}$	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
qqg	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
gg	1	×	1	×	1	×	1	×	1	×	1	×	1	×	1	×

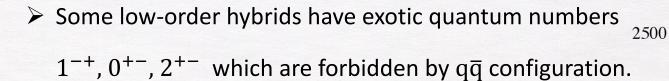


Glueball searches

- Provide critical information on the gluon field.
- Obtain quantitative understanding of confinement.
- \succ 0⁺⁺, 0⁻⁺ and 2⁺⁺ glueballs predicted by LQCD:

Mass under 3.0 GeV.

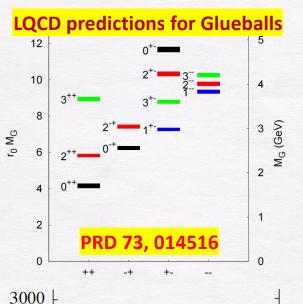
Exotic hybrids



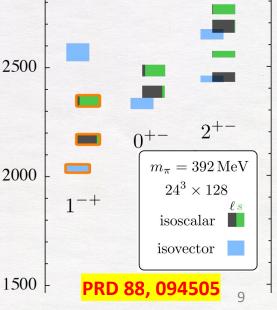
> 1^{-+} hybrids predicted by LQCD: Mass around 2.0 GeV.

An excellent opportunity for studying them using J/ψ or $\psi(2S)$ data samples!

2023-8-30



LQCD predictions for Exotic Hybrids

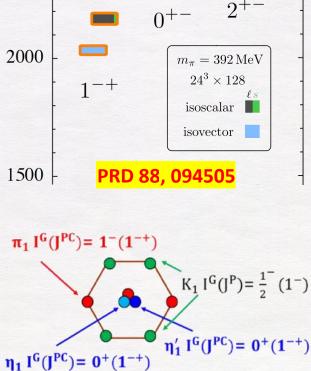


Exotic hybrids

- $\mathbf{1}^{-+}$ nonet of hybrids is predicted to be the lightest.
- Isovector 1^{-+} have been observed including $\pi_1(1400)$, $\pi_1(1600)$ and $\pi_1(2015)$.
- Hybrid is the most popular interpretation.
- $\pi_1(1400)$ and $\pi_1(1600)$ can be explained as one pole. (PRD 105, 012005)
- Isoscalar 1^{-+} is critical for establishing the hybrid nonet >
- Can be produced in the gluon-rich J/ψ radiative decays.
- Can decay to $\eta \eta'$ in P-wave.

(PRD 83, 014021; PRD 83, 014006; EPJP 135, 945)

Observation of states with exotic quantum numbers, such as 1^{-+} , has been of great experimental interest!



LQCD predictions for Exotic Hybrids

3000

2500

2000

1500

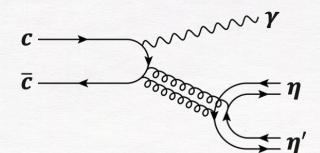
Observation of Exotic Isoscalar State $\eta_1(1855)$ in $J/\psi \rightarrow \gamma \eta \eta'$

PWA of $J/\psi \rightarrow \gamma \eta \eta'$ using 10 Billion of J/ψ data @ BESIII

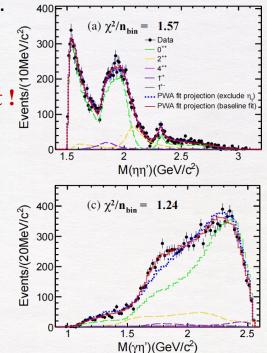
PRL 129, 192002 (2022) PRD 106, 072012 (2022)

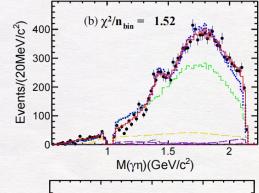
- $\succ \eta \rightarrow \gamma \gamma$ and $\eta' \rightarrow \gamma \pi^+ \pi^- / \eta \pi^+ \pi^-$.
- > An isoscalar 1^{-+} state, $\eta_1(1855)$, has been observed with statistical significance larger than 19σ .
- Consistent with LQCD calculation for 1⁻⁺ hybrid (Mass 1.7 – 2.1 GeV). EPJA 16, 537

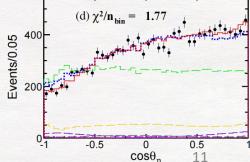
Critical to establish the 1^{-+} hybrid nonet !



Decay mode	Resonance	$M \; ({\rm MeV}/c^2)$	Γ (MeV)	$M_{\rm PDG}~({\rm MeV}/c^2)$	$\Gamma_{\rm PDG}~(MeV)$	B.F. $(\times 10^{-5})$	Sig.
	$f_0(1500)$	1506	112	1506	112	$1.81{\pm}0.11^{+0.19}_{-0.13}$	$\gg 30\sigma$
	$f_0(1810)$	1795	95	1795	95	$0.11{\pm}0.01^{+0.04}_{-0.03}$	11.1σ
	$f_0(2020)$	$2010{\pm}6^{+6}_{-4}$	$203{\pm}9^{+13}_{-11}$	1992	442	$2.28{\pm}0.12^{+0.29}_{-0.20}$	24.6σ
$J/\psi \to \gamma X \to \gamma \eta \eta'$	$f_0(2330)$	$2312{\pm}7^{+7}_{-3}$	$65{\pm}10^{+3}_{-12}$	2314	144	$0.10{\pm}0.02^{+0.01}_{-0.02}$	13.2σ
	$\eta_1(1855)$	$1855 \pm 9^{+6}_{-1}$	$188{\pm}18^{+3}_{-8}$	-	-	$0.27 \pm 0.04^{+0.02}_{-0.04}$	21.4σ
	$f_2(1565)$	1542	122	1542	122	$0.32{\pm}0.05{}^{+0.12}_{-0.02}$	8.7σ
	$f_2(2010)$	$2062{\pm}6^{+10}_{-7}$	$165{\pm}17^{+10}_{-5}$	2011	202	$0.71{\pm}0.06^{+0.10}_{-0.06}$	13.4σ
	$f_4(2050)$	2018	237	2018	237	$0.06{\pm}0.01^{+0.03}_{-0.01}$	4.6σ
	0 ⁺⁺ PHSP		-	-	-	$1.44{\pm}0.15^{+0.10}_{-0.20}$	15.7σ
$J/\psi \to \eta' X \to \gamma \eta \eta'$	$h_1(1415)$	1416	90	1416	90	$0.08{\pm}0.01^{+0.01}_{-0.02}$	10.2σ
$J/\psi o \eta' X o \gamma \eta \eta'$	$h_1(1595)$	1584	384	1584	384	$0.16{\pm}0.02^{+0.03}_{-0.01}$	9.9 <i>σ</i>







