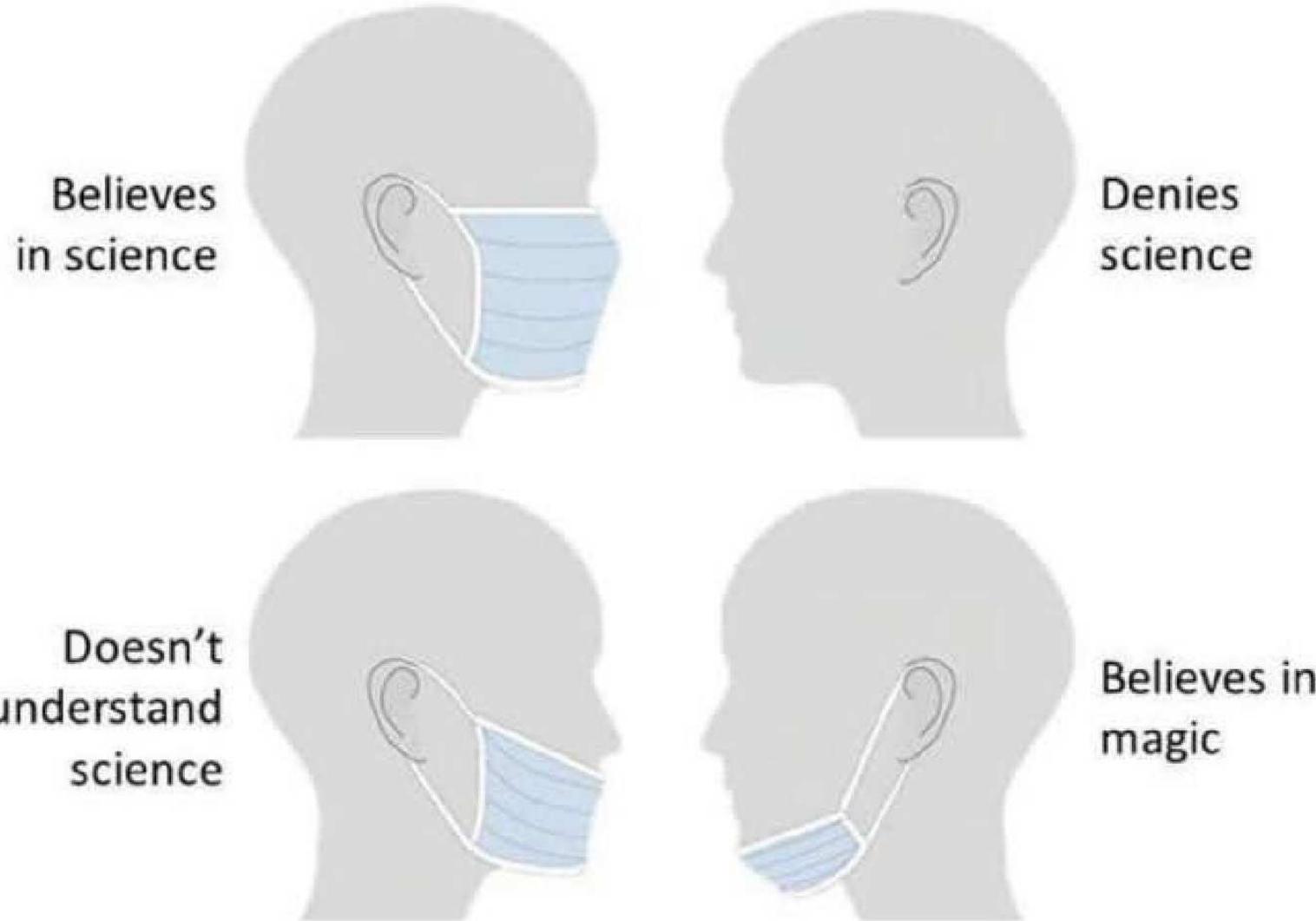


The Four COVID Personality Types

(spotted at the local grocery store)



BSM Physics in the light of the new “MUON G-2” Result

Sven Heinemeyer, IFT (CSIC, Madrid)

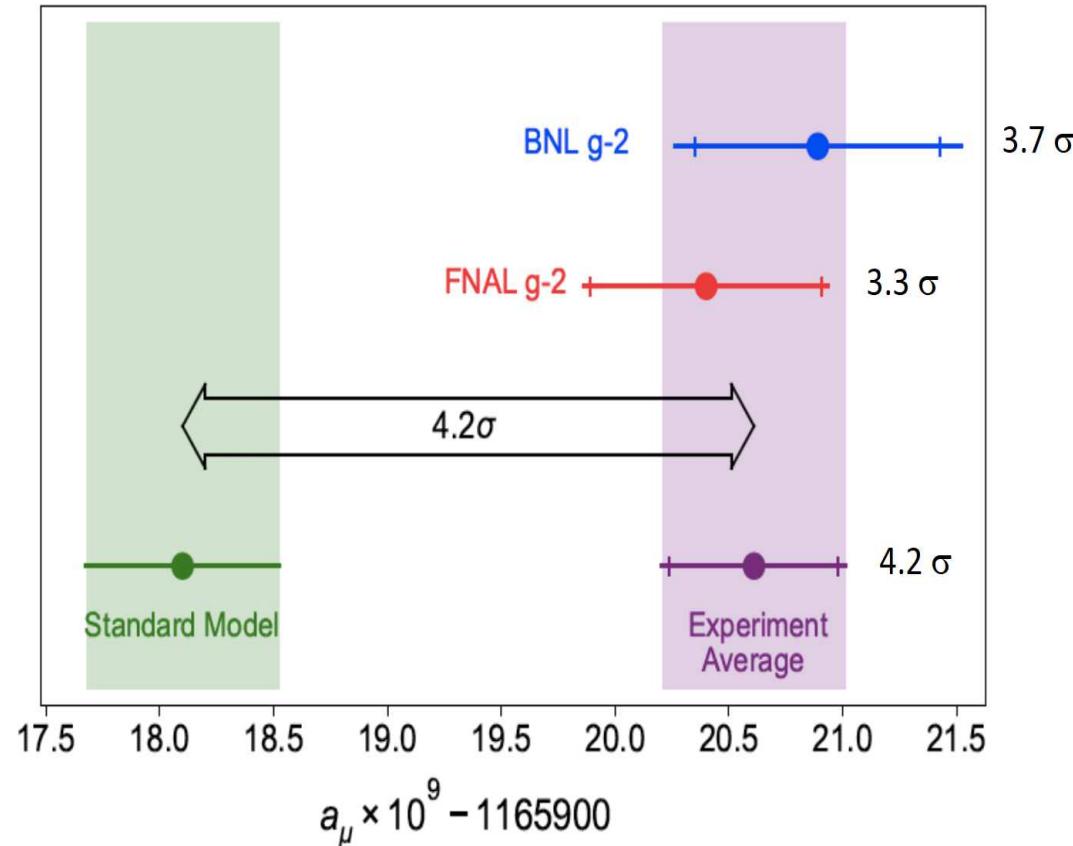
virtual, 07/2021

- 1. Introduction**
- 2. “Two Higgs Doublets are as natural as one”**
- 3. “SUSY fits like a glove”**
- 4. “There’s always the axion”**
- 5. Conclusions**

