CONSTRAINTS ON MODIFIED THEORIES OF GRAVITY FROM THE LATEST LIGO-VIRGO RINGDOWN OBSERVATIONS





Phys. Rev. D 103, 124043 (2021)

Gregorio Carullo Lomonosov Conference August 2021



UNIVERSITÀ DI PISA



- Current status of **Observational black hole spectroscopy**;
- The "So what?" test
- Beyond the LVC parametrisation: **ParSpec**
- ParSpec constraints on **modified gravity** from black holes spectra
- Comparison to other experimental bounds
- A non-perturbative beyond-Kerr case: Kerr-Newman black holes

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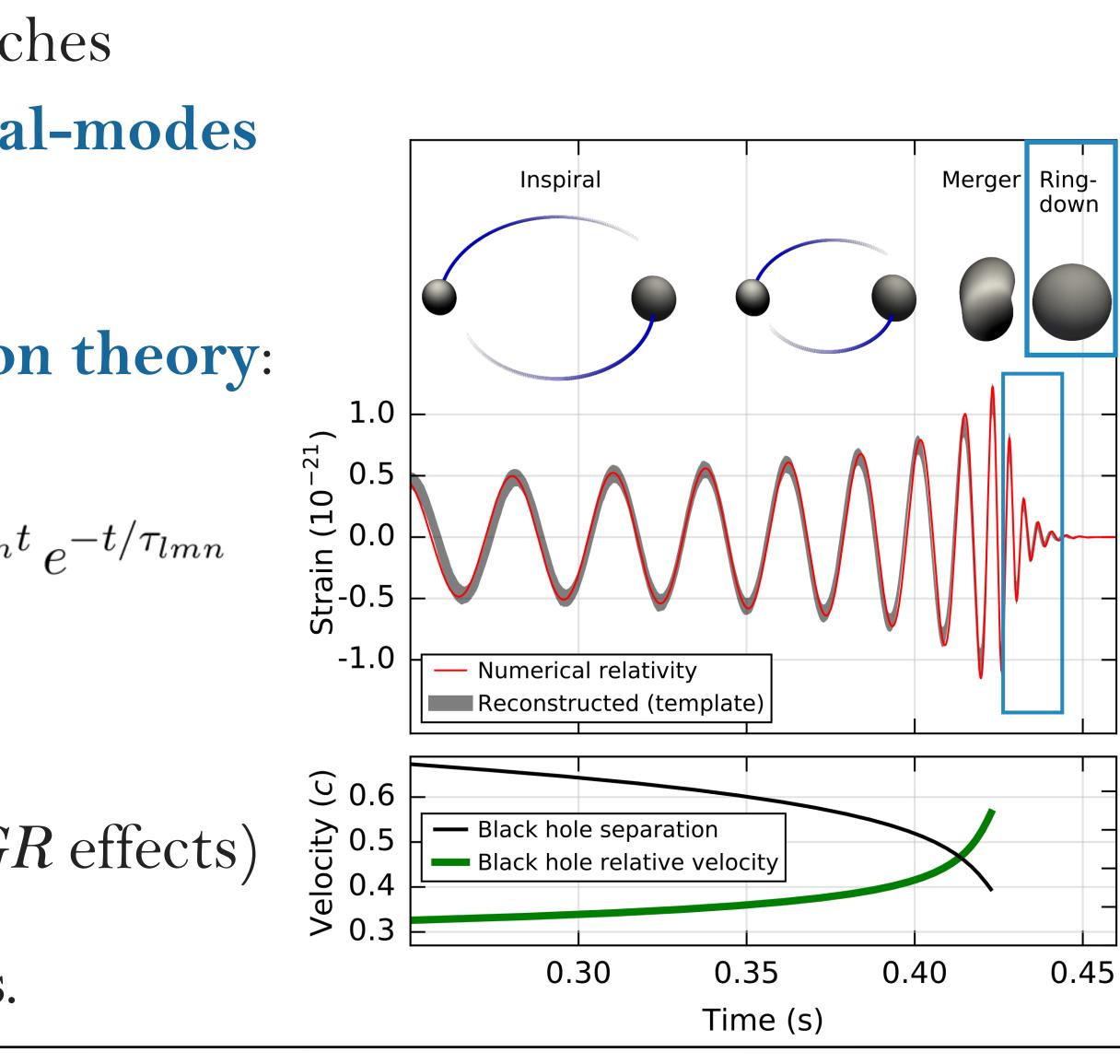
CONTENTS

- Binary black holes **remnant** approaches equilibrium emitting damped normal-modes ("spacetime signature").
- **Spectrum** predicted by **perturbation theory**:

$$h_{+} - i h_{\times} = -\frac{M}{r} \sum_{l,m,n} \mathcal{A}_{lmn} S(\theta, \phi) e^{i\omega_{lmn}}$$

- "Simple" prediction ("easy" to incorporate *beyond-GR* effects)
- Sensitive to ~ *light ring* physics.

BLACK HOLE RINGDOWN



LVC (2016) PRL 116, 061102





- **Testing** the emission:
 - Remnant compact object **nature**; Are we really observing black holes?
 - General Relativity predictions for **spectral emission**; Is General Relativity a correct description of gravity at high curvatures?
 - Black Hole Uniqueness Theorems; Do non-extremal black holes have additional hairs?
 - Is our classical description of black holes valid?

• Possible quantum horizon effects and classical **BH thermodynamics**.

MISSING PHYSICS?

- Why GW merger-ringdown observations compared to other experiments? (Solar System, Binary Pulsars, X-ray binaries)
 - Largest curvature regime experimentally accessible;
 - Gauge the contribution of missing physics for precision tests:
 - Cosmological expansion
 - Darz matter
 - Accretion, magnetic fields, electric charge
- GW ringdown tests are **clean**!

