

Self-similar growth of axion stars



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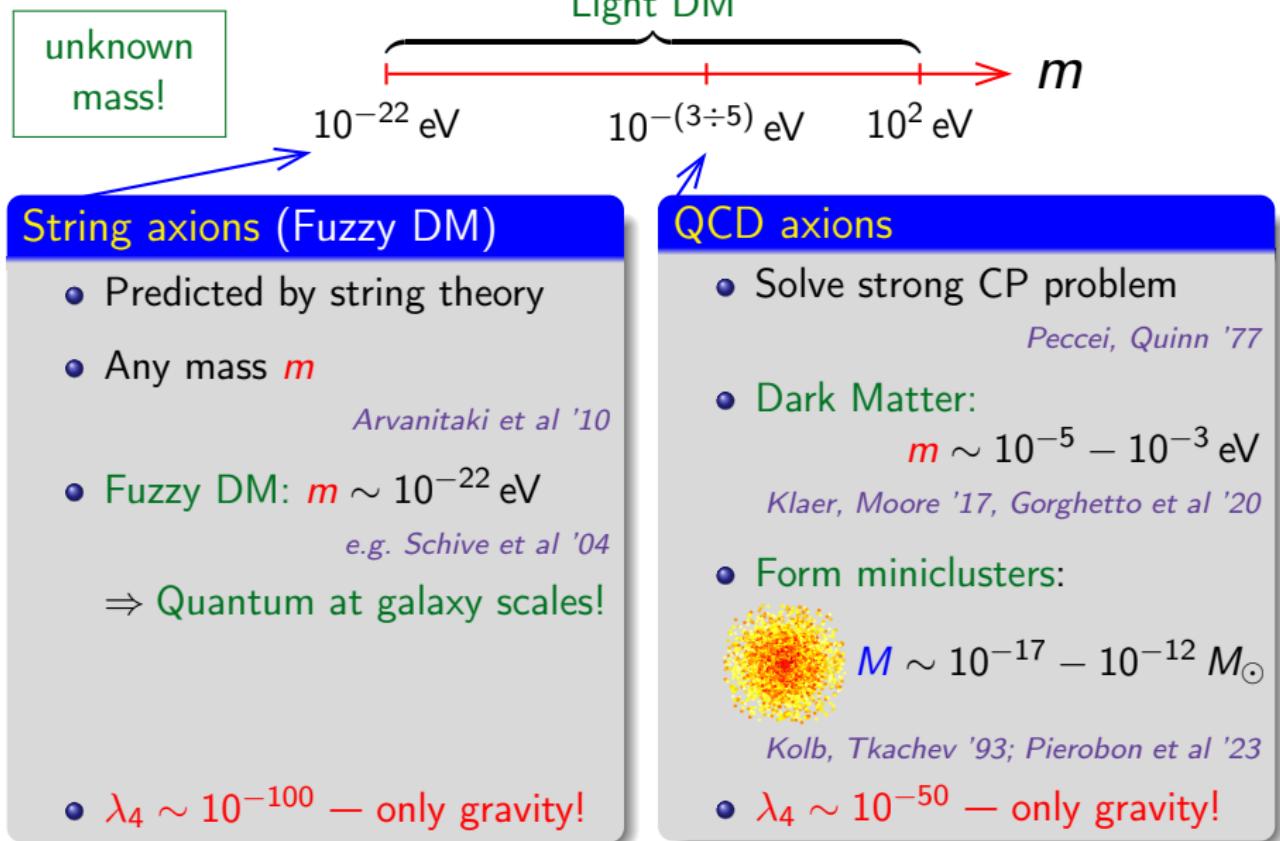


**TWENTY-FIRST LOMONOSOV
CONFERENCE** August, 24-30, 2023
ON ELEMENTARY PARTICLE PHYSICS
MOSCOW STATE UNIVERSITY

DL, A. Panin, I. Tkachev, PRL 121 (2018) 151301

A. Dmitriev, DL, A. Panin, I. Tkachev, arXiv:2305.01005

Light bosonic (axion-like) dark matter

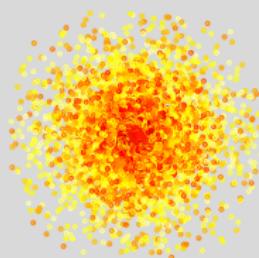


Any small structure

a dwarf galaxy



or axion minicluster



ρ, \mathbf{v} — known!
density, velocity

- Large phase-space density!

$$m \ll 10^2 \text{ eV} \Rightarrow f \sim \frac{\rho/m}{(m\mathbf{v})^3} \gg 1$$

 \Rightarrow classical field $\psi(t, \mathbf{x})$

- Nonrelativistic approximation

$$v \ll 1 \Rightarrow$$

Schrödinger-Poisson (SP) eqs

$$i\partial_t \psi = -\Delta \psi / 2m + m \mathbf{U} \psi$$

$$\Delta \mathbf{U} = 4\pi G m |\psi|^2$$

grav. potential $\mathbf{U}(t, \mathbf{x})$ field $\psi(t, \mathbf{x})$

Rich wave (quantum) phenomena?