2023-8-30

Further checks on the 1^{-+} state $\eta_1(1855)$

Veight sum/(10 MeV/c²)

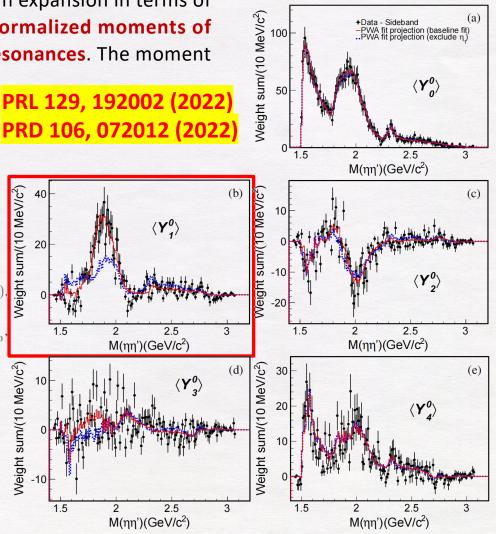
> The $\cos(\theta_{\eta})$ distribution can be expressed as an expansion in terms of Legendre polynomials; the coefficients (unnormalized moments of expansion) characterize the spin of the $\eta\eta$ resonances. The moment for the kth bin of $M(\eta \eta')$ is:

$$\langle Y_l^0 \rangle \equiv \sum_{i=1}^{N_k} W_i Y_l^0 (\cos \theta_\eta^i)$$

The moments are related to the spin-O(S), spin-1(P) and spin-2(D) amplitudes by: $\sqrt{4\pi}\langle Y_0^0 \rangle = S_0^2 + P_0^2 + P_1^2 + D_0^2 + D_1^2 + D_2^2,$ $\sqrt{4\pi}\langle Y_1^0 \rangle = 2S_0P_0 \cos \phi_{P_0} + \frac{2}{\sqrt{5}}(2P_0D_0 \cos(\phi_{P_0} - \phi_{D_0}) + \sqrt{3}P_1D_1 \cos(\phi_{P_1} - \phi_{D_1})),$ $\sqrt{4\pi}\langle Y_2^0 \rangle = \frac{1}{7\sqrt{5}} (14P_0^2 - 7P_1^2 + 10D_0^2 + 5D_1^2 - 10D_2^2) + 2S_0D_0 \cos \phi_{D_0},$ $\sqrt{4\pi}\langle Y_3^0 \rangle = \frac{6}{\sqrt{35}} (\sqrt{3}P_0 D_0 \cos(\phi_{P_0} - \phi_{D_0}) - P_1 D_1 \cos(\phi_{P_1} - \phi_{D_1})),$ $\sqrt{4\pi}\langle Y_4^0 \rangle = \frac{1}{7}(6D_0^2 - 4D_1^2 + D_2^2).$

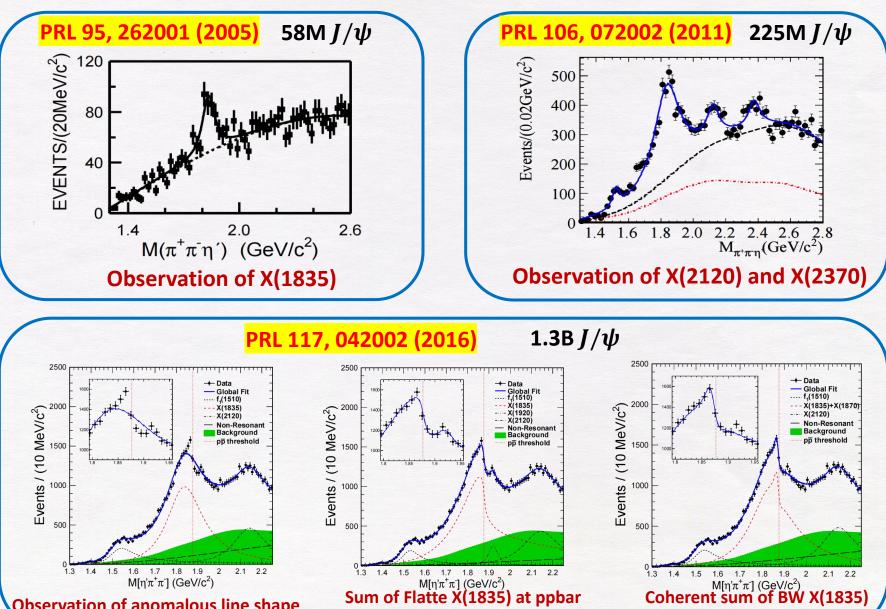
 $< Y_1^0 >$ indicates significant P-wave.

 \blacktriangleright In $\eta\eta$ system, the narrow structure should be $\eta_1(1855)$ P-wave component.