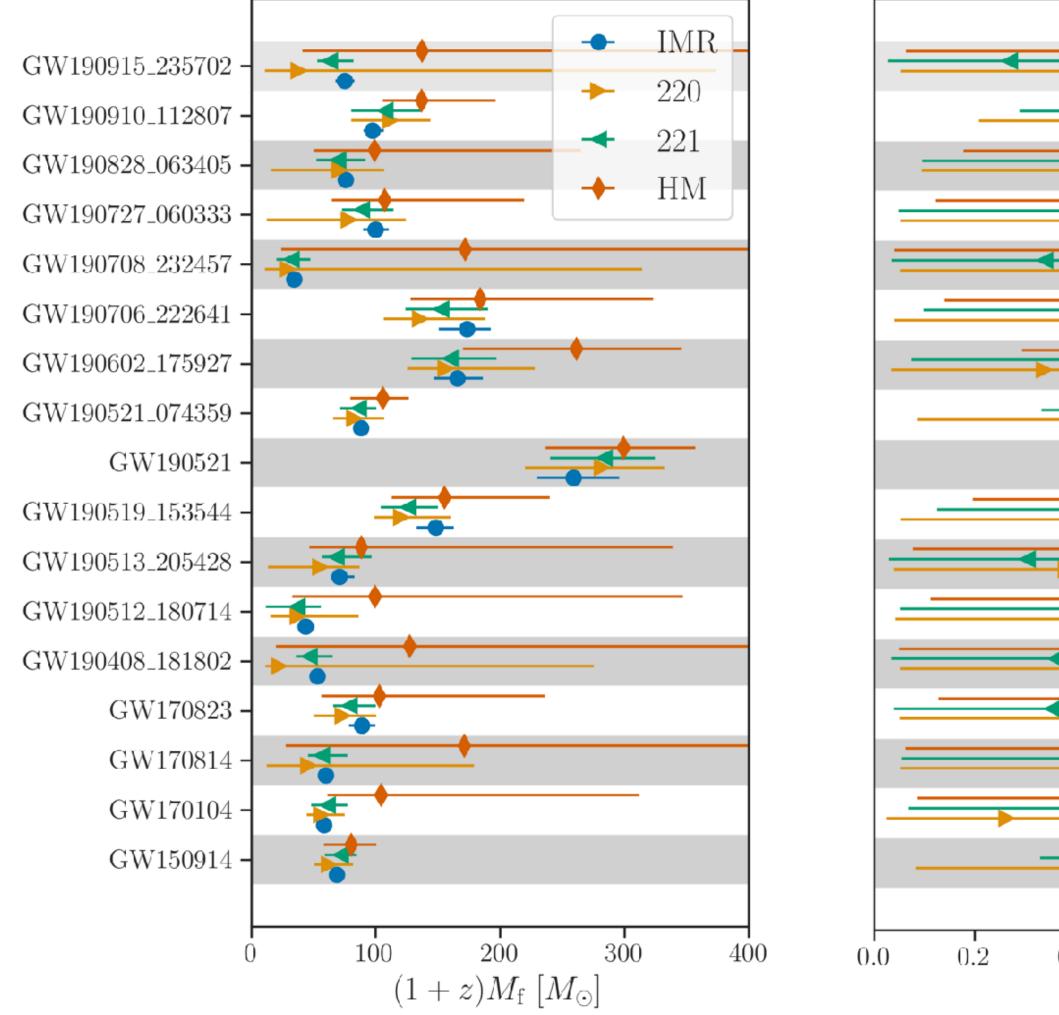
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BLACK HOLE RINGDOWN CATALOG

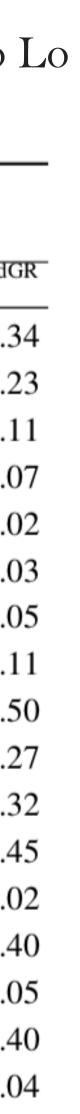
• First catalog of ringdown-only observations:



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Image credit: LIGO-Virgo/Rico Lo

	- GW190915_235702	Higher modes	Overtones	
	$-GW190910_{-112807}$	$\log_{10}\mathcal{B}_{220}^{\mathrm{HM}}$	$\log_{10} \mathcal{B}^{221}_{220}$	$\log_{10} O_{\rm GR}^{\rm modG}$
	- GW190828_063405	0.03	0.63	-0.3
	-GW190727_060333	0.26	-0.20	-0.2
	-GW190708_232457	0.04	-0.19	-0.1
	-GW190706_222641	0.02	-0.98	-0.0
	$-GW190602_175927$	-0.05	-1.02	-0.0
	-GW190521_074359	0.09	-0.42	0.0
	- GW190521	0.09	-0.54	-0.0
	- GW190519_153544	0.21	-0.00	-0.1
	- GW190513_205428	0.12	-0.86	-0.5
	-GW190512_180714	-0.04	1.29	-0.2
	-GW190408_181802	0.61	-1.56	0.3
	- GW170823	-0.06	-0.64	-0.4
	-GW170814	-0.11	-0.17	-0.0
	-GW170104	-0.02	-1.65	-0.4
	-GW150914	0.05	-0.72	-0.0
	- G W 100914	-0.10	-0.64	-0.4
0.4 0.6 0.8 1.	0	0.06	-0.37	-0.0
χ_{f}	0			



BLACK HOLE RINGDOWN CATALOG

• First catalog of ringdown-only observations:



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- GW190915_235702	Higher modes	Overtones	
- GW190910_112807	$\log_{10} \mathcal{B}_{220}^{\text{HM}}$	$\log_{10} \mathcal{B}_{220}^{221}$	$\log_{10} O_{\rm GR}^{\rm modGR}$
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= $=$ $=$ $=$ $=$ $=$ $=$ $=$ $=$ $=$		-0.98	-0.07
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	0.	-0.00	-0.11
GV 05 96 8	0.	-0.86	-0.50
90512 0714	0.01	1.29	-0.27
-GW190408_181802	0.61	-1.56	0.32
- GW170823	-0.06	-0.64	-0.45
-GW170814	-0.11	-0.17	-0.02
-GW170104	-0.02	-1.65	-0.40
-GW150914	0.05	-0.72	-0.05
	-0.10	-0.64	-0.40
0.4 0.6 0.8 1.0	0.06	-0.37	-0.04
$\chi_{ m f}$			

GWTC-2 Testing GR, LVC (2020), 2010.14529

- **Bounds on deviations** from the GR spectrum. Data consistent with GR.
- **Deviations** parameterized as:

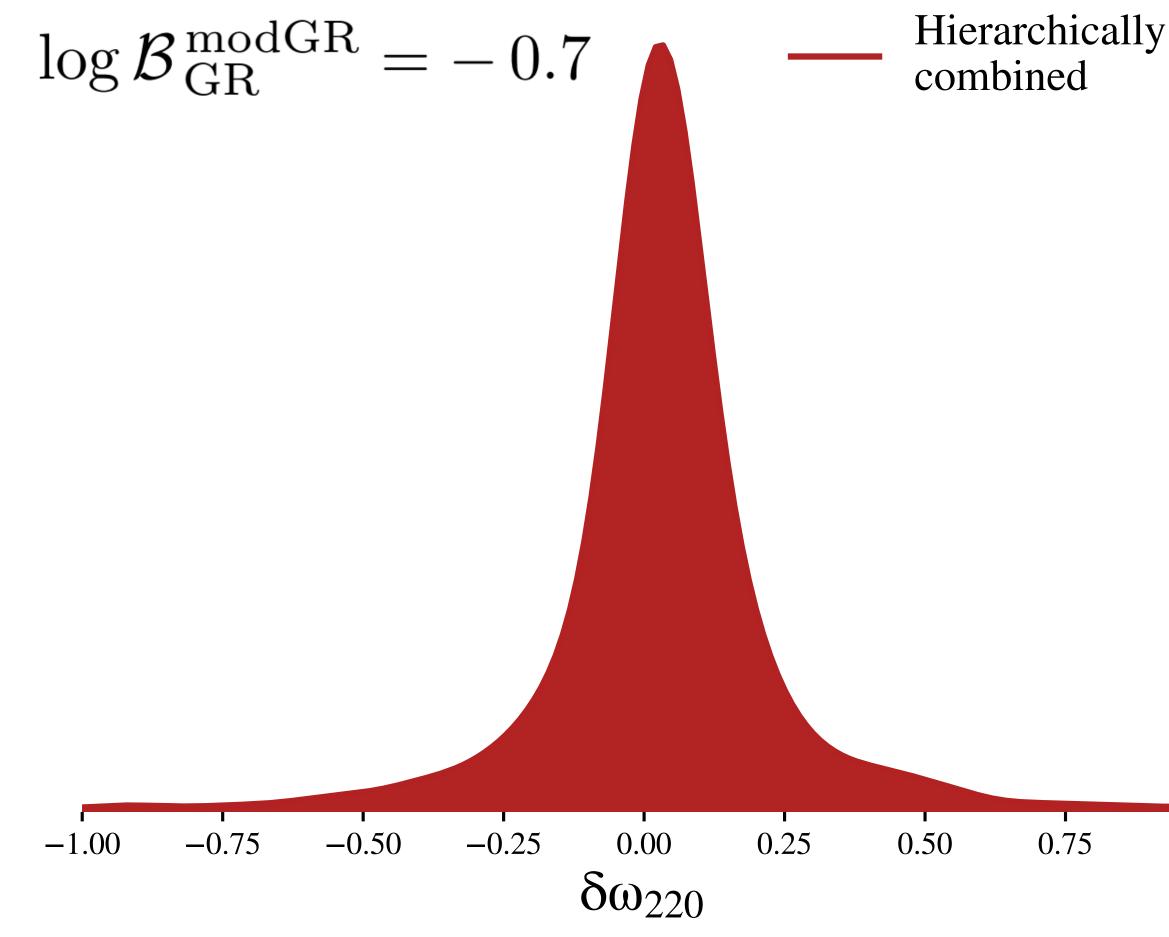
$$\omega = \omega^{Kerr} \cdot (1 + \delta\omega)$$
$$\tau = \tau^{Kerr} \cdot (1 + \delta\tau).$$

$$\delta\omega_{220} = 0.03^{+0.38}_{-0.35}$$

See also: Ghosh+, arXiv:2104.01906

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TESTS OF GENERAL RELATIVITY



GWTC-2 Testing GR, LVC (2020), 2010.14529





THE "SO WHAT?" TEST

- Do these results pass the "So what?" test?
 - Take an urn full of a gaussian mixture of theoretical particle physicists and general relativists, with a possibly non-null intersection.
 - Draw them at random and have **them look** separately at **these results**.
 - **Count** how many people say: "So what?"
 - If it's more than 42% of them, the test is **failed**.

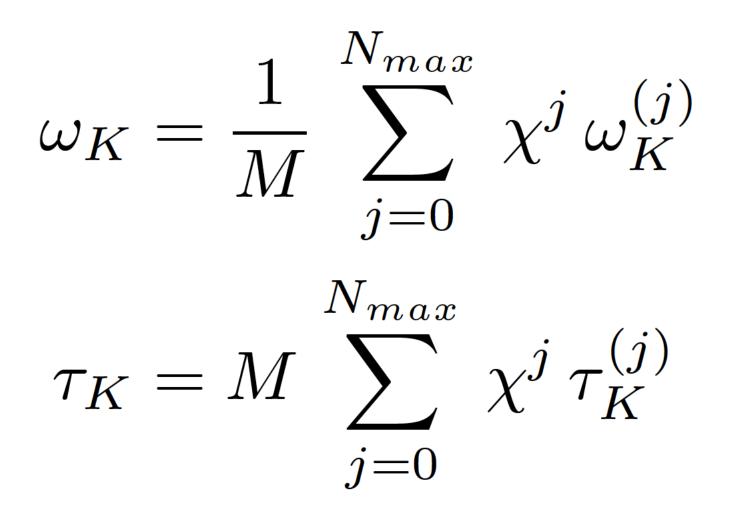




- Implications of LVK results to specific alternative theories of gravity not immediate (unlike, e.g. graviton mass);
- Why doesn't the LVK constrain specific and well-motivated beyond-GR theories?
 - Most theories are **not proven to be well-posed**. Difficult to rigorously justify certain approximations or understand merger-ringdown in a specified theory.
 - Huge amount of possibilities and of effects to take into account (isospectrality breaking, parity violations, additional modes, scalar-fields dynamics...)
- **Policy** of the Testing GR group: **do not consider** specific theories for which a **rigorous** framework has not been established.







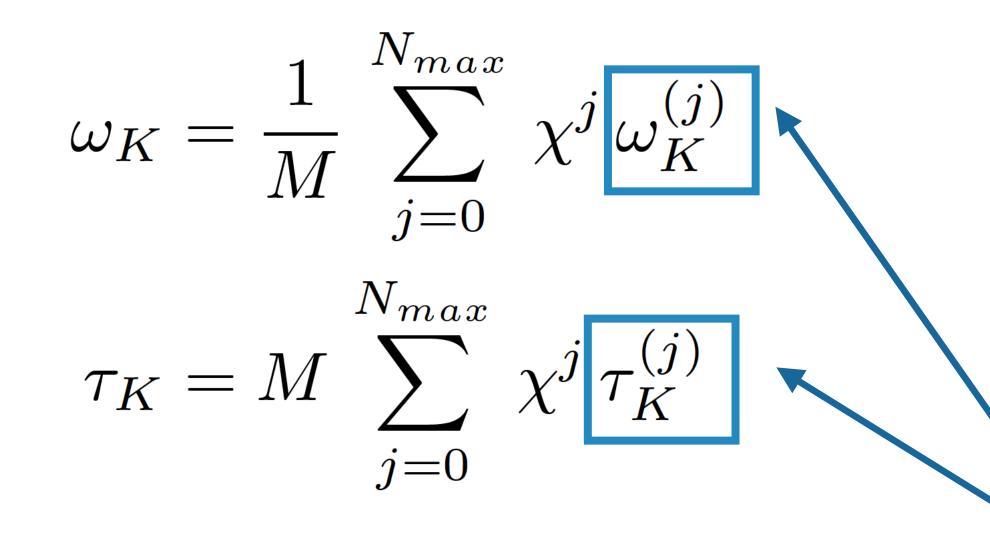
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Expand each QNM parameter in a polynomial expansion in the remnant spin.