1. Introduction

The anomalous magnetic moment of the muon: $a_\mu \equiv (g - 2)_\mu / 2$

Overview about the current experimental and SM (theory) result:



$$a_\mu^{\text{exp}} - a_\mu^{\text{theo,SM}} \approx (25.1 \pm 5.9) \times 10^{-10} : 4.2 \sigma$$

$$a_\mu^{\text{exp}} - a_\mu^{\text{theo,SM}} \approx (25.1 \pm 5.9) \times 10^{-10} : 4.2\sigma$$

discrepancy $\approx 2 \times a_\mu^{\text{SM,weak}}$
but: expect $a_\mu^{\text{NP}} \sim a_\mu^{\text{SM,weak}} \times \left(\frac{M_W}{M_{\text{NP}}}\right)^2 \times \text{couplings}$

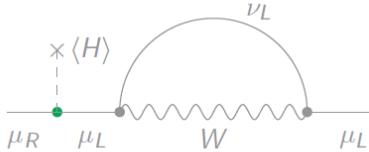
$$a_\mu \sim \frac{m_\mu \times (\text{some VEV}) \times (\mu_{L \leftrightarrow R}\text{-flipping parameter})}{M_{\text{typical}}^2}$$

$$a_\mu^{\text{exp}} - a_\mu^{\text{theo,SM}} \approx (25.1 \pm 5.9) \times 10^{-10} : 4.2\sigma$$

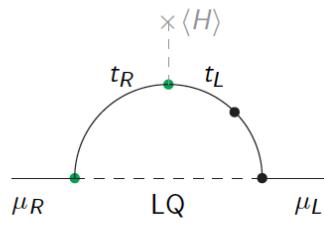
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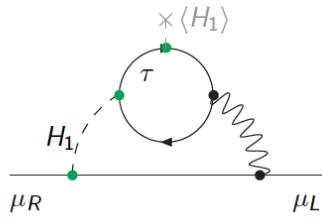
- EWSM: $\alpha \frac{m_\mu^2}{M_W^2}$



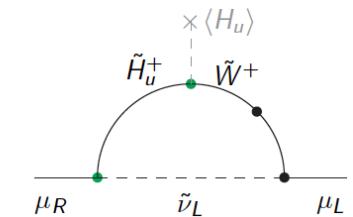
- LQ: $g_L g_R \frac{m_\mu m_t}{M_{\text{LQ}}^2}$



- 2HDM: $\alpha^2 \tan^2 \beta \frac{m_\mu^2}{M_H^2}$



- SUSY: $\alpha \frac{m_\mu^2 \tan \beta}{M_{\text{SUSY}}^2} \frac{\mu}{M_{\text{SUSY}}}$



- rad. $m_\mu \sim \frac{m_\mu^2}{M_{\text{NP}}^2}$

2. “Two Higgs Doublets are as natural as one” (2HDM)

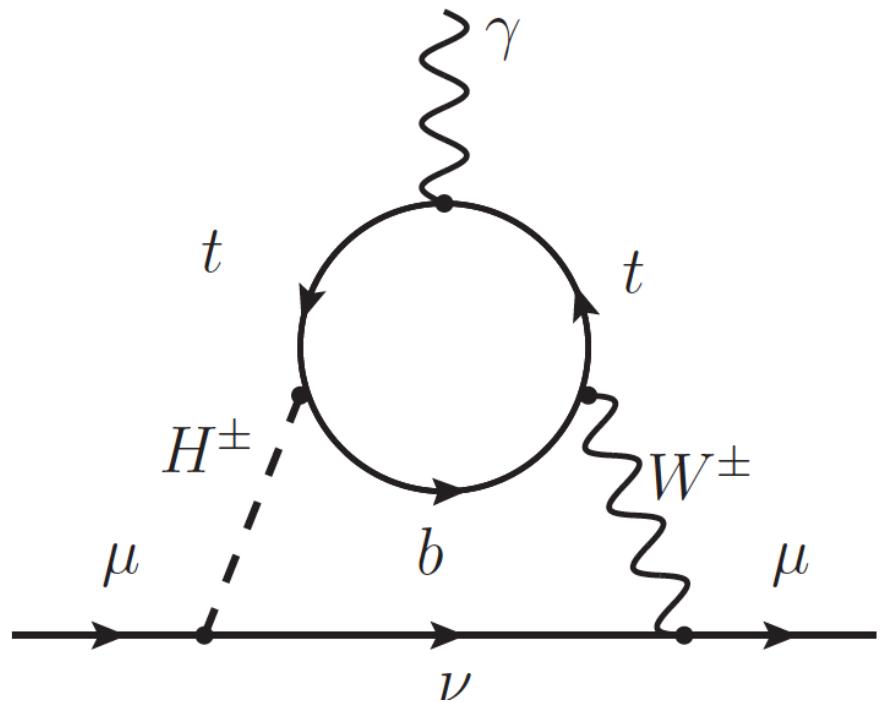
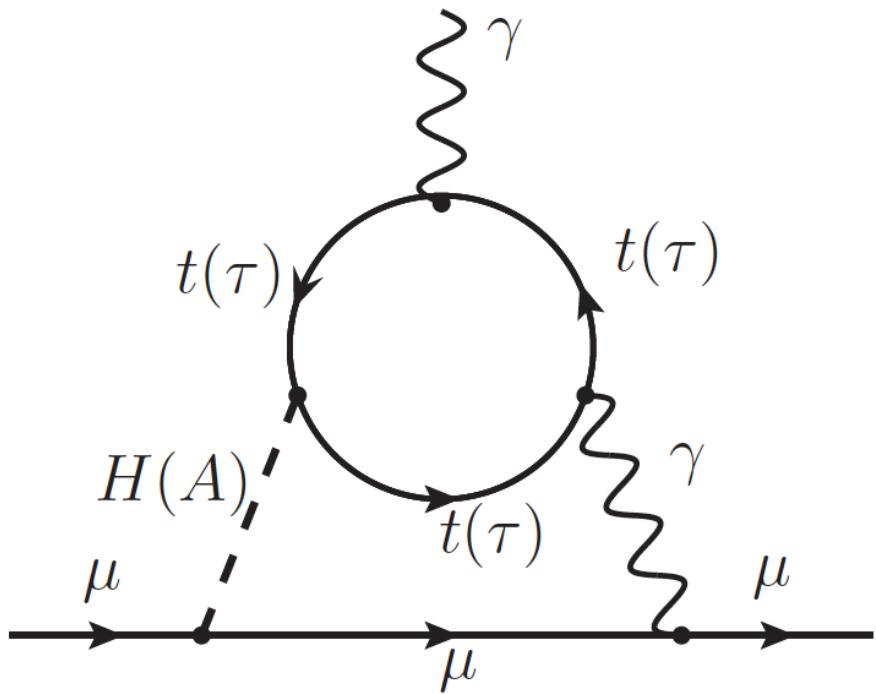
Extension of the Z_2 symmetry to fermions determines four types:

	u -type	d -type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
type III (lepton-specific)	Φ_2	Φ_2	Φ_1
type IV (flipped)	Φ_2	Φ_1	Φ_2

Couplings to fermions:

	u -type (c_{Att})	d -type (c_{Abb})	leptons ($c_{A\tau\tau}$)
type I	$\cot \beta$	$\cot \beta$	$\cot \beta$
type II	$\cot \beta$	$\tan \beta$	$\tan \beta$
type III/X (lepton-specific)	$\cot \beta$	$\cot \beta$	$\tan \beta$
type IV/Y (flipped)	$\cot \beta$	$\tan \beta$	$\cot \beta$