Good agreement between the PWA fit results and data. 12

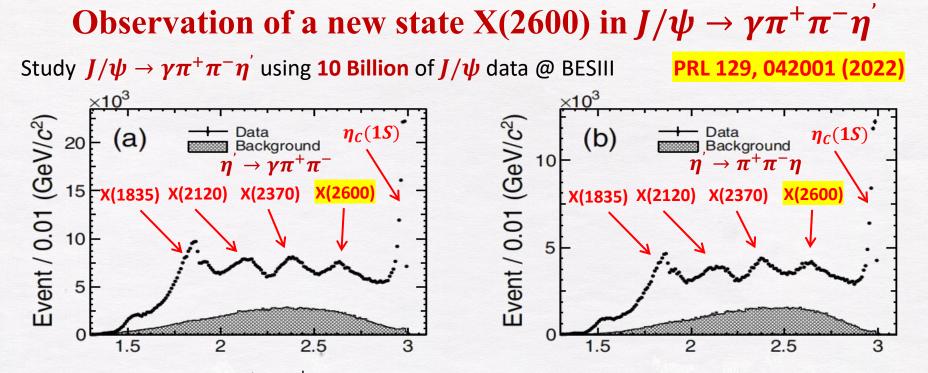
States in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta^-$



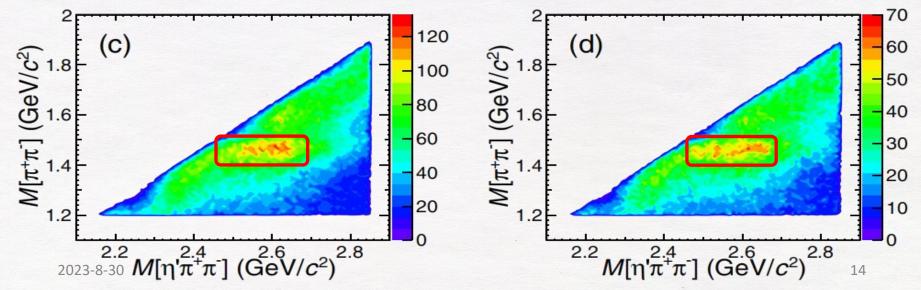
and BW X(1920)



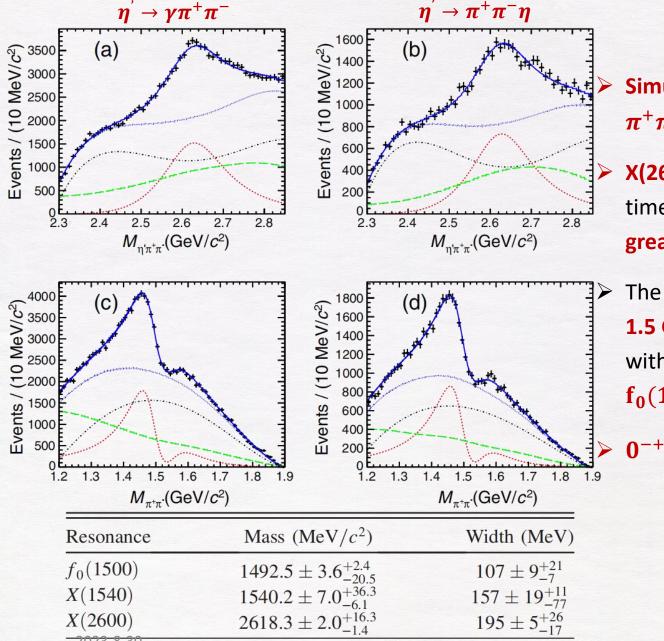
Observation of anomalous line shape



> A new state in M($\pi^+\pi^-\eta'$) invariant mass is observed around 2.6 GeV, which is correlated to a structure in M($\pi^+\pi^-$) at 1.5 GeV.



Observation of a new state X(2600) in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$



Simultaneous fit to $\pi^+\pi^-\eta'$ and $\pi^+\pi^-$ mass spectra is performed. X(2600) observed for the first time with a statistical significance

PRL 129, 042001

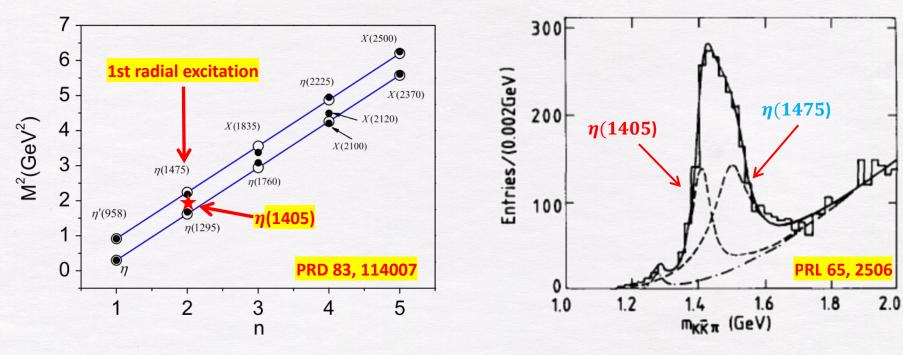
greater than 20σ .

The structure in $M(\pi^+\pi^-)$ around **1.5 GeV** can be well described with the interference between $f_0(1500)$ and X(1540).

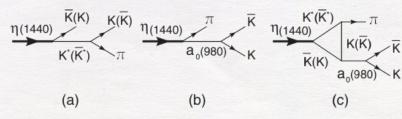
0⁻⁺ **or 2**⁻⁺ is favored for X(2600).

 η radial excitation? an exotic hadron? J^{PC} ?

"η(1405/1475)" Confusion / "ι" Puzzle



- ➢ Mark III suggested there should exist two 0^{−+} structures called $\eta(1405)$ and $\eta(1475)$.
- > $\eta(1405/1475)$ Confusion:
- Two states or just one? (Triangle Singularity Mechanism)
- Their nature? Glueball / 1st radial excitation states of η' ? 2023-8-30



Long standing puzzle!