> Extract the mass and spin structure in polynomial form.





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Expand each QNM parameter in a polynomial expansion in the remnant spin.

Extract the mass and spin structure in polynomial form.





$$\omega_K = \frac{1}{M} \sum_{j=0}^{N_{max}} \chi^j \,\omega_K^{(j)} \left(1 + \gamma \,\delta\right)$$
$$\tau_K = M \sum_{j=0}^{N_{max}} \chi^j \,\tau_K^{(j)} \left(1 + \gamma \,\delta\right)$$

Add deviations at each given order.

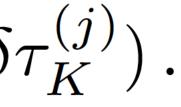
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 $\delta\omega_K^{(j)}$),

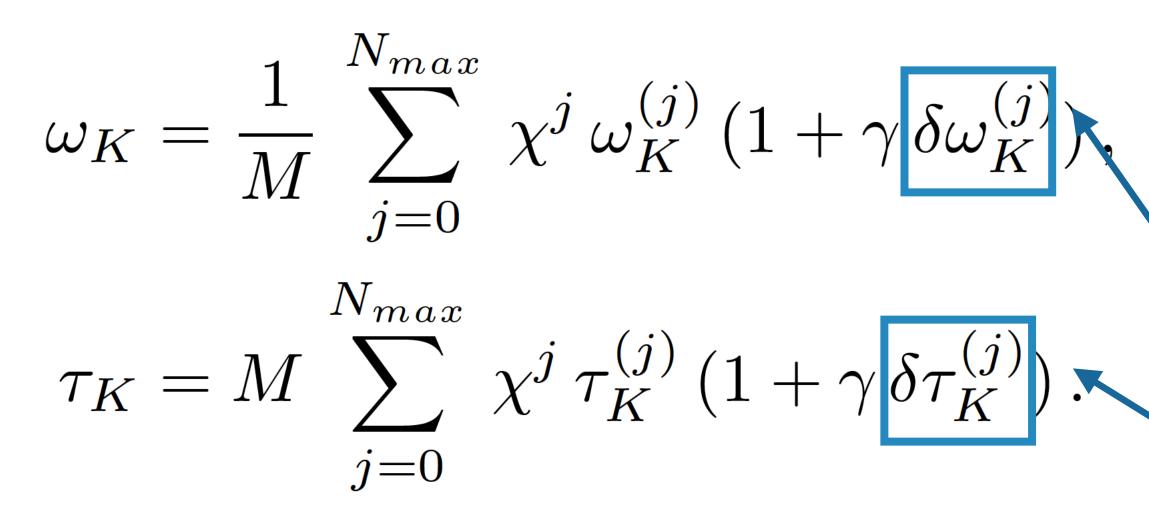
Proportional to action coupling(s):

$$\gamma := \left(\frac{\ell c^2 \left(1+z\right)}{G M}\right)^p$$



Maselli+, PRD 101, 024043 (2020)





Add deviations at each given order.

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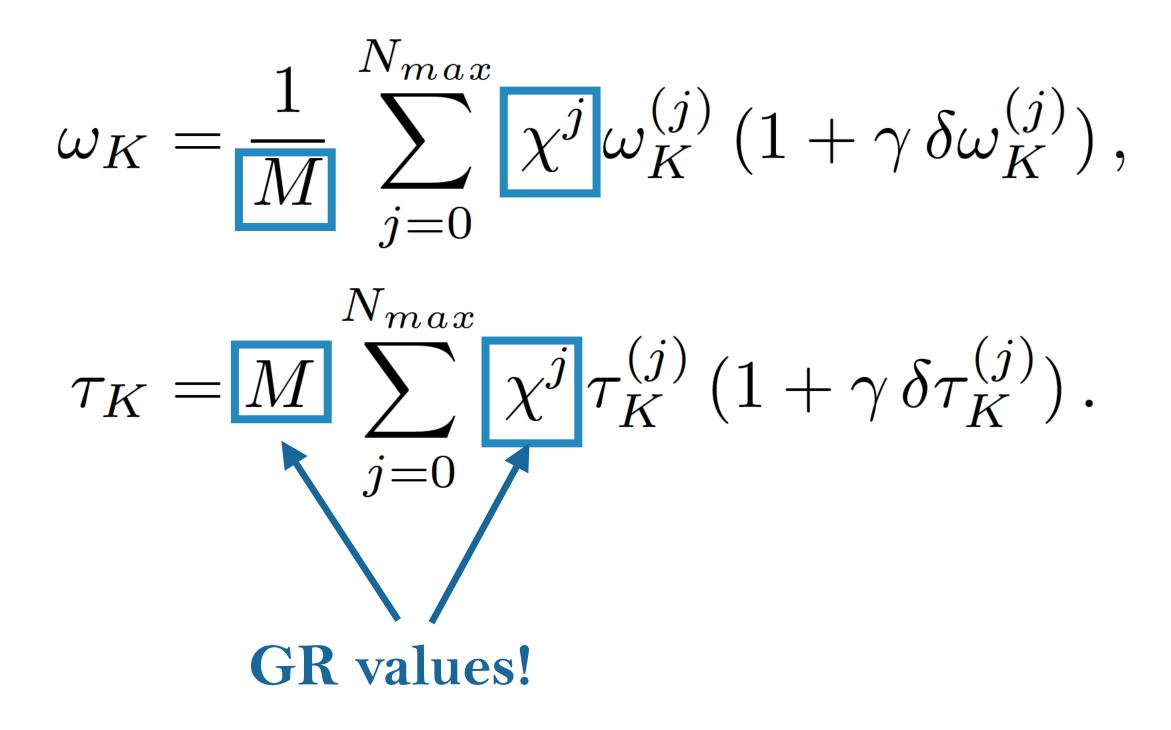


Proportional to action coupling(s):

$$\gamma := \left(\frac{\ell \, c^2 \, (1+z)}{G \, M}\right)^p$$

Also numerical constants! Independent of specific signal.





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Proportional to action coupling(s):

$$\gamma := \left(\frac{\ell c^2 \left(1+z\right)}{G M}\right)^p$$





- Advantages:
 - Applicable to large classes of **beyond-GR** theories;
 - **Source-independent** parameters: optimal constraints when combining events;

 - Direct connection with **parameters** appearing in the **action**.
- modifications from GR predictions.