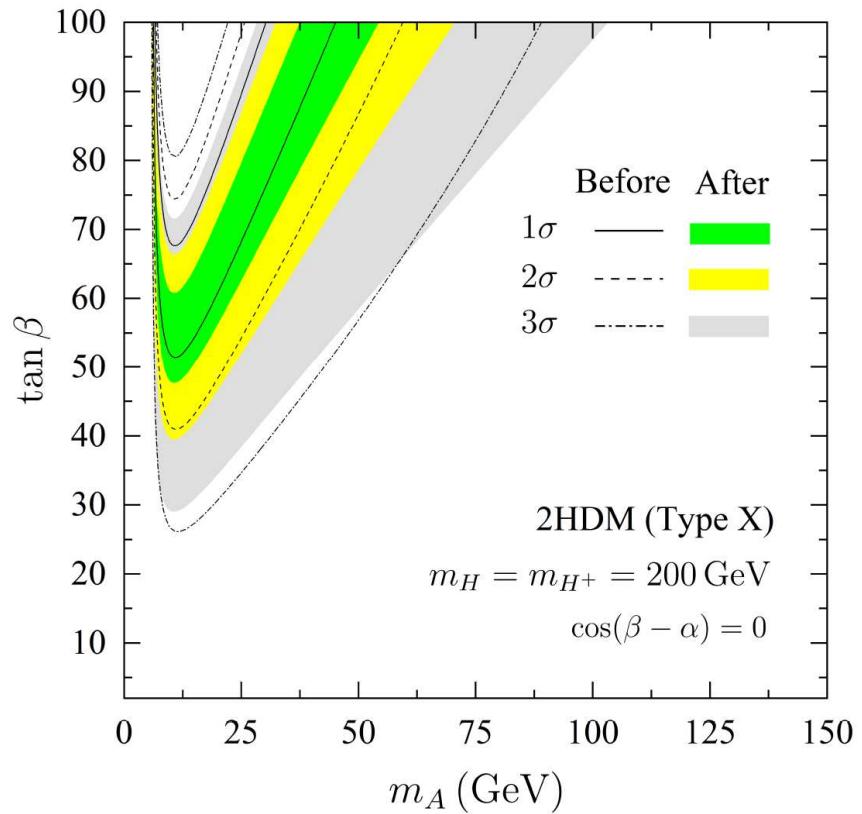
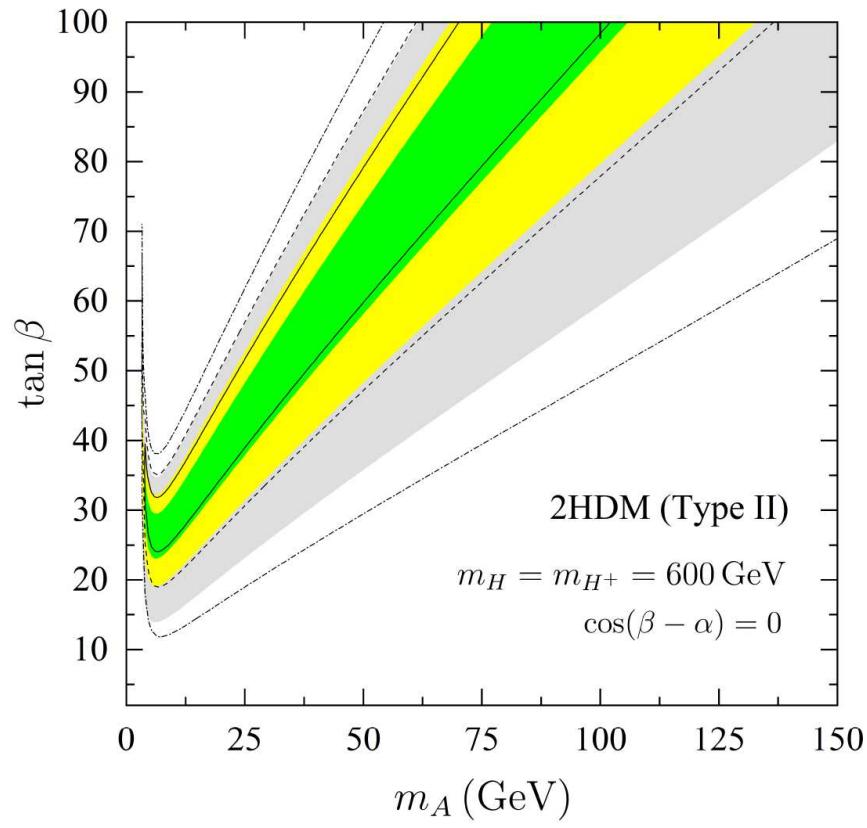
Barr-Zee diagrams:



Known since long:

Light A at high $\tan\beta$ can explain the $(g-2)_\mu$ result

⇒ in type II and X: enhanced couplings to μ 's



Taking EWPO into account: $\Rightarrow M_{H^\pm} \lesssim 200$ GeV

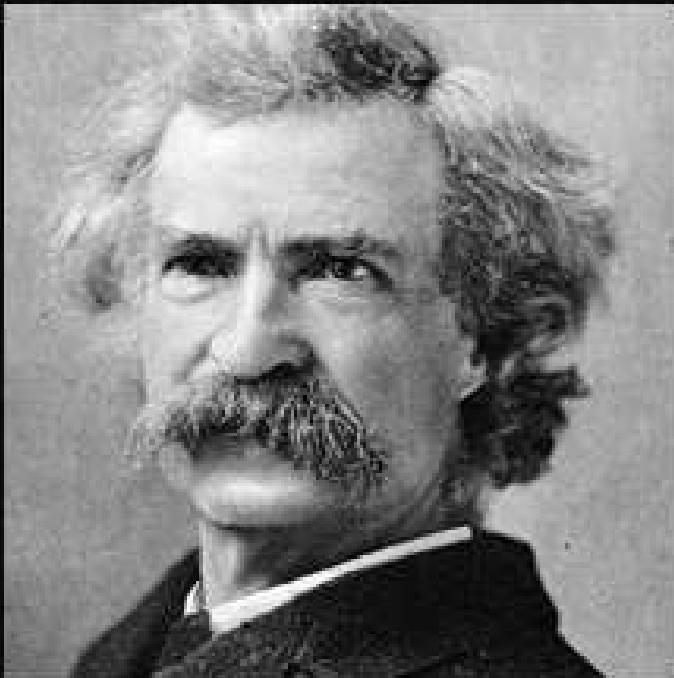
type II: problems with $b \rightarrow s\gamma$

\Rightarrow only type X is in agreement with $(g-2)_\mu$, EWPO and flavor observables

\Rightarrow Not taken into account here: $h_{125} \rightarrow AA \dots$

3. “SUSY fits like a glove”