PRL 108, 081803

Partial Wave Analysis of $J/\psi \rightarrow \gamma K_s K_s \pi^0$

PWA of $J/\psi \rightarrow \gamma K_s K_s \pi^0$ using 10 Billion of J/ψ data @ BESIII

JHEP 03(2023)121

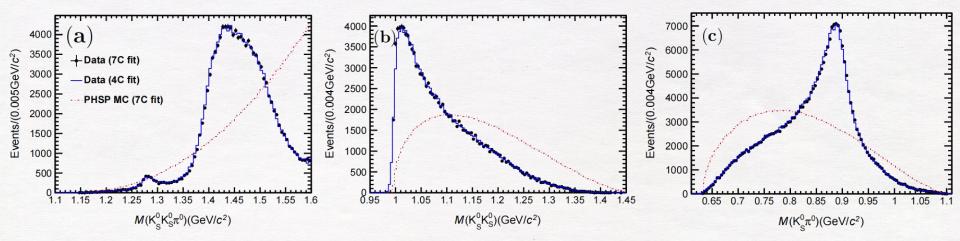
- $\succ \pi^0 \rightarrow \gamma \gamma$ and $K_S^0 \rightarrow \pi^+ \pi^-$.
- Study of $\eta(1405)$ and $\eta(1475)$.
- Prominent structure around 1.45 GeV and a clear bump around 1.28 GeV.

Mass Dependent PWA:

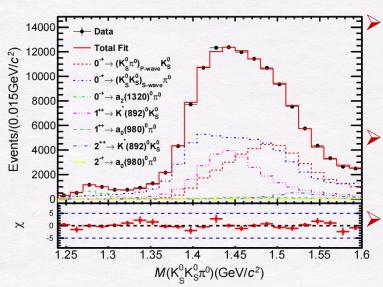
Extract the resonance parameters of the intermediate states by choosing particular dynamic models.

Mass Independent PWA:

Explore the lineshape of $K_s K_s \pi^0$ invariant mass for the different decay modes, and minimize bias from particular dynamic models of the intermediate states.



$J/\psi \rightarrow \gamma K_s K_s \pi^0$ - Mass Dependent PWA

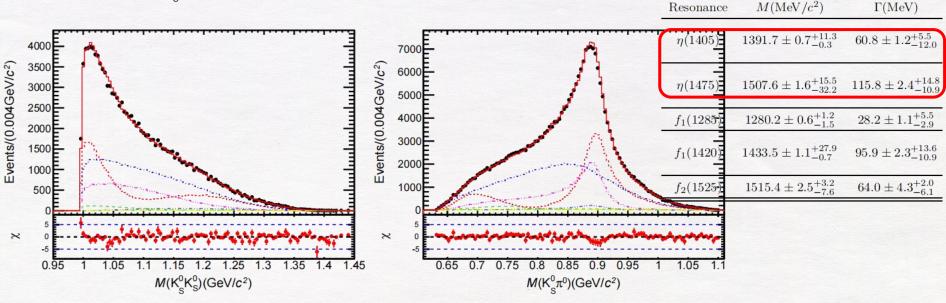


JHEP 03(2023)121

The pseudoscalar components: two states are needed in MD-PWA using relativistic BW model, and both of them can decay by $(K_sK_s)_{S-wave}\pi^0$ and $(K_s\pi^0)_{P-wave}K_s$ modes.

 $f_1(1285)$, $f_1(1420)$ and $f_2(1525)$ are observed for the first time in this process.

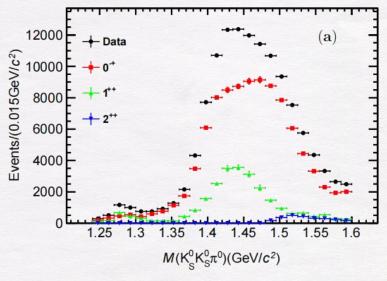
TSM and **one** 0^{-+} **state assumptions:** the description of $K_s K_s \pi^0$ invariant mass spectrum **deteriorates significantly**.



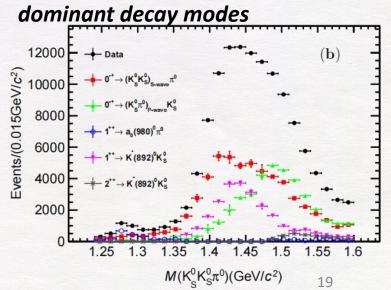
$J/\psi \rightarrow \gamma K_s K_s \pi^0$ - Mass Independent PWA

- > Datasets of MI-PWA: dividing $M(K_sK_s\pi^0)$ from **1.24GeV to 1.6GeV** into 24 bins, the width of each bin is **15MeV**.
- > MI-PWA in each bin: the dynamic function of $M(K_sK_s\pi^0)$ in each bin is approximate to be a constant. Perform PWA in each bin and extract components with significance greater than 5σ .
- > The 0^{-+} component is the dominant contribution:
- The relatively flat lineshape around 1.4 GeV rules out one standard resonance parameterization.
- (K_sK_s)_{S-wave}π⁰ & (K_sπ⁰)_{P-wave}K_s waves are of comparable magnitudes, but with different lineshape and peaks.