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• Perturbative expansion: can retain **GR remnant mass and spin** parameters;

• Consequence: much **smaller number** of signals (or smaller SNR) to **detect**



THEORY PARAMETER SPACE

• **p=0** (e.g. certain scalar-tensor or Lorentz-violating)

$$S_{\text{E}} = \frac{1}{16\pi G_{\text{E}}} \int \sqrt{-g} \left(R - M^{\alpha\beta}{}_{\mu\nu} \nabla_{\alpha} u^{\mu} \nabla_{\beta} u^{\nu} \right) d^4x$$

• **p=2** (e.g. **Kerr-Newman** or Dark photon)

$$\mathcal{L} = \sqrt{-g} \left(\frac{R}{16\pi} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + 4\pi e j_{\rm em}^{\mu} A_{\mu} + 4\pi e_h j_h^{\mu} B_{\mu} + 4\pi \epsilon e j_h^{\mu} A_{\mu} \right)$$

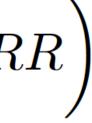
• p=4 (e.g. Einstein-scalar-Gauss-Bonnet or dynamical Chern-Simons)

$$S \equiv \int \frac{m_{\rm pl}^2}{2} d^4 x \sqrt{-g} \left[R - \frac{1}{2} (\partial \vartheta)^2 + 2\alpha_{\rm GB} f(\vartheta) \mathcal{R}_{\rm GB} \right], \quad S \equiv \int d^4 x \sqrt{-g} \left(\frac{m_{\rm pl}^2}{2} R - \frac{1}{2} (\partial \vartheta)^2 - \frac{m_{\rm pl}}{8} \ell^2 \vartheta^* R \right)$$

p=6 (e.g. Effective Field Theories)

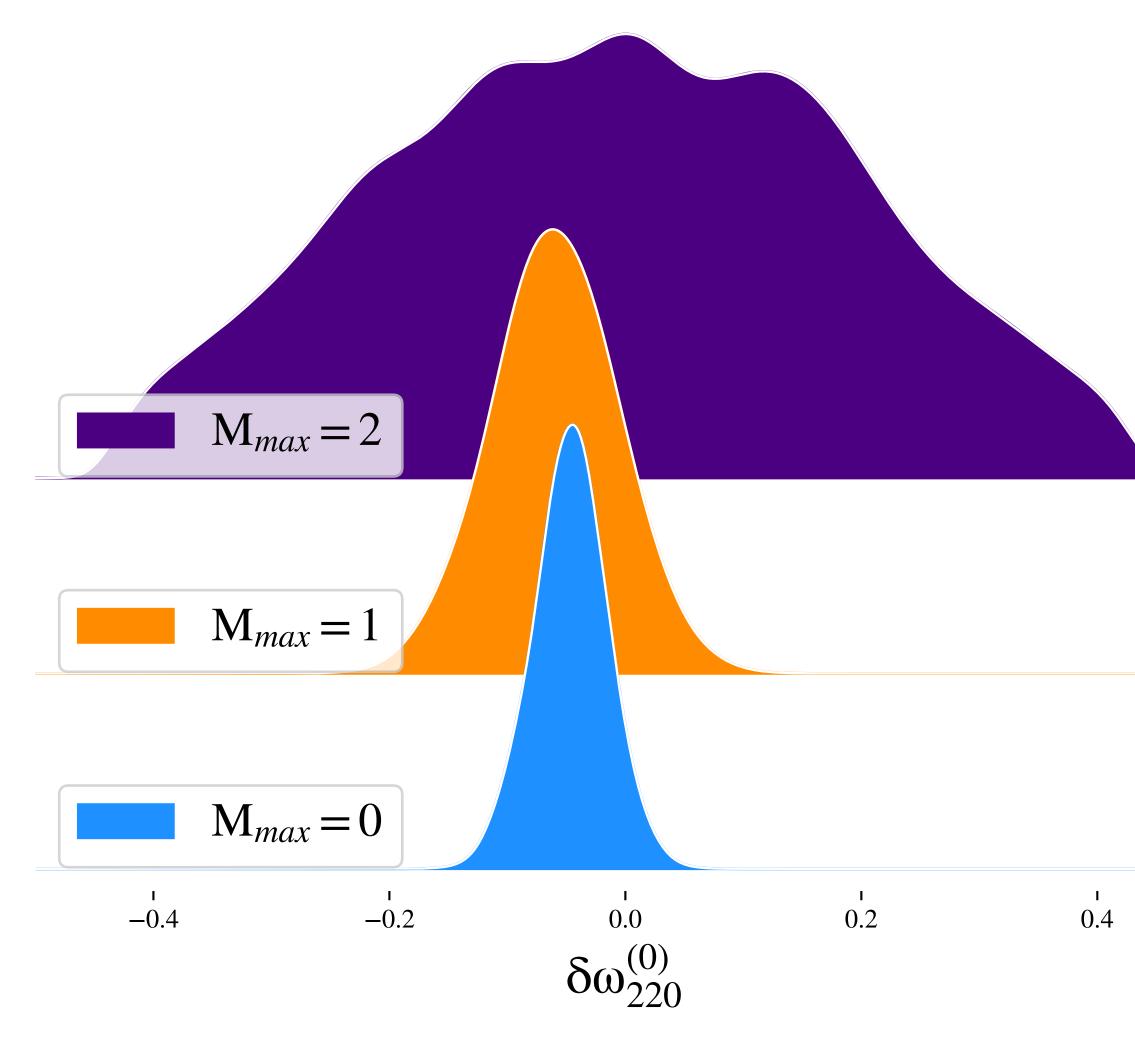
$$S_{\rm eff} = \int d^4 x \sqrt{-g} 2M_{\rm pl}^2 \left(R - \frac{C^2}{\Lambda^6} - \frac{\tilde{C}^2}{\tilde{\Lambda}^6} - \frac{\tilde{C}C}{\Lambda^6} \right)$$

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RESULTS FOR SCALAR DEVIATIONS

• Constraints on theories with scalar coupling in the action:



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$\delta\omega_{220}^0 = -0.05^{+0.05}_{-0.05}$