Light DM Bose-condenses by gravitational scattering!

- **Gravitational kinetic relaxation:**

Levkov, Panin, Tkachev '18

$$t_{gr} = \frac{4\sqrt{2}m}{\sigma_{gr}} \frac{b}{v \rho f}$$

Rutherford cross section

$$\sigma_{gr} \propto (mG)^2 \Lambda / v^4$$

phase-space density

$$f \propto (\rho/m) / (mv)^3$$

$$t_{gr} = \frac{b\sqrt{2}}{12\pi^3} \frac{m^3 v^6}{G^2 \Lambda \rho^2}, \quad b \approx 0.9$$

Coulomb logarithm:
 $\Lambda \equiv \log(mvR)$

- $v \ll 1$
- $f \gg 1$

gravity is enhanced &
beats self-coupling $\sigma_{gr} \gg \sigma_\lambda$

- Long-range $\sigma_{gr} \propto v^{-4}$
- Bose-factor f

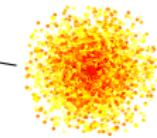
- Fuzzy DM in dwarf galaxies:

$$t_{gr} \gtrsim 10^3 \text{ yr}$$



- QCD axions in miniclusters:

$$t_{gr} \gtrsim hr$$



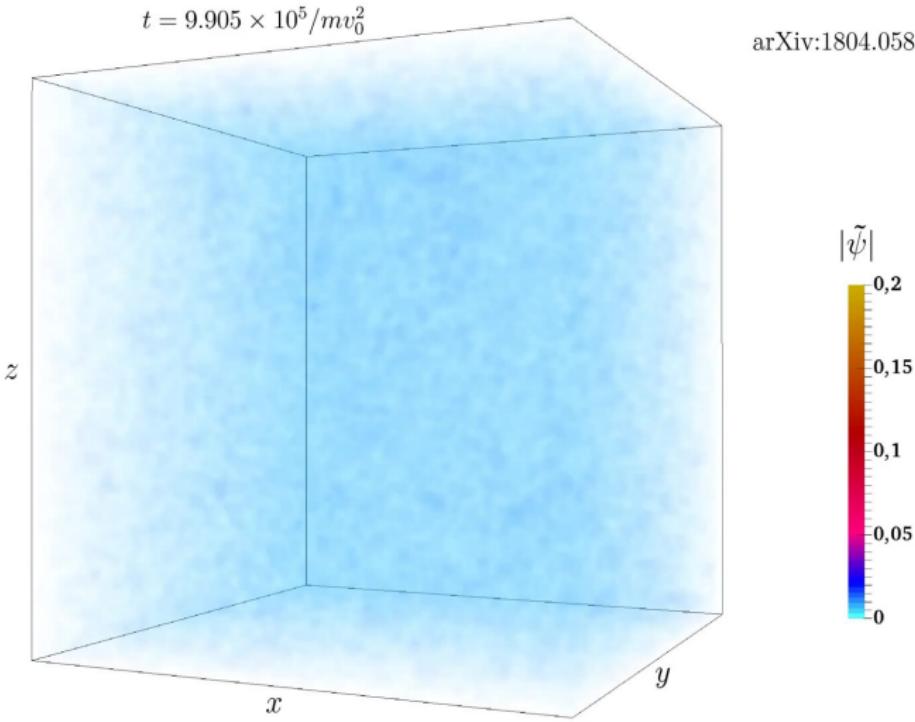
- $f \gg 1$: Relaxation \Rightarrow Bose-condensation

Simulation: solve SP equations!