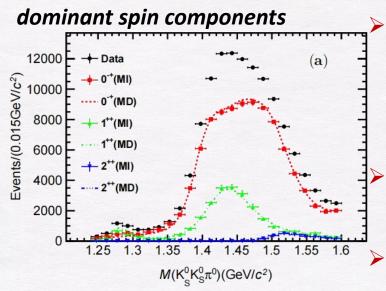
dominant spin components



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$J/\psi \rightarrow \gamma K_s K_s \pi^0$ - Mass Independent PWA

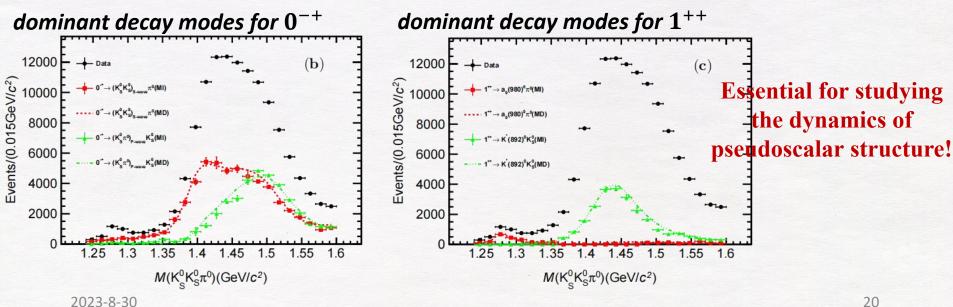


JHEP 03(2023)121

Axial vector component at 1.28 GeV and 1.42 GeV in MI-PWA corresponding to $f_1(1285)$ and $f_1(1420)$ in MD-PWA clearly. The first one decays by $a_0(980)\pi^0$ and the second one decays by $K^{*0}(892)K_s$.

Tensor component around **1.52 GeV** decays by $K^{*0}(892)K_s$ in MI-PWA corresponding to $f_2(1525)$ in MD-PWA clearly.

Consistency between MI and MD results.



Partial Wave Analysis of $J/\psi \rightarrow \gamma \eta' \eta'$

PWA of $J/\psi \rightarrow \gamma \eta' \eta'$ using 10 Billion of J/ψ data @ BESIII

2000 $\blacktriangleright \eta \rightarrow \gamma \gamma$ and $\eta' \rightarrow \gamma \pi^+ \pi^- / \eta \pi^+ \pi^-$. χ^{2} /Nbins = 1.21 χ^2 /Nbins = 1.29 (1200 Ge/(c²) 1000 + Data — Fit Projections GeV(c²) 1500 ** PHSP 800 2020 Events / (0.02 (0.02 ((0.02 $\succ I^{PC} = 0^{++}, 2^{++}, 4^{++} \text{ in } \eta' \eta'.$ 600 2480 1415 Events 400 500 $\succ I^{PC} = 1^{+-}, 1^{--} \text{ in } \gamma \eta'.$ 200 Ng Ng > $f_2(2340)$ observed in $\eta'\eta'$ mode $M_{\eta'\eta'}$ (GeV/c^{2.5}) 1.4 1.6 M_{vn'} (GeV/c² for the first time, its significance is 10 16.1σ. GeV/c²) 10² Events / (0.02 GeV/c²) 10² > $f_2(2340)$ is one potential tensor g_{10} glueball candidate which is close to LQCD prediction for 2^{++} glueballs. 10-1 2.4 22 2.6 1.3 1.4 1.5 1.6 M_{n'n'} (GeV/c²) M_{vm} (GeV/c²) $\Gamma(J/\psi\to\gamma f_2(2340)){\sim}3.0\times10^{-4}$ including:

- $J/\psi \rightarrow \gamma K_S K_S$ PRD 87, 092009
- $J/\psi \rightarrow \gamma \eta \eta$ PRD 98, 072003
- $J/\psi \rightarrow \gamma \phi \phi$ 2023-8-30

PRD 93, 112011

LQCD: $\Gamma(J/\psi \to \gamma G_{2^{++}})/\Gamma_{tot} = 1.1 \times 10^{-2}$ PRL 111, 091601

PRD 105, 072002(2022)

Partial Wave Analysis of $J/\psi \rightarrow \gamma \eta' \eta'$

- > $f_0(2020)$, $f_0(2330)$ observed in $\eta' \eta'$ mode for the first time.
- > $f_0(2480)$: new scalar state observed.
- \succ *f*₀(2020):

PRD 87, 092009 PRD 98, 072003

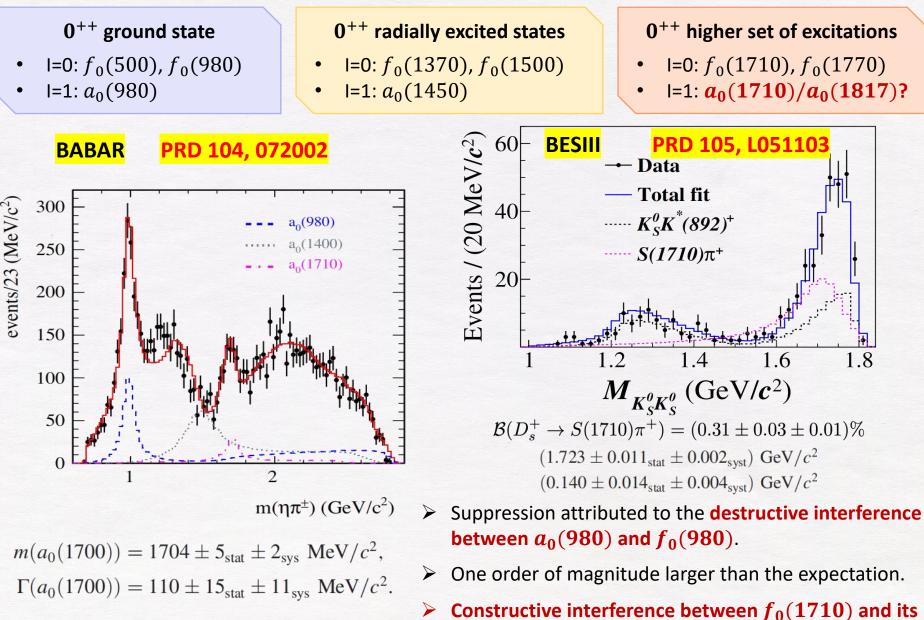
• Production compatible with that of $f_0(2100)$ in $J/\psi \rightarrow \gamma \eta \eta$ and $f_0(2200)$ in $J/\psi \rightarrow \gamma K_S K_S$.