→ quote by Bill Marciano after the BNL result came out in 2001

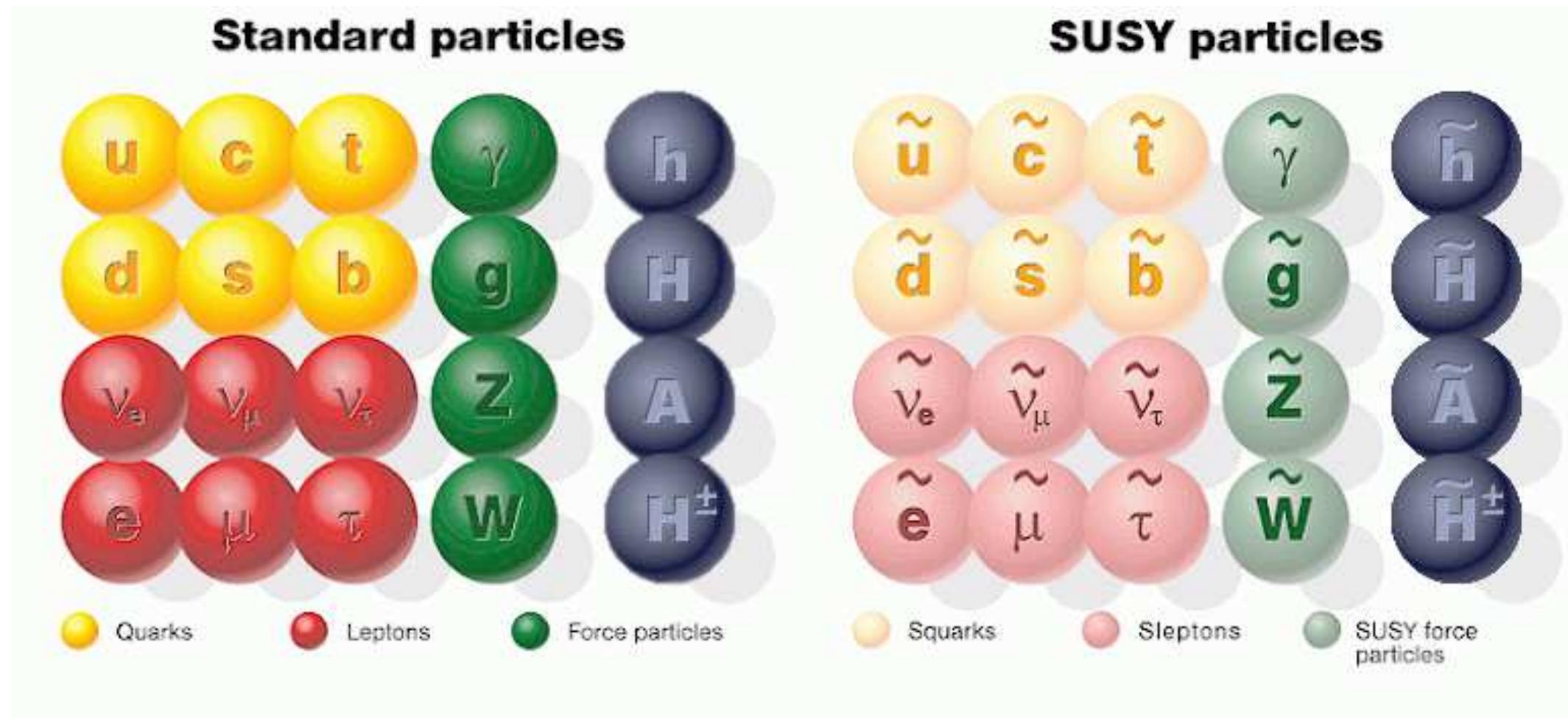


The reports of my death have
been greatly exaggerated.

~ Mark Twain

Simplest SUSY realization: the MSSM

Superpartners for Standard Model particles



⇒ large uncolored / EW sector

Neutralinos and charginos:

Higgsinos and electroweak gauginos mix

charged:

$$\tilde{W}^+, \tilde{h}_u^+ \rightarrow \tilde{\chi}_1^+, \tilde{\chi}_2^+, \quad \tilde{W}^-, \tilde{h}_d^- \rightarrow \tilde{\chi}_1^-, \tilde{\chi}_2^-$$

Diagonalization of the mass matrix:

$$\mathbf{X} = \begin{pmatrix} M_2 & \sqrt{2} \sin \beta M_W \\ \sqrt{2} \cos \beta M_W & \mu \end{pmatrix},$$

$$\mathbf{M}_{\tilde{\chi}^-} = \mathbf{V}^* \mathbf{X}^\top \mathbf{U}^\dagger = \begin{pmatrix} m_{\tilde{\chi}_1^\pm} & 0 \\ 0 & m_{\tilde{\chi}_2^\pm} \end{pmatrix}$$

⇒ charginos: mass eigenstates

mass matrix given in terms of M_2 , μ , $\tan \beta$

neutral:

$$\underbrace{\tilde{\gamma}, \tilde{Z}, \tilde{h}_u^0, \tilde{h}_d^0}_{\tilde{W}^0, \tilde{B}^0} \rightarrow \tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$$

Diagonalization of mass matrix:

$$Y = \begin{pmatrix} M_1 & 0 & -M_Z s_W \cos \beta & M_Z s_W \sin \beta \\ 0 & M_2 & M_Z c_W \cos \beta & -M_Z c_W \sin \beta \\ -M_Z s_W \cos \beta & M_Z c_W \cos \beta & 0 & -\mu \\ M_Z s_W \sin \beta & -M_Z c_W \sin \beta & -\mu & 0 \end{pmatrix},$$

$$M_{\tilde{\chi}^0} = N^* Y N^\dagger = \text{diag}(m_{\tilde{\chi}_1^0}, m_{\tilde{\chi}_2^0}, m_{\tilde{\chi}_3^0}, m_{\tilde{\chi}_4^0})$$

⇒ neutralinos: mass eigenstates

mass matrix given in terms of M_1 , M_2 , μ , $\tan \beta$

⇒ only one additional parameter

⇒ MSSM predicts mass relations between neutralinos and charginos

Scalar lepton sector of the MSSM

Charged slepton mass matrices

$$M_{\tilde{l}}^2 = \begin{pmatrix} M_{\tilde{l}_L}^2 + m_l^2 + DT_{l_1} & m_l X_l \\ m_l X_l & M_{\tilde{l}_R}^2 + m_l^2 + DT_{l_2} \end{pmatrix} \xrightarrow{\theta_{\tilde{l}}} \begin{pmatrix} m_{\tilde{l}_1}^2 & 0 \\ 0 & m_{\tilde{l}_2}^2 \end{pmatrix}$$

with

$$X_l = A_l - \mu \tan \beta$$

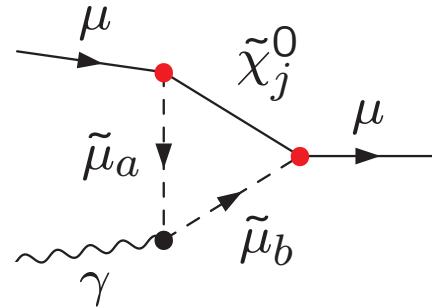
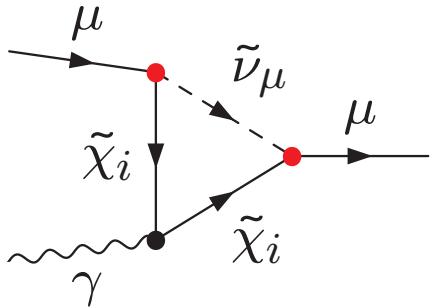
Sneutrino mass

$$m_{\tilde{\nu}_l}^2 = M_{\tilde{l}_L}^2 + DT_\nu$$

Simplifying assumption: $M_{\tilde{l}_L}$ and $M_{\tilde{l}_R}$ identical for all three generations

SUSY can easily explain the deviation in a_μ :

Feynman diagrams for MSSM 1L corrections:



- Diagrams with chargino/sneutrino exchange
- Diagrams with neutralino/smuon exchange

Enhancement factor as compared to SM:

$$\mu - \tilde{\chi}_i^\pm - \tilde{\nu}_\mu : \sim m_\mu \tan \beta$$

$$\mu - \tilde{\chi}_j^0 - \tilde{\mu}_a : \sim m_\mu \tan \beta$$

$$\text{SM, EW 1L: } \frac{\alpha}{\pi} \frac{m_\mu^2}{M_W^2}$$

$$\text{MSSM, 1L: } \frac{\alpha}{\pi} \frac{m_\mu^2}{M_{\text{SUSY}}^2} \times \tan \beta$$

SUSY corrections at 1L:

$$a_\mu^{\text{SUSY},1\text{L}} \approx 13 \times 10^{-10} \left(\frac{100 \text{ GeV}}{M_{\text{SUSY}}} \right)^2 \tan \beta \text{ sign}(\mu)$$

$M_{\text{SUSY}} (= m_{\tilde{\mu}} = m_{\tilde{\nu}} = m_{\tilde{\chi}})$: generic SUSY mass scale

$$a_\mu^{\text{SUSY},1\text{L}} = (-100 \dots + 100) \times 10^{-10}$$

$$a_\mu^{\text{exp}} - a_\mu^{\text{theo,SM}} \approx (28 \pm 7.4) \times 10^{-10}$$

⇒ SUSY could easily explain the “discrepancy”

⇒ a_μ can provide upper limits on the EW masses
(by requiring agreement at the 95% C.L.)