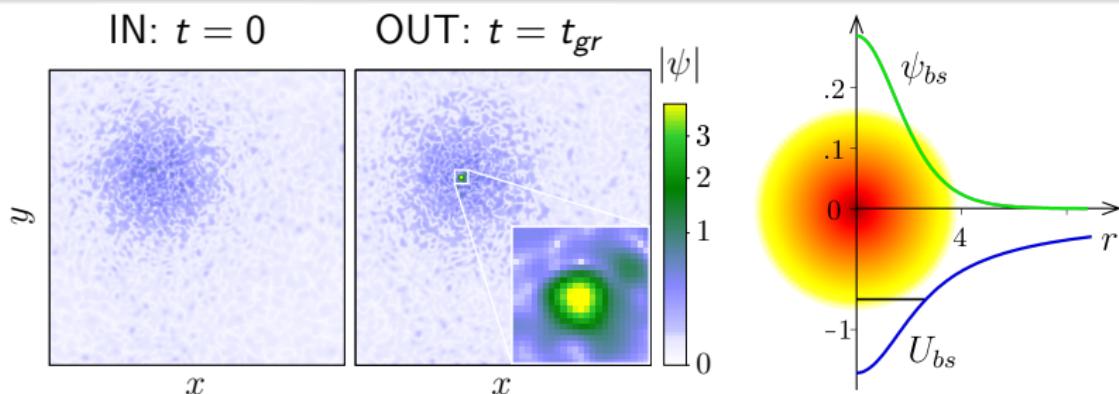
starting from random (virialized) waves

$$t = 9.905 \times 10^5 / mv_0^2$$

arXiv:1804.05857



This is a Bose star



- **Bose star** = Bose-condensate on a single level of U_{bs}
- DM is light \Rightarrow The Universe is packed with Bose stars!

How do the Bose stars grow?

- A difficult problem
- Numerical simulations: conflicting results

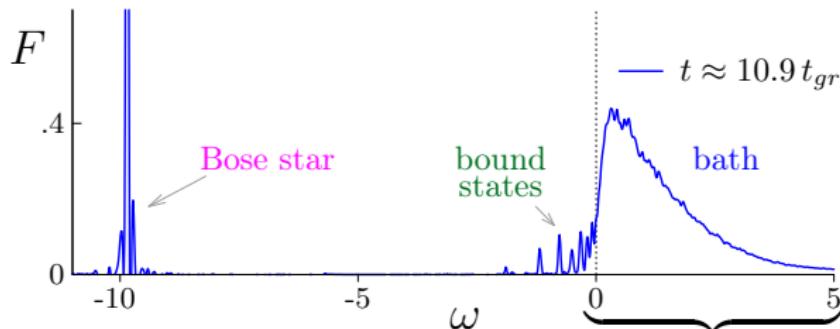
$$M_{bs} \propto t^{1/2}, t^{1/4}, t^{1/8}$$

e.g. Levkov et al '18, Eggemeier et al '19, Chan et al '22

But the solution is simple and analytical!

Distribution of particles over energies

$$F(t, \omega) \equiv \frac{1}{N} \frac{dN}{d\omega} = \int \frac{dt_1 d^3x}{2\pi N} \psi(t, x) \psi^*(t + t_1, x) e^{i\omega t_1 - t_1^2/\Delta t^2}$$

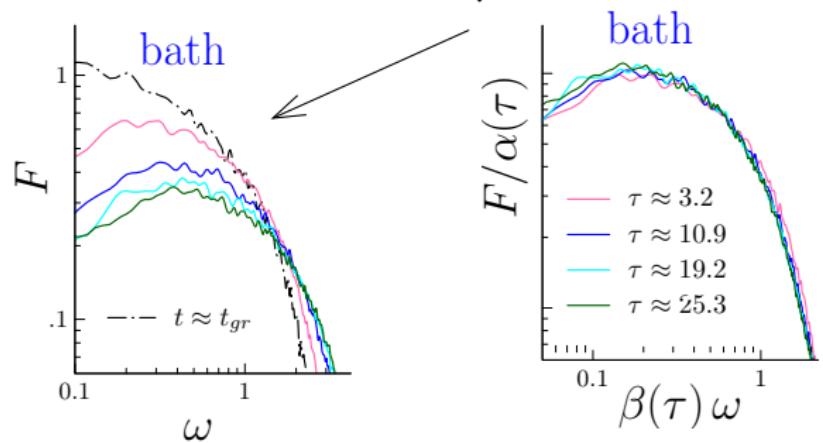


Scaling symmetry

$$F = \alpha F_s(\beta \omega)$$

$\alpha = \tau^{-1/D}$	$\beta = \tau^{2/D-1}$
------------------------	------------------------

$$\tau \equiv t/t_{gr}, \quad D = 2.8$$



Self-similar bath

- Consider the bath $\omega > 0$
 - Ignore the Bose star potential $U_{bs}(r)$
- $\} \Rightarrow \boxed{\partial_t F = St F}$
bath kinetic eq.