- Its large production rate in radiative J/ψ decay suggest a large **overlap with scalar glueball**.
- Consistent with previous analysis results, but its mass is lower than the mass of the first excitation of scalar glueball from the LQCD prediction. PLB 309, 378 PRD 60, 034509

Resonance	$M(MeV/c^2)$	$\Gamma(MeV)$	B.F.	Significance (σ)
$f_0(2020)$ $f_0(2330)$	$1982 \pm 3^{+54}_{-0} \\ 2312 \pm 2^{+10}_{-0}$	$\begin{array}{c} 436 \pm 4^{+46}_{-49} \\ 134 \pm 5^{+30}_{-9} \end{array}$	$\begin{array}{c} (2.63 \pm 0.06 \substack{+0.31 \\ -0.46}) \times 10^{-4} \\ (6.09 \pm 0.64 \substack{+4.00 \\ +4.68}) \times 10^{-6} \end{array}$	
$f_0(2480)$	$2470 \pm 4^{+4}_{-6}$	$75\pm9^{+11}_{-8}$	$(8.18 \pm 1.77^{+3.73}_{-2.23}) \times 10^{-7}$	new 0 ⁺⁺ state 5.2
$h_1(1415)$	$1384 \pm 6^{+9}_{-0}$	$66 \pm 10^{+12}_{-10}$	$(4.69 \pm 0.80^{+0.74}_{-1.82}) \times 10^{-7}$	5.3
$f_2(2340)$	$2346\pm8^{+22}_{-6}$	$332 \pm 14^{+26}_{-12}$	$(8.67 \pm 0.70^{+0.61}_{-1.67}) \times 10^{-6}$	16.1
0 ⁺⁺ PHSP			$(1.17 \pm 0.23^{+4.09}_{-0.70}) \times 10^{-5}$	15.7

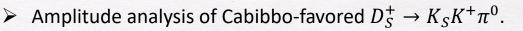
PRD 105, 072002(2022)

Light hadrons in open-charm decays



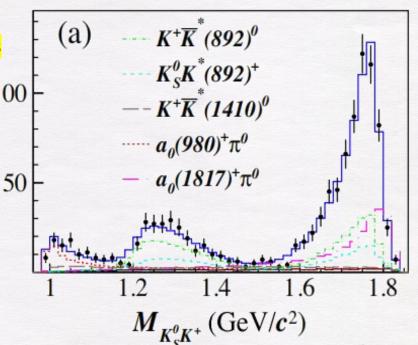
isospin=1 partner.

Light hadrons in open-charm decays



> Together with BESIII analysis $D_S^+ \rightarrow K_S K^+ \pi^0$, this result support the existence of a new a_0 triplet. PRD 105, L051103

- ► BF of $D_S^+ \to a_0^+ \pi^0$ with $a_0^+ \to K_S K^+$ is roughly consistend 100 with the prediction. EPJC 82, 225
- M(a₀) about 100 MeV/c2 greater than the expectation for a₀(1710). Mass = 1.817 ± 0.008 ± 0.020 GeV/c2 Width = 0.097 ± 0.022 ± 0.015 GeV/c2
 a₀(1817) could be the isospin=1 partner of the
- a₀(1817) could be the isospin=1 partner of the f₀(1800) [X(1812)]. PRD 105, 114014



The possibility of the f₀(1710) as the scalar glueball still cannot be excluded.

Amplitude	Phase (rad)	FF (%)	BF (10 ⁻³)	σ
$\overline{D_s^+ \to \bar{K}^*(892)^0 K^+}$	0.0 (fixed)	$32.7 \pm 2.2 \pm 1.9$	$4.77 \pm 0.38 \pm 0.32$	> 10
$D_s^+ \to K^*(892)^+ K_S^0$	$-0.16 \pm 0.12 \pm 0.11$	$13.9 \pm 1.7 \pm 1.3$	$2.03 \pm 0.26 \pm 0.20$	> 10
$D_s^+ \to a_0(980)^+ \pi^0$	$-0.97 \pm 0.27 \pm 0.25$	$7.7 \pm 1.7 \pm 1.8$	$1.12 \pm 0.25 \pm 0.27$	6.7
$D_{\rm s}^+ \to \bar{K}^* (1410)^0 K^+$	$0.17 \pm 0.15 \pm 0.08$	$6.0 \pm 1.4 \pm 1.3$	$0.88 \pm 0.21 \pm 0.19$	7.6
$D_s^+ \to a_0(1817)^+ \pi^0$	$-2.55 \pm 0.21 \pm 0.07$	$23.6 \pm 3.4 \pm 2.0$	$3.44 \pm 0.52 \pm 0.32$	> 10

PRL 129, 182001

Summary

- The BESIII experiment is an excellent laboratory to study light hadron physics, while the charmonium provides an unique opportunity to search for glueball and exotic states.
- > Selection of latest physics results on light exotic states are presented:
 - First observation of exotic isoscalar 1^{-+} state $\eta_1(1855)$ in $J/\psi \rightarrow \gamma \eta \eta'$. PRL 129, 192002 (2022) PRD 106, 072012 (2022)
 - X(1835), X(2120) and X(2370) are confirmed in $J/\psi \rightarrow \gamma \pi^+ \pi^- \eta'$ and X(2600) is observed for the first time. PRL 129, 042001 (2022)
 - Two states of $\eta(1405)$ and $\eta(1475)$ are needed in $J/\psi \rightarrow \gamma K_s K_s \pi^0$ and both MD-PWA and MI-PWA results have been povided. JHEP 03(2023)121
 - The tensor glueball candidate $f_2(2340)$ is observed in $\eta' \eta'$ mode for the first time in $J/\psi \rightarrow \gamma \eta' \eta'$ and a new scalar state $f_0(2480)$ is observed as well. PRD 105, 072002(2022)
 - Light hadron $a_0(1817)$ is observed for the first time in open-charm decays. PRD 105, L051103 PRL 129, 182001
- > The BESIII detector is still in very good status, more surprises at BESIII are expected!