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⇒ SUSY could easily explain the “discrepancy”

⇒ a_μ can provide upper limits on the EW masses
(by requiring agreement at the 95% C.L.)

If SUSY exists, it should fix $(g-2)_\mu$!
⇒ there must be light EW SUSY particles!

The general idea:

[*M. Chakraborti, S.H., I. Saha '20, '21*]

- scan the relevant EW SUSY parameter space
- impose all relevant experimental constraints:
 - $(g - 2)_\mu$
 - Dark Matter relic density
 - Dark Matter direct detection
 - LHC searches for EW particles
- Dark Matter relic density requires a mechanism to reduce the density in the early universe
 - bino/wino DM with **chargino** co-annihilation
 - bino/wino DM with **slepton** co-annihilation
 - **higgsino** DM
 - **wino** DM
- obtain **lower and upper limits** on the various **EW particle masses**
- evaluate the prospects for future searches

$(g - 2)_\mu$ constraint: (GM2Calc)

$$\text{old: } \Delta a_\mu^{\text{old}} = (28.0 \pm 7.4) \times 10^{-10}$$

$$\text{new: } \Delta a_\mu^{\text{new}} = (25.1 \pm 5.9) \times 10^{-10}$$

\Rightarrow some results for $\Delta a_\mu^{\text{new}} (\equiv \Delta a_\mu)$

some results only available for $\Delta a_\mu^{\text{old}}$

Note: $\Delta a_\mu^{\text{old}} - 2\sigma^{\text{old}} \approx \Delta a_\mu^{\text{new}} - 2\sigma^{\text{new}}$

\Rightarrow upper limits on SUSY masses are not expected to change

Dark Matter relic density: MicrOmegas

$$\Omega_{\text{CDM}} h^2 = 0.120 \pm 0.001$$

or $\Omega_{\text{CDM}} h^2 \leq 0.122$

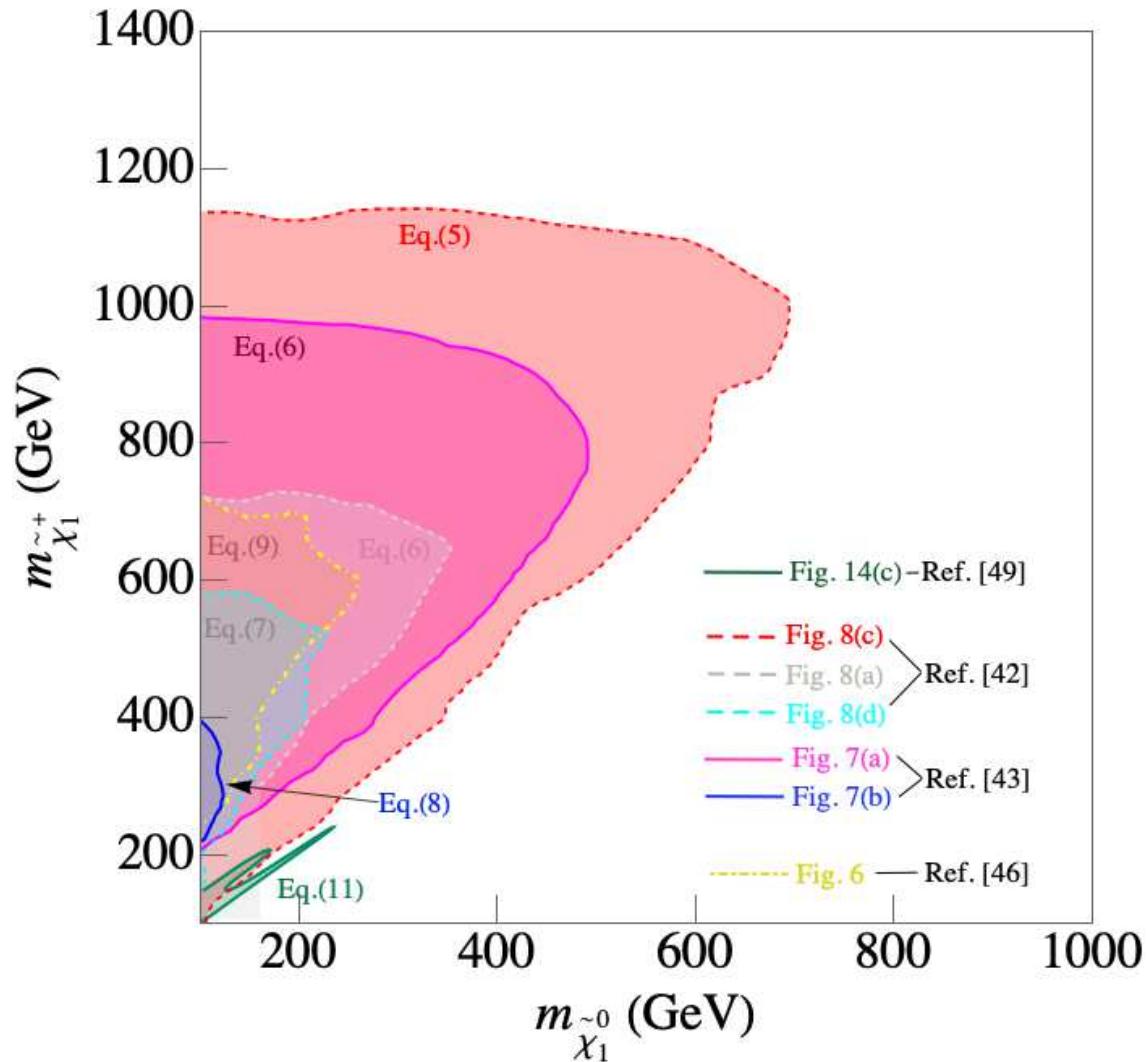
(as taken from [Planck '18])

Dark Matter direct detection: MicrOmegas

limit on spin independent scattering cross section (Xenon1T)