 $\log \mathcal{B}_{\rm GR}^{\rm modGR} = -14.55$

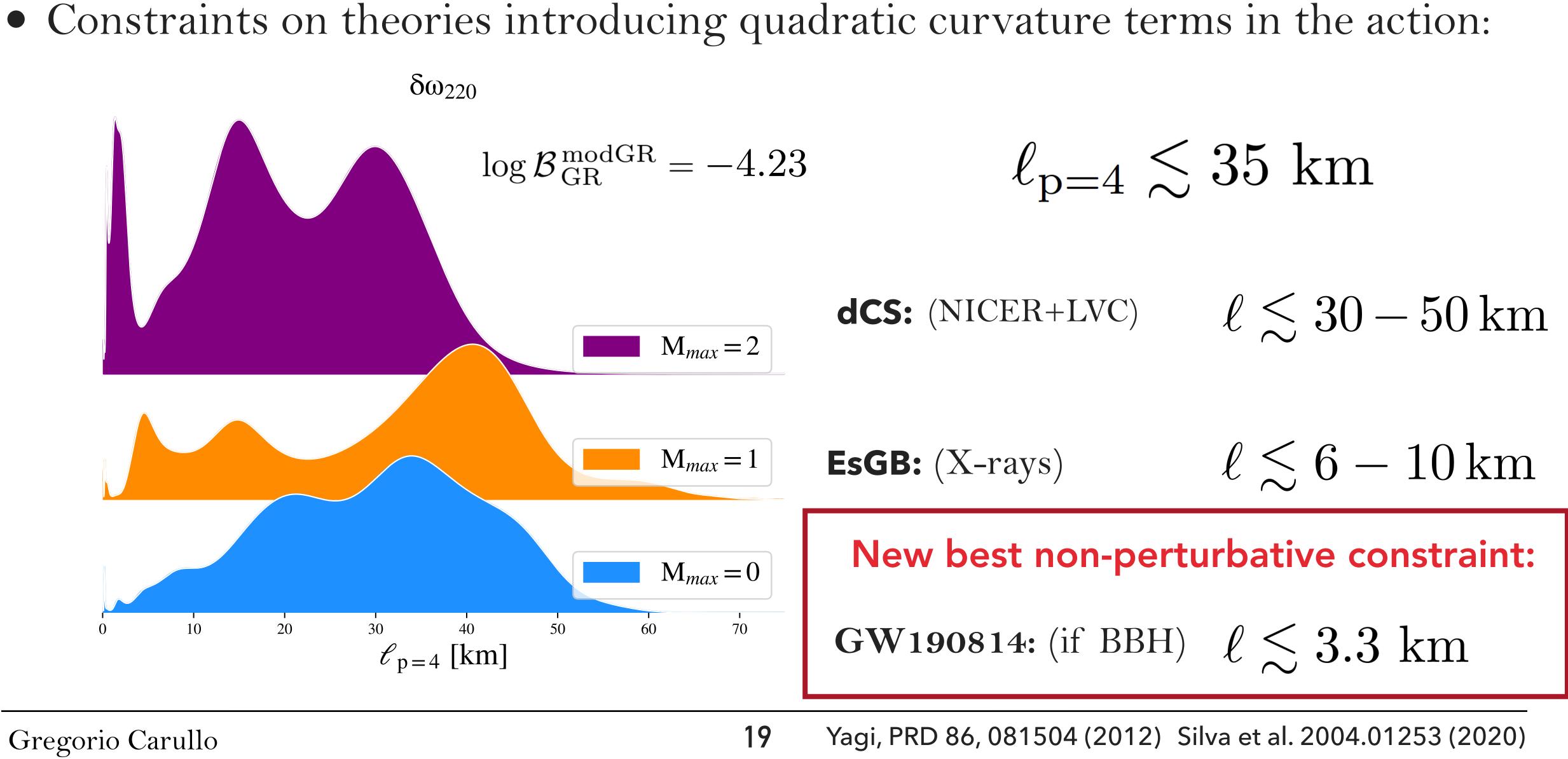
Improvement of a **factor of** ~2 wrt LVC parametrization.

Reduction factor of ~4 in N_{events} to detect a violation.

 M_{max} = spin-expansion order.



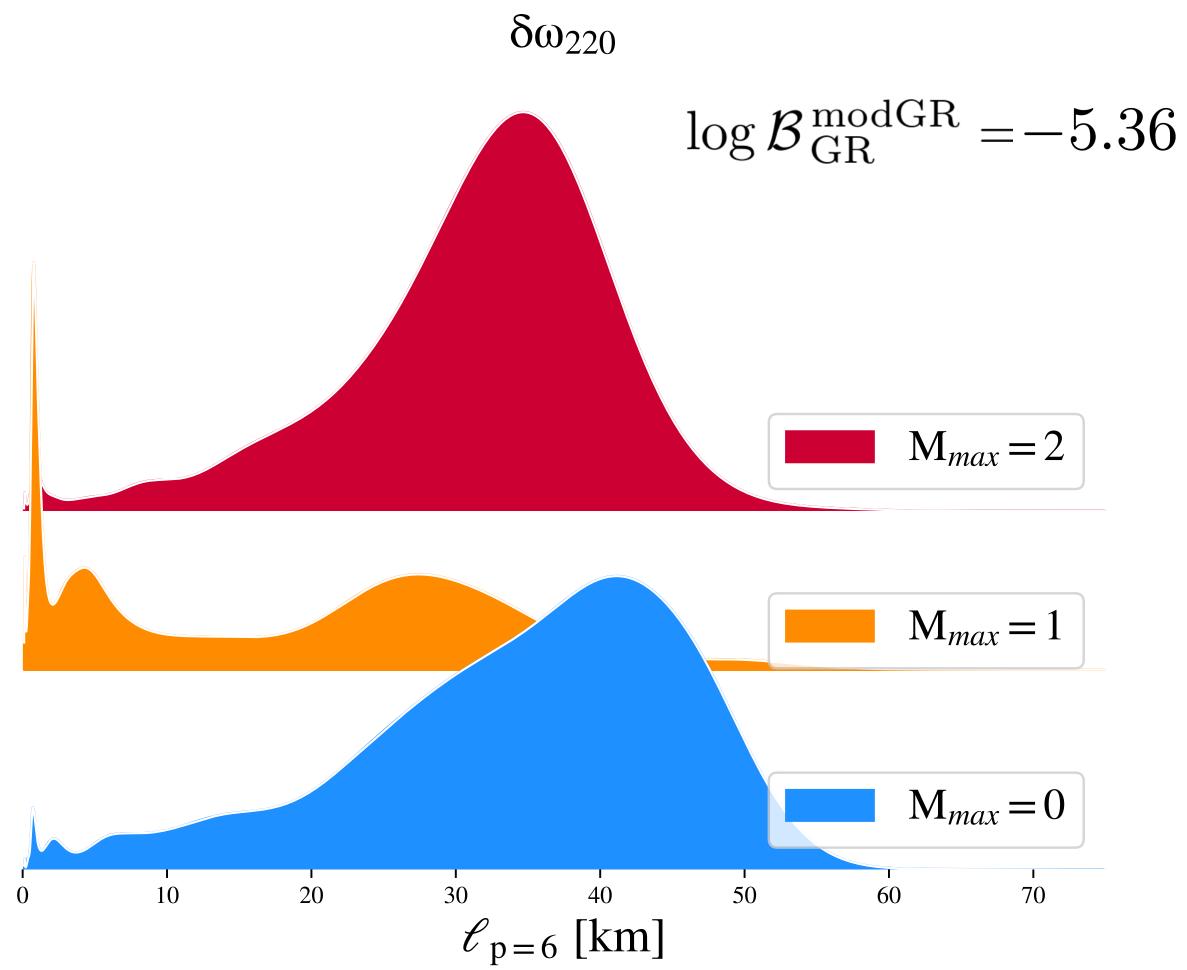
RESULTS FOR QUADRATIC GRAVITY



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RESULTS FOR EFFECTIVE FIELD THEORIES