Thanks for your attention!

Backup slide (I)

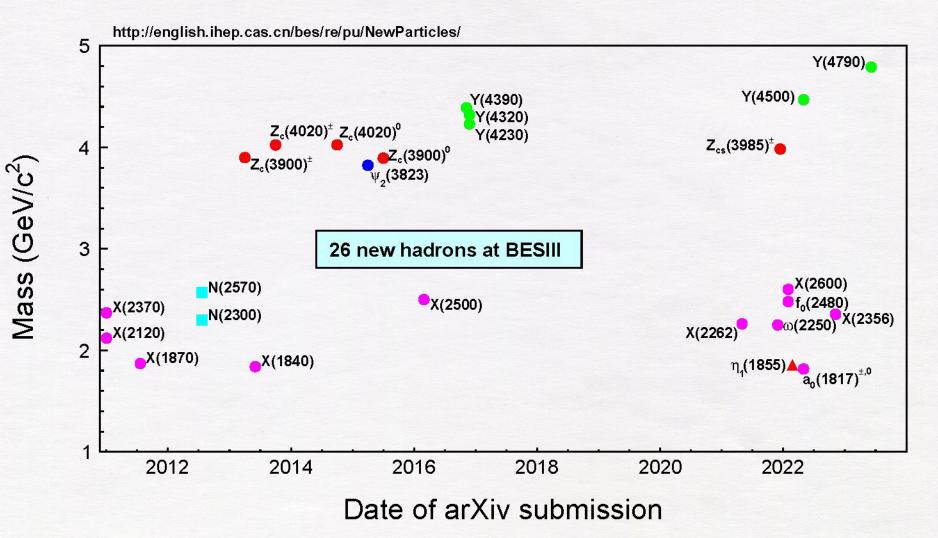
	0+	2+	0-
J/ψ→γPP	$J/\psi \rightarrow \gamma \eta \eta \text{ (pr}$ $J/\psi \rightarrow \gamma \pi^0 \pi^0 \text{ (pr}$ $J/\psi \rightarrow \gamma K_S K_S \text{ (fr}$ $J/\psi \rightarrow \gamma \eta \eta^2 \text{ (pr}$ $J/\psi \rightarrow \gamma \eta^2 \eta^2 \text{ (pr}$	RD92,052003) PRD98,072003) RL129,192002)	
$J/\psi \rightarrow \gamma VV$		J	$\psi \rightarrow \gamma \omega \phi$ (PRD87,032008) $\psi \rightarrow \gamma \phi \phi$ (PRD93,112011) $\psi \rightarrow \gamma \omega \omega$ (PRD100,052012)
J/ ψ →γPPP			$J/\psi \rightarrow \gamma \eta^{2} \pi \pi \text{ (PRL106,072002, noPWA)}$ J/ψ $\rightarrow \gamma K_{S} K_{S} \eta \text{ (PRL115,091803)}$ J/ψ $\rightarrow \gamma K_{S} K_{S} \pi^{0} \text{ (JHEP 03,121)}$

- $J/\psi \to \gamma PP: 0^{++}, 2^{++}, ...$
- $J/\psi \rightarrow \gamma PPP, \gamma VV: 0^{-+}$
- Neutral channel is much cleaner than the charged ones

Amplitude Analysis: toll to extract the complex amplitudes from experimental data

- Models with free parameters
- Consider the kinematic of final states particles
- Vary the parameters to maximize the likelihood
- Mass Dependent (MD) PWA: model the dynamics of particle interactions as coherent sum of resonances
- Mass Independent (MI) PWA: make minimal model assumptions and measure the dynamical amplitudes independently in small regions of two-meson invariant mass (JHEP 03,121 (2023))

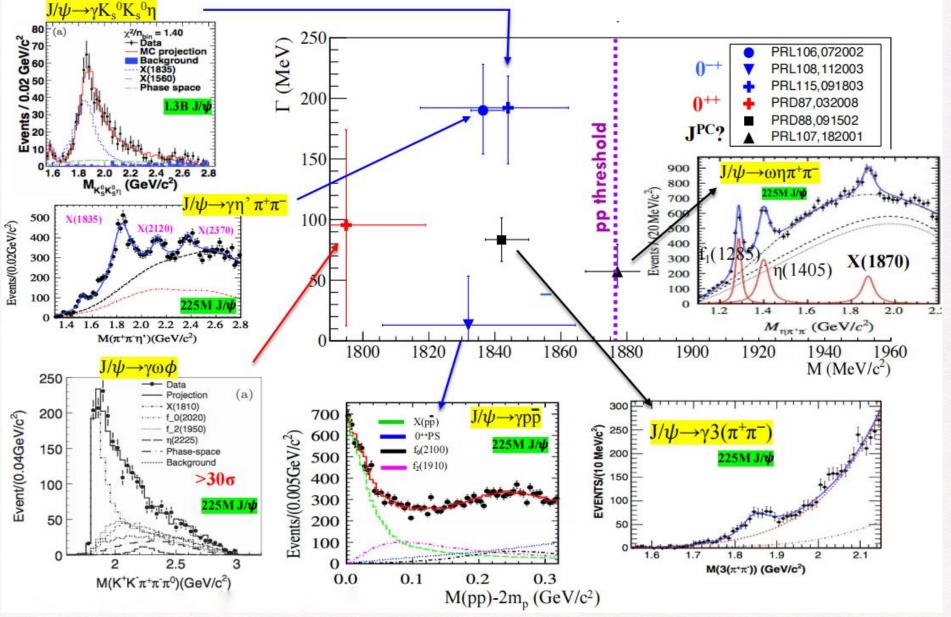
Backup slide (II)