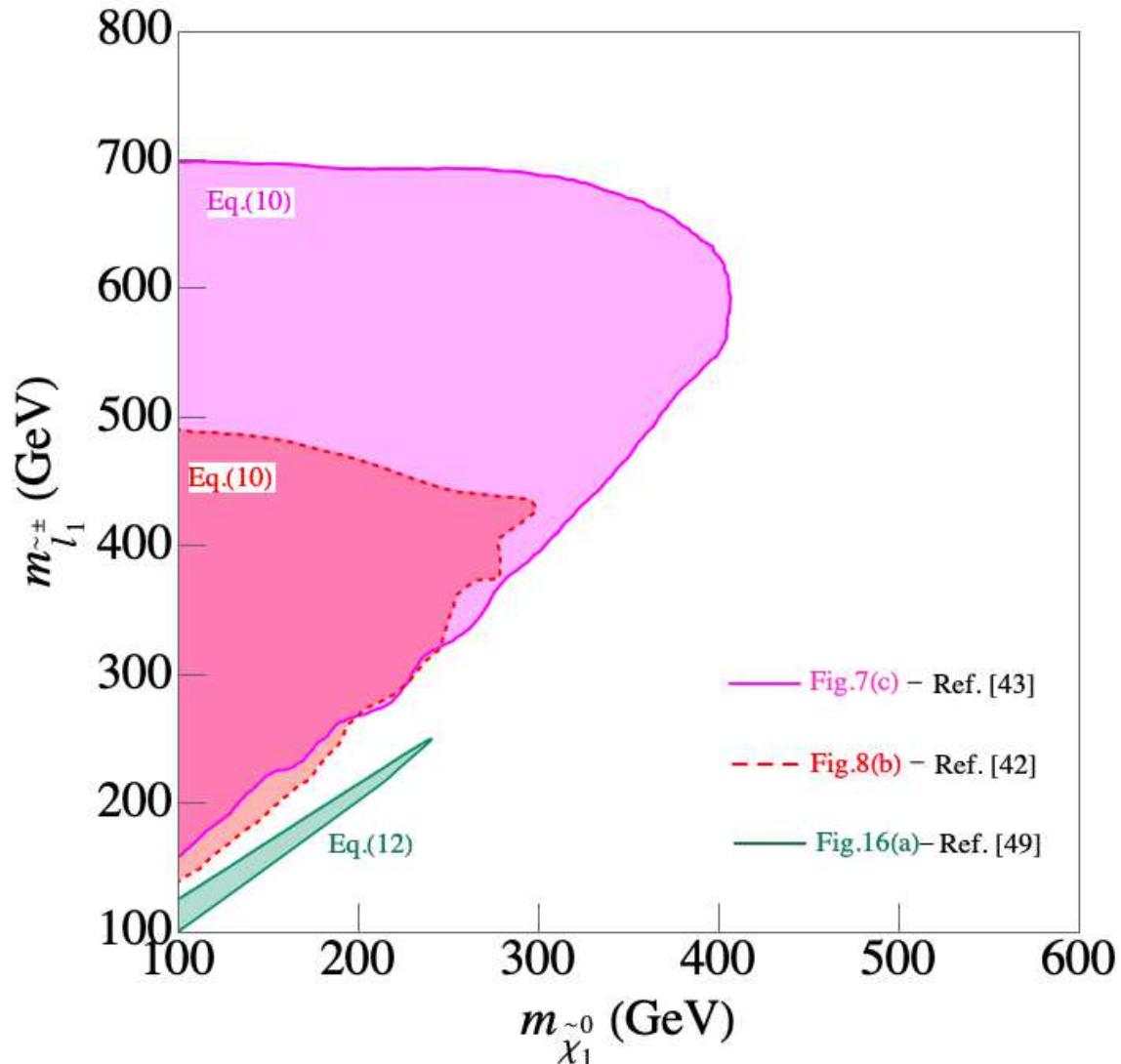
[Xenon collab. '18]

LHC exclusion bounds (I) (as given for Simplified Model Spectra (SMS))



⇒ all newly included into CheckMate [M.C & I.S.]
Exception: compressed spectra ⇒ direct application

LHC exclusion bounds (II) (as given for Simplified Model Spectra (SMS))



⇒ all newly included into CheckMate [M.C & I.S.]
Exception: compressed spectra ⇒ direct application

Possible scenarios:

1. bino/wino DM with **chargino** co-annihilation

\Rightarrow full relic DM constraint \Rightarrow updated with $\Delta a_\mu^{\text{new}}$

2. bino DM with **slepton** co-annihilation case-L

\Rightarrow full relic DM constraint \Rightarrow updated with $\Delta a_\mu^{\text{new}}$

3. bino DM with **slepton** co-annihilation case-R

\Rightarrow full relic DM constraint \Rightarrow updated with $\Delta a_\mu^{\text{new}}$

4. higgsino DM: $m_{\tilde{\chi}_1^0} \approx m_{\tilde{\chi}_2^0} \approx m_{\tilde{\chi}_1^\pm}$

full relic DM constraint $\Rightarrow m_{\tilde{\chi}_1^0} \sim 1 \text{ TeV} \Rightarrow (g - 2)_\mu$ not ok

\Rightarrow relic DM upper bound

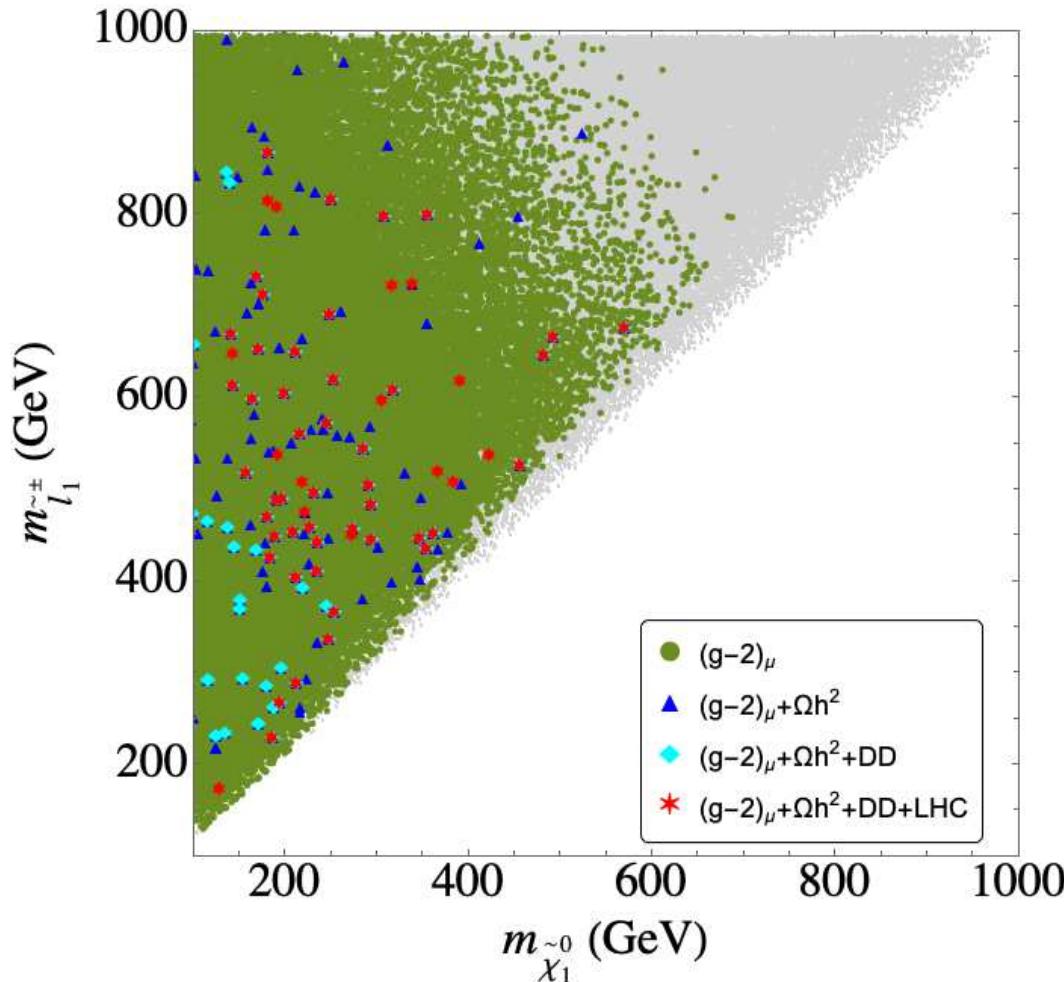
5. wino DM: $m_{\tilde{\chi}_1^0} \approx m_{\tilde{\chi}_1^\pm}$

full relic DM constraint $\Rightarrow m_{\tilde{\chi}_1^0} \sim 3 \text{ TeV} \Rightarrow (g - 2)_\mu$ not ok