Ansatz passes the equation!

$$F(\tau, \omega) = \underbrace{\alpha(\tau)}_{\text{self-similar profile}} \underbrace{F_s}_{(\beta(\tau)\omega)} \quad \begin{aligned} \alpha &= \tau^{-1/D} \\ \beta &= \tau^{2/D-1} \end{aligned}$$

⇒ Profile Eq:

$$(2/D - 1)\omega_s \partial_{\omega_s} F_s - F_s/D = St F_s$$

⇒ Power-law bath mass & energy

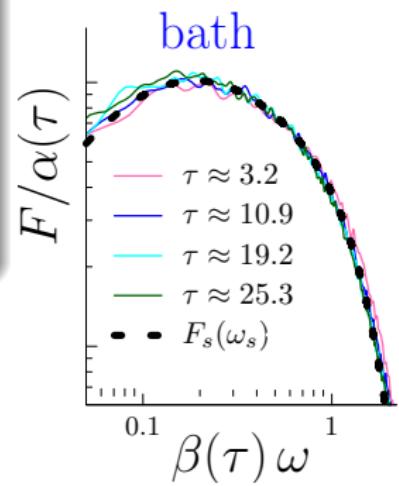
$$M_b \propto \tau^{1-3/D}, \quad E_b \propto \tau^{2-5/D}$$

non-autonomous system!

- Nontrivial BC: Particle Flux $\neq 0$ at $\omega = 0$

Then the solution exists
And it is a kinetic attractor

$$\tau \equiv t/t_{gr}$$



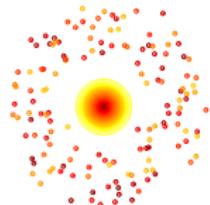
Bose star growth

Assumptions:

- Quasi-stationary & self-similar bath:

$$D = D(t) \quad \text{but still}$$

$$E_b^3 / M_b^5 \propto \tau$$



- Energy & mass conservation:

$$E_b = E - E_{bs} - E_e,$$

Bose star
↑
ex. states

$$M_b = M - M_{bs} - M_e$$

Bose star
↑
ex. states

- Bose star energy:

$$E_{bs} = -\gamma M_{bs}^3, \quad \gamma \approx 0.0542 m^2 G^2$$

constant

- Low occupancies of bound states:

$x_e \equiv M_e / M \approx \text{const}$, $E_e \approx 0$ — confirmed by simulations

$$\frac{(1 + x^3/\epsilon^2)^3}{(1 - x_e - x)^5} = \frac{\tau - \tau_i}{\tau_*}$$

- $x(\tau) \equiv M_{bs}/M$
- $\tau \equiv t/t_{gr}$

A simple and predictive law!

Compare with simulations

Equation

$$\frac{(1 + x^3/\epsilon^2)^3}{(1 - x_e - x)^5} = \frac{\tau - \tau_i}{\tau_*}$$

Variables

- $x(\tau) \equiv M_{bs}/M$
- $\tau \equiv t/t_{gr}$

Parameters

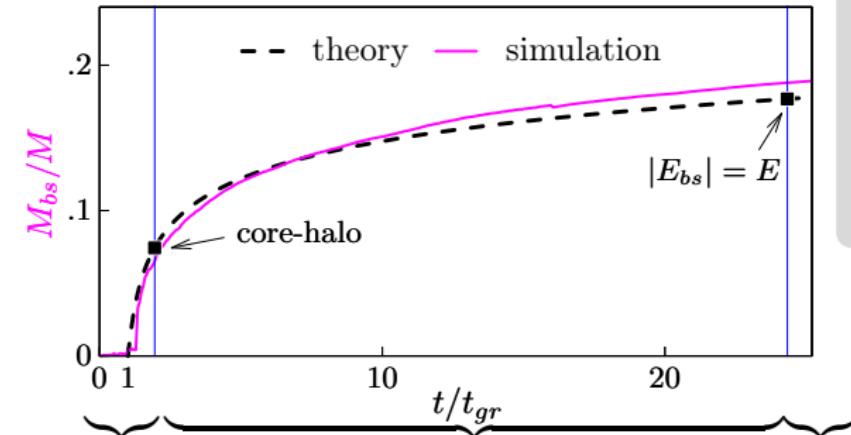
To fit:

- $x_e \equiv M_e/M$ — small
- $\tau_i \approx -0.1$ — universal

Known:

- $\epsilon^2 \equiv E/\gamma M^3$
- $\tau_* = \frac{1-\tau_i}{(1-x_e)^5} \leftrightarrow M_{bs}(1) = 0$

$$\leftarrow \epsilon = 0.074, x_e = 0.043, \tau_i = -0.1$$



$M_{bs} \propto t$ $t^{1/3}$ $t^{1/9}$ — like in simulations!