• Constraints on viable Effective Field Theories of beyond-GR gravity:



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$\ell_{\rm p=6} \lesssim 42 \ \rm km$

Previous best bound from GW inspiral:

 $\ell \lesssim 150\,\mathrm{km}$

(Now probing finite-size effects)



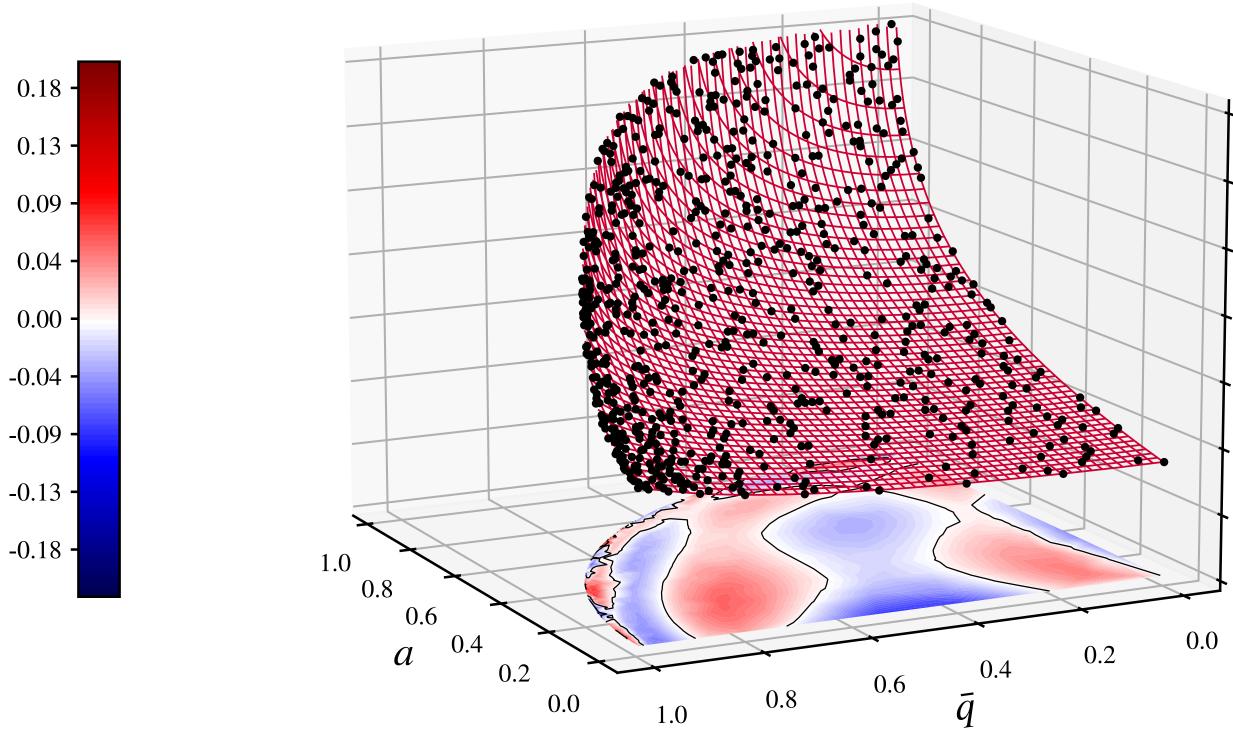


- Application to **charged BHs**: the **Kerr-Newman (KN)** case.
- beyond-GR effects.
- Non-standard scenarios: minicharged dark matter, magnetic monopoles, scalar-tensor-vector gravity, ...

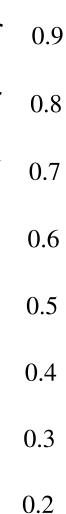
Carullo, Del Pozzo, Dias, Godazgar, Johnson-McDaniel, Laghi, Santos, Veitch (In prep.)

KERR-NEWMAN SPECTRUM

• Kerr spectrum deviations known for arbitrary spin: valuable **test-bed** for



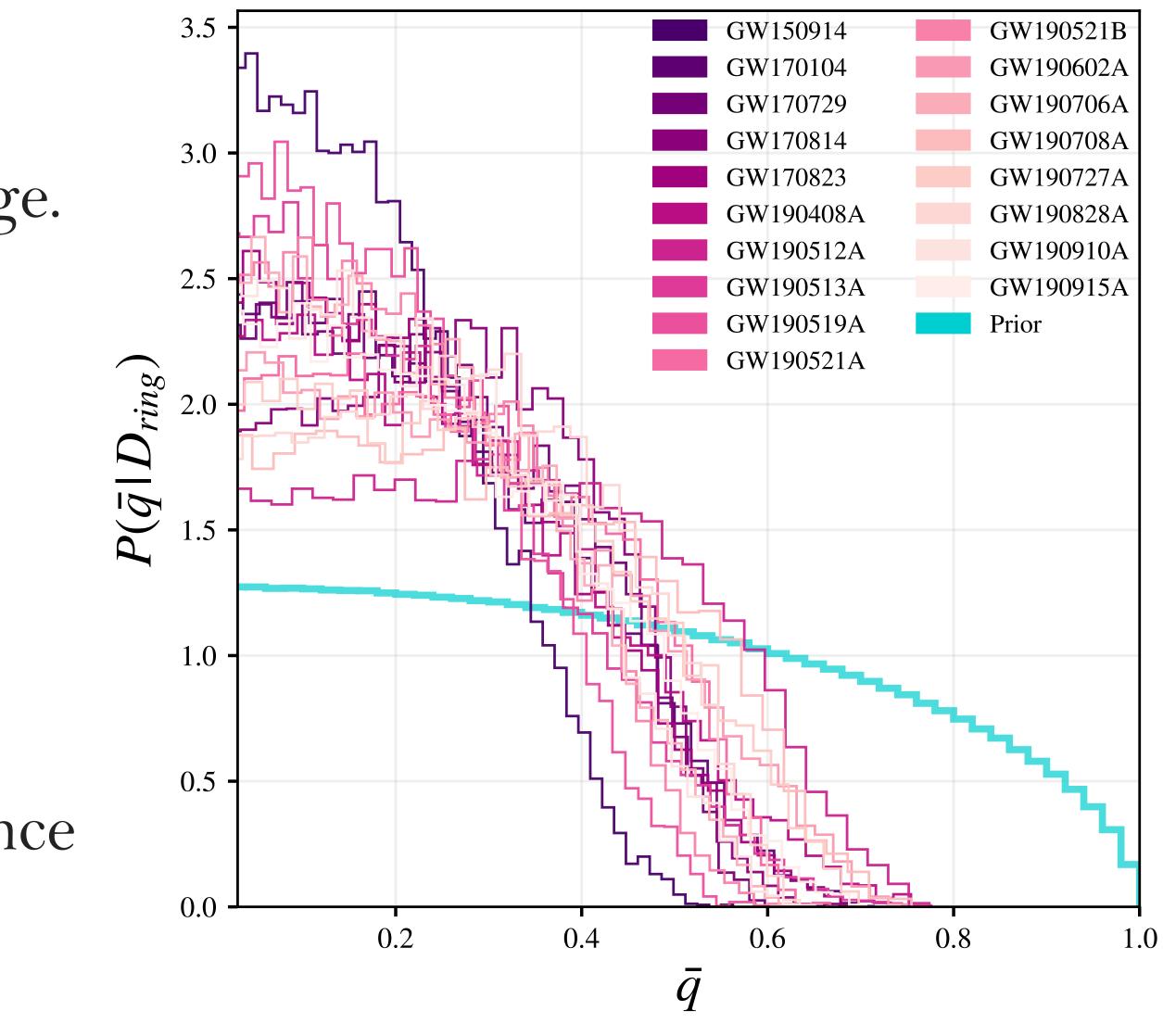
Residuals (%)