\Rightarrow relic DM upper bound

Example I: $\tilde{\chi}_1^\pm$ -coannihilation: $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1^\pm}$ plane:

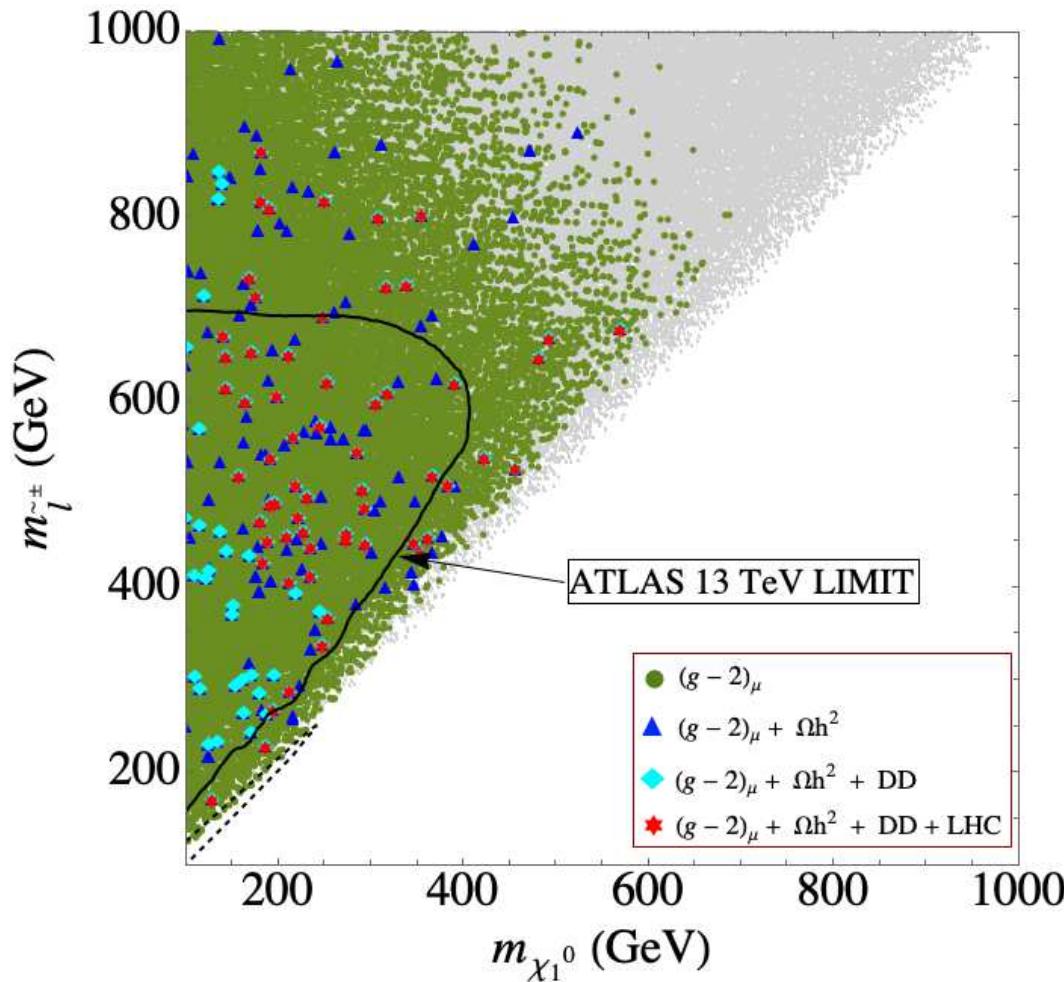
new $(g - 2)_\mu$



⇒ important: \tilde{l} -pair production searches (10)

Example I: $\tilde{\chi}_1^\pm$ -coannihilation: $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:

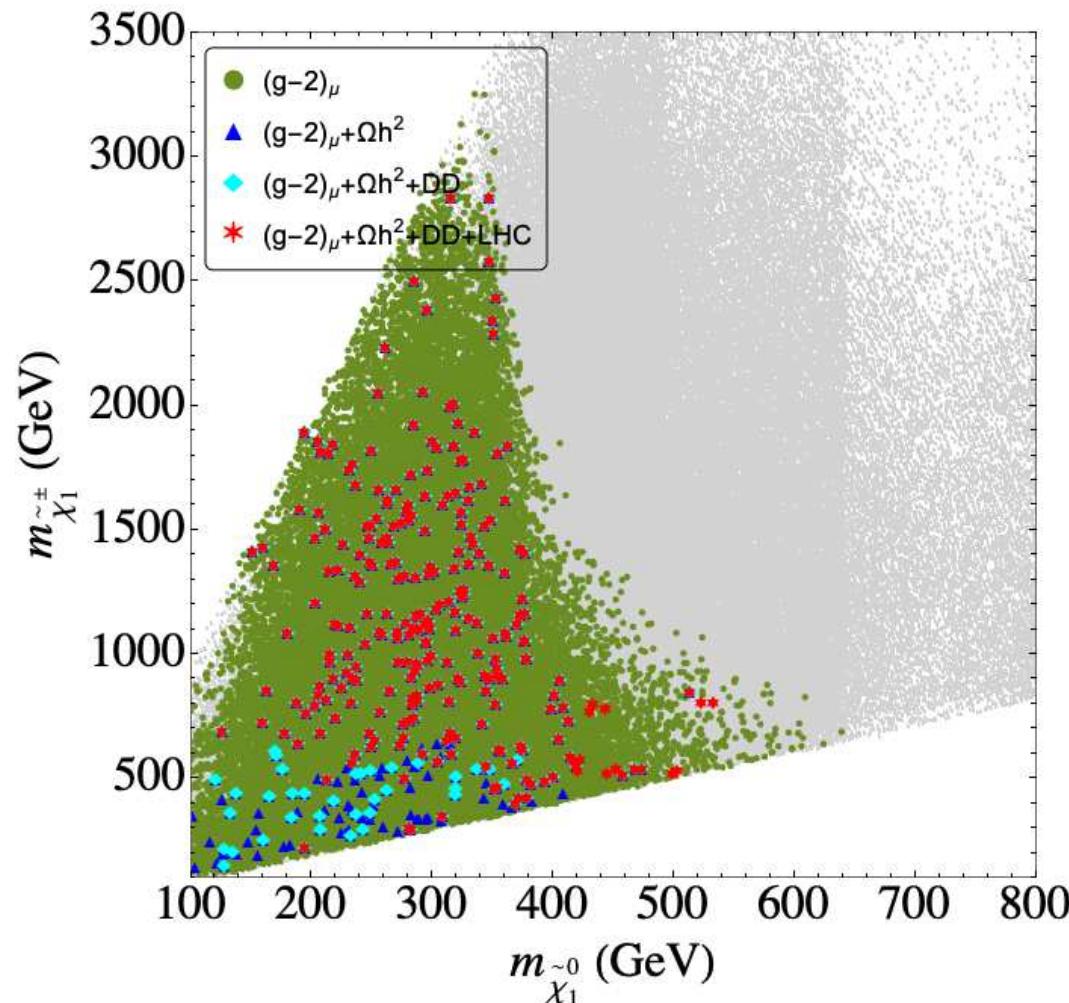
old $(g - 2)_\mu$



⇒ important: \tilde{l} -pair production searches (10)
⇒ naive application of LHC bounds fails