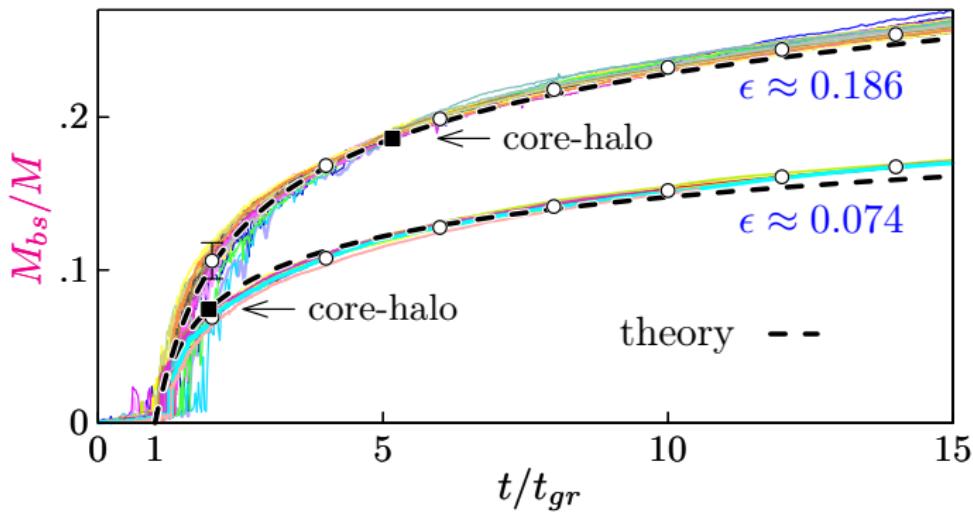
cf. Levkov et al '18, Eggemeier et al '19, Chan et al '22

core-halo: $|E_{bs}/M_{bs}| = E/M$ — stars in simulations “stop growing”

Schive et al '14

Statistical test

- 33 simulations
- Essentially different t_{gr} & two values of ϵ



Other performed tests

- + Nonzero self-interaction $\lambda_4 \neq 0$
- + Condensation in miniclusters

Applications to cosmology: String axions

$$\frac{(1 + \textcolor{magenta}{x}^3/\textcolor{blue}{\epsilon}^2)^3}{(1 - \textcolor{magenta}{x})^5} \sim \frac{t}{t_{gr}}$$

$$x(t) \equiv \frac{M_{bs}}{M}$$

Core of Fornax dwarf



$$M \sim 10^8 M_\odot$$

$$v_{\text{vir}} \sim 20 \text{ km/s}$$

- Parameters:

$$t_{gr} \sim 0.05 \frac{m^3 (GM)^4}{\Lambda v_{\text{vir}}^6}$$

$$\epsilon \sim 3 \frac{v_{\text{vir}}}{G m M}$$

- Time to “core-halo” slowdown:

$$t_{\text{c/h}} \sim 9\epsilon t_{gr} \sim \Lambda^{-1} m^2 (GM)^3 / v_{\text{vir}}^5$$

- Fuzzy DM in Fornax Dwarf ($m \sim 10^{-22} \text{ eV}$)

$$t_{\text{c/h}} \sim 10^7 \text{ yr} - \text{form \& grow}$$

- Experimental bound: $m \gtrsim 2 \cdot 10^{-20} \text{ eV}$

$$t_{\text{c/h}} \gtrsim 10^{11} \text{ yr} - \text{do not grow! (in Fornax)}$$

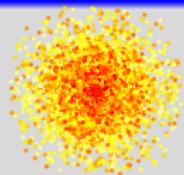
Larger galaxies are worse!