 ω_{220}

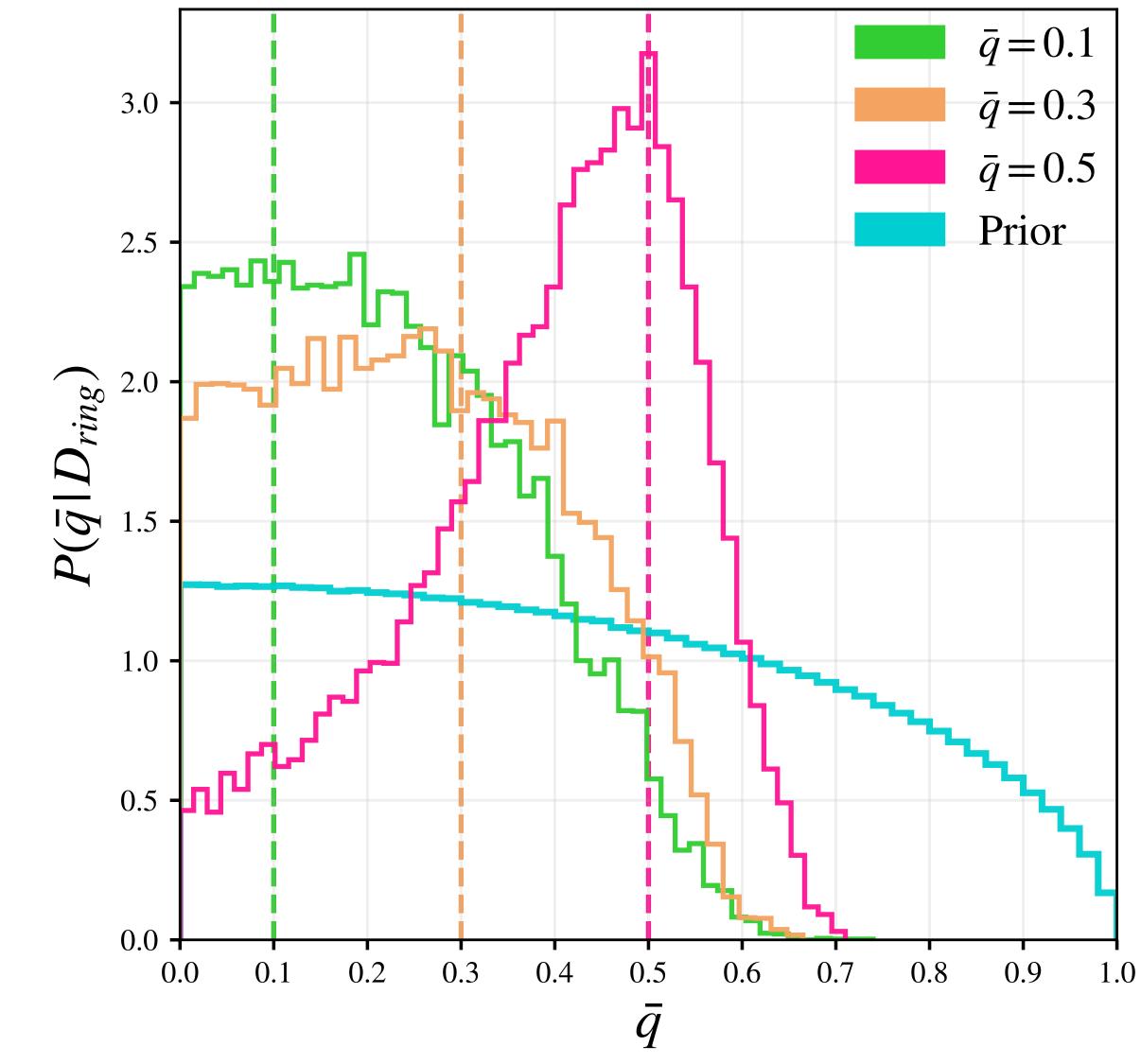
- "Standard" parameter estimation uninformative, due to strong **correlation** between spin and charge.
- Null test: maximum amount of charge compatible with current observations
- Best event (GW150914) gives: $\bar{q} < 0.33$
- No evidence for/against the presence of charge from Bayes' Factors

KERR-NEWMAN CONSTRAINTS



KERR-NEWMAN FUTURE CONSTRAINTS

- Can future observations from current detector network **discriminate** the presence of a **charge**?
- Simulate KN signals in LIGO-Virgo data at design sensitivity
- Recover using KN templates
- Charge confidently measured **only** for **high values**
- Need more info to break
 spin-charge correlations



- **Black hole perturbations** have unquestionably transitioned from a mathematical problem to an **observational reality**.
- our current gravity paradigm and search for signs of new physics.
- LIGO-Virgo ringdown observations can already probe high-curvature on beyond-GR extensions.

This material is based upon work supported by NSF's LIGO Laboratory which is a major facility fully funded by the National Science Foundation and by the European Gravitational Observatory, composed by CNRS, INFN and Nikhef.

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CONCLUSIONS AND PROSPECTS

• The analysis of gravitational black holes spectra is a powerful tool to test

unexplored regimes of beyond-GR theories or place competitive bounds

• Future detections promise to quickly overcome other experimental bounds.



Credits to: Jani, Ghonge