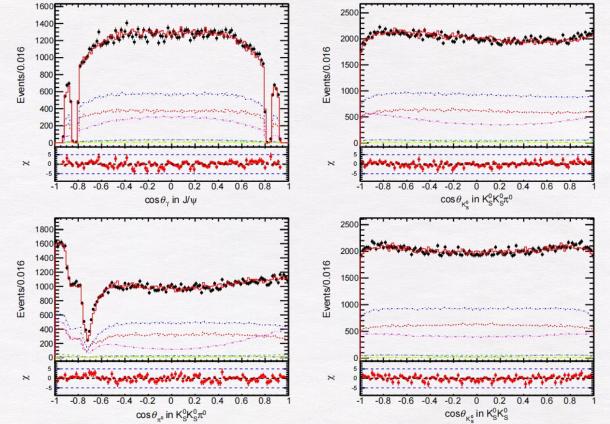
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Backup slide (III)

X(18xx) between 1.8-1.9 GeV



Backup slide (IV)



Resonance	$M({ m MeV}/c^2)$	$\Gamma({ m MeV})$	Decay Mode	B.F.	$\operatorname{Sig.}(\sigma)$
$\eta(1405)$	$1391.7 \pm 0.7^{+11.3}_{-0.3}$	$60.8 \pm 1.2^{+5.5}_{-12.0}$	$J/\psi \to \gamma \eta(1405) \to \gamma K^0_S (K^0_S \pi^0)_{\text{P-wave}} \to \gamma K^0_S K^0_S \pi^0$	$(5.84 \pm 0.12^{+2.03}_{-3.36}) \times 10^{-5}$	$\gg 35$
η(1403)	$1391.7 \pm 0.7 - 0.3$		$J/\psi \to \gamma \eta (1405) \to \gamma (K_S^0 K_S^0)_{\text{S-wave}} \pi^0 \to \gamma K_S^0 K_S^0 \pi^0$	$(2.88 \pm 0.04^{+1.64}_{-0.38}) \times 10^{-5}$	18.4
$\eta(1475)$ 1507.6 $\pm 1.6^{+1}_{-3}$	$1507.6 \pm 1.6^{+15.5}_{-32.2}$	$115.8 \pm 2.4^{+14.8}_{-10.9}$	$J/\psi \to \gamma \eta (1475) \to \gamma K^0_S (K^0_S \pi^0)_{\text{P-wave}} \to \gamma K^0_S K^0_S \pi^0$	$(6.58 \pm 0.12^{+3.98}_{-2.82}) \times 10^{-5}$	$\gg 35$
η(1475)	$1507.0 \pm 1.0_{-32.2}$	$115.0 \pm 2.4 - 10.9$	$J/\psi \to \gamma \eta (1475) \to \gamma (K_S^0 K_S^0)_{\text{S-wave}} \pi^0 \to \gamma K_S^0 K_S^0 \pi^0$	$(3.99 \pm 0.09^{+0.41}_{-0.66}) \times 10^{-5}$	$\gg 35$
$f_1(1285)$	$1280.2 \pm 0.6^{+1.2}_{-1.5}$	$28.2 \pm 1.1^{+5.5}_{-2.9}$	$J/\psi \rightarrow \gamma f_1(1285) \rightarrow \gamma a_0(980)^0 \pi^0 \rightarrow \gamma K^0_S K^0_S \pi^0$	$(8.55 \pm 0.41^{+3.42}_{-1.04}) \times 10^{-6}$	$\gg 35$
$f_1(1420)$	$1433.5 \pm 1.1^{+27.9}_{-0.7}$	$95.9 \pm 2.3^{+13.6}_{-10.9}$	$J/\psi \to \gamma f_1(1420) \to \gamma K^*(892)^0 K^0_S \to \gamma K^0_S K^0_S \pi^0$	$(7.25 \pm 0.12^{+0.73}_{-1.25}) \times 10^{-5}$	$\gg 35$
$J_1(1420)$ 1433.3 ± 1	$1455.5 \pm 1.1_{-0.7}$	$95.9 \pm 2.3_{-10.9}$	$J/\psi \to \gamma f_1(1420) \to \gamma a_0(980)^0 \pi^0 \to \gamma K^0_S K^0_S \pi^0$	$(4.62 \pm 0.36^{+2.36}_{-1.94}) \times 10^{-6}$	17.8
$f_2(1525)$	$1515.4 \pm 2.5^{+3.2}_{-7.6}$	$64.0 \pm 4.3^{+2.0}_{-6.1}$	$J/\psi \to \gamma f_2(1525) \to \gamma K^*(892)^0 K^0_S \to \gamma K^0_S K^0_S \pi^0$	$(9.47 \pm 0.43^{+1.51}_{-0.66}) \times 10^{-6}$	23.8
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Backup slide (V)

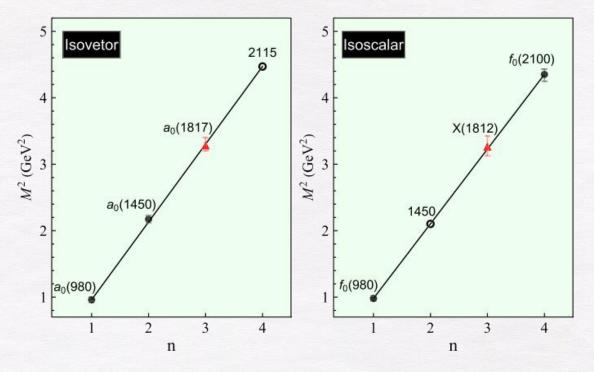


FIG. 1. Two Regge trajectories for the isovector and isoscalar scalar states. Here, the red triangles denote the $a_0(1817)$ and X(1812), while the solid points and empty circles are the experimental and predicted states, respectively. Except the $a_0(1817)$, X(1812) and predicted $a_0(2115)$, $f_0(1450)$, the masses of the other scalar states are taken from PDG. Here, the established states are marked by the black solid points, while the predicted states are denoted by the black circles. The error bars present total experimental uncertainties.