Only small Bose stars exist in the Universe!

$$\frac{(1 + \textcolor{magenta}{x}^3/\epsilon^2)^3}{(1 - \textcolor{magenta}{x})^5} \sim \frac{t}{t_{gr}}$$

$$x(t) \equiv \frac{M_{bs}}{M}$$

Axion minicluster



$$M \sim 10^{-17 \div 12} M_\odot$$

$$\Phi = \delta \rho_a / \bar{\rho}_a \Big|_{RD}$$

$$= 0 \div 10^3$$

Hogan, Rees '88; Kolb, Tkachev '93

- Parameters:

$$t_{gr} \sim \frac{5 \cdot 10^8 \text{ yr}}{\Phi^4} \left(\frac{M}{10^{-14} M_\odot} \right)^2 \left(\frac{m}{10^{-4} \text{ eV}} \right)^3$$

$$\epsilon \sim 0.02 \Phi^{2/3} \left(\frac{M}{10^{-14} M_\odot} \right)^{-2/3} \left(\frac{m}{10^{-4} \text{ eV}} \right)^{-1}$$

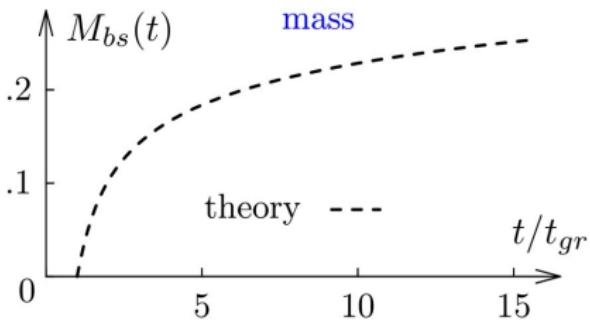
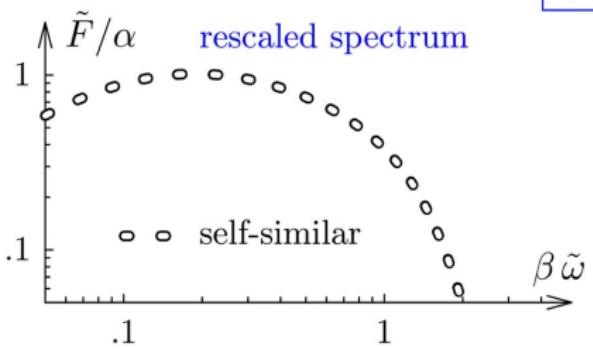
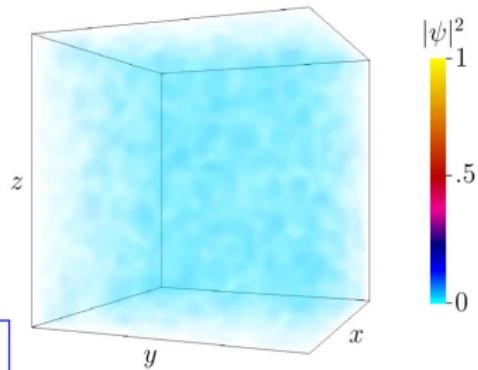
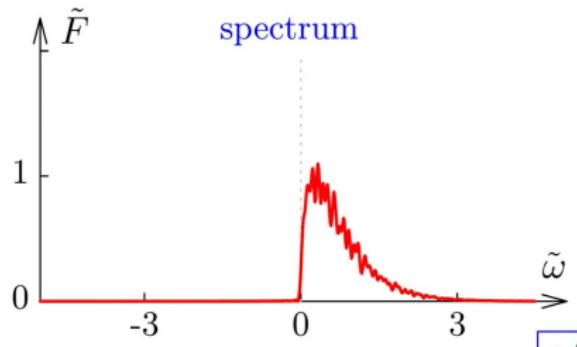
- Time to eat 10% of minicluster:

$$t_{10} \sim t_{gr} \frac{(10\%)^9}{\epsilon^6} < 10^{10} \text{ yr} \quad \text{if} \quad \boxed{\Phi \gtrsim 1}$$

All denser miniclusters turn into axion stars

The Universe is packed with grown-up axion stars!

Conclusions I: the movie



Conclusions II: Implications of Bose stars in axion cosmology

- Less diffuse DM \Rightarrow weaker signals in DM detectors
- Gravitational microlensing and femtolensing

Kolb, Tkachev '96; Fairbairn et al '17

- Radio lines from transient axion stars

Witte et al '22

- Parametric resonance: radio explosions of heavy stars — explain FRB?

Levkov, Panin, Tkachev '20; Chung-Jukko et al '22

- Radio-emitting stars reionize the cosmological medium

Escudero et al '23

- Bosenovas: heavy stars collapse & emit relativistic axions

Levkov, Panin, Tkachev '17; Eby et al '22

THANK YOU FOR ATTENTION!

This work was supported by the RSF grant 22-12-00215