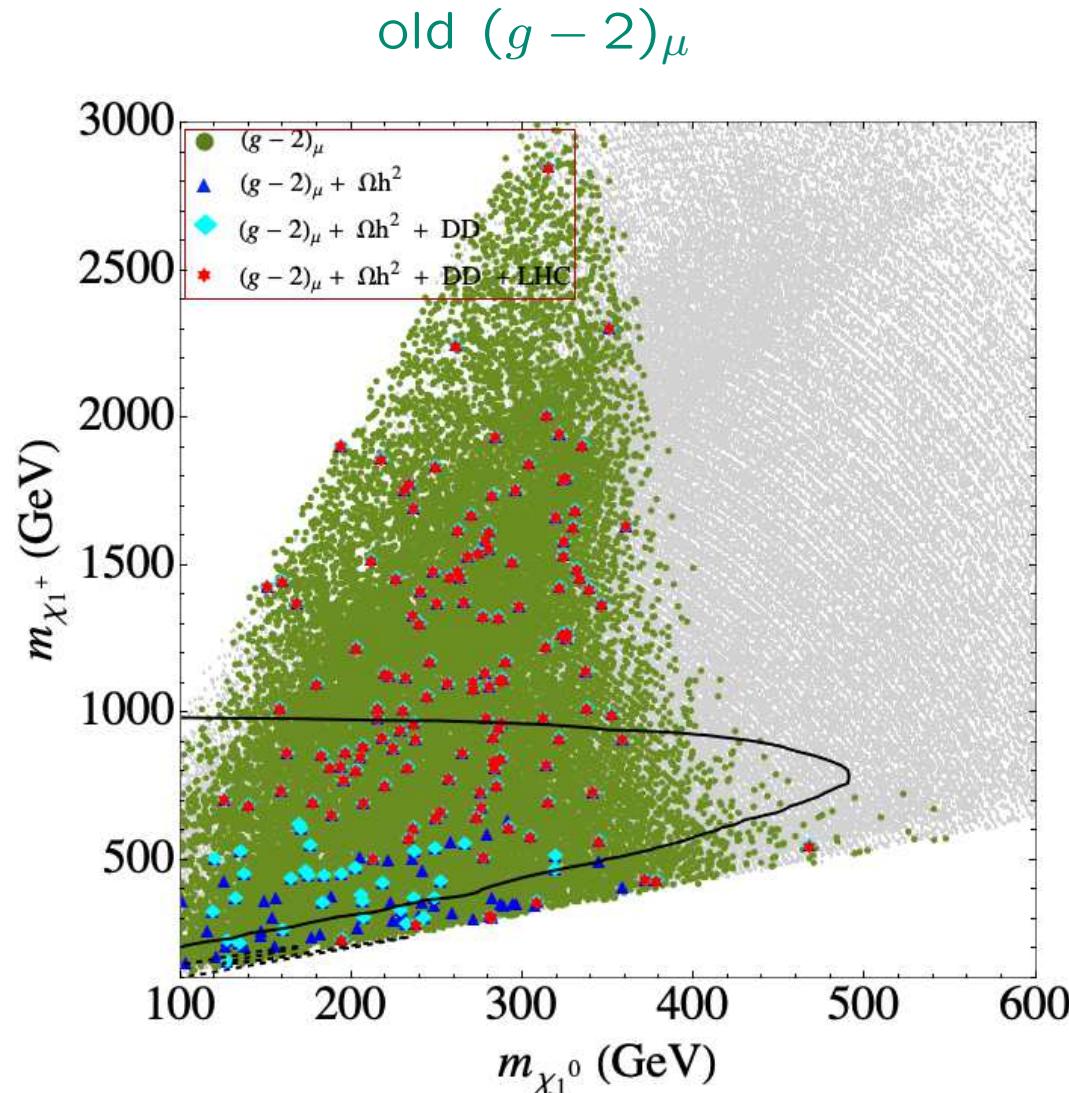
Example II: \tilde{t} -coannihilation: $m_{\tilde{\chi}_1^0} - m_{\tilde{\chi}_1^\pm}$ plane:

new $(g - 2)_\mu$



⇒ important: $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production searches (5)

Example II: \tilde{t} -coannihilation: $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{\chi}_1^\pm}$ plane:

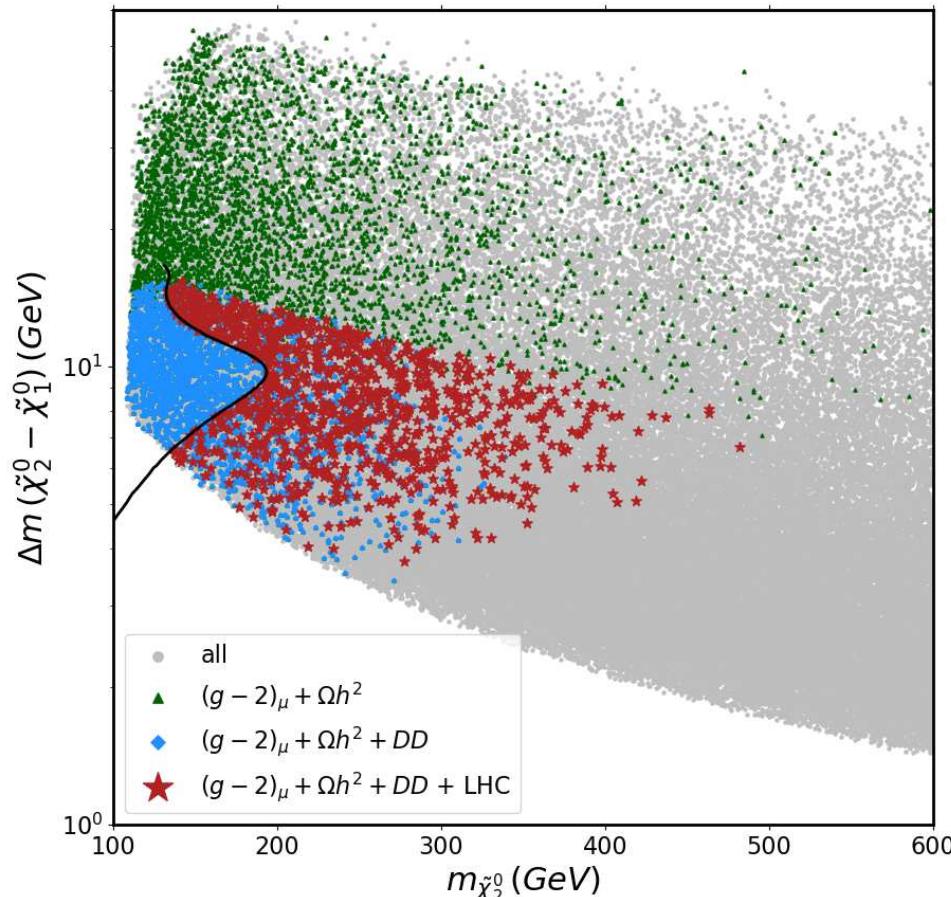


→ important: $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production searches (5)

→ naive application of LHC bounds fails

Example III: higgsino DM: $m_{\tilde{\chi}_2^0}$ - Δm plane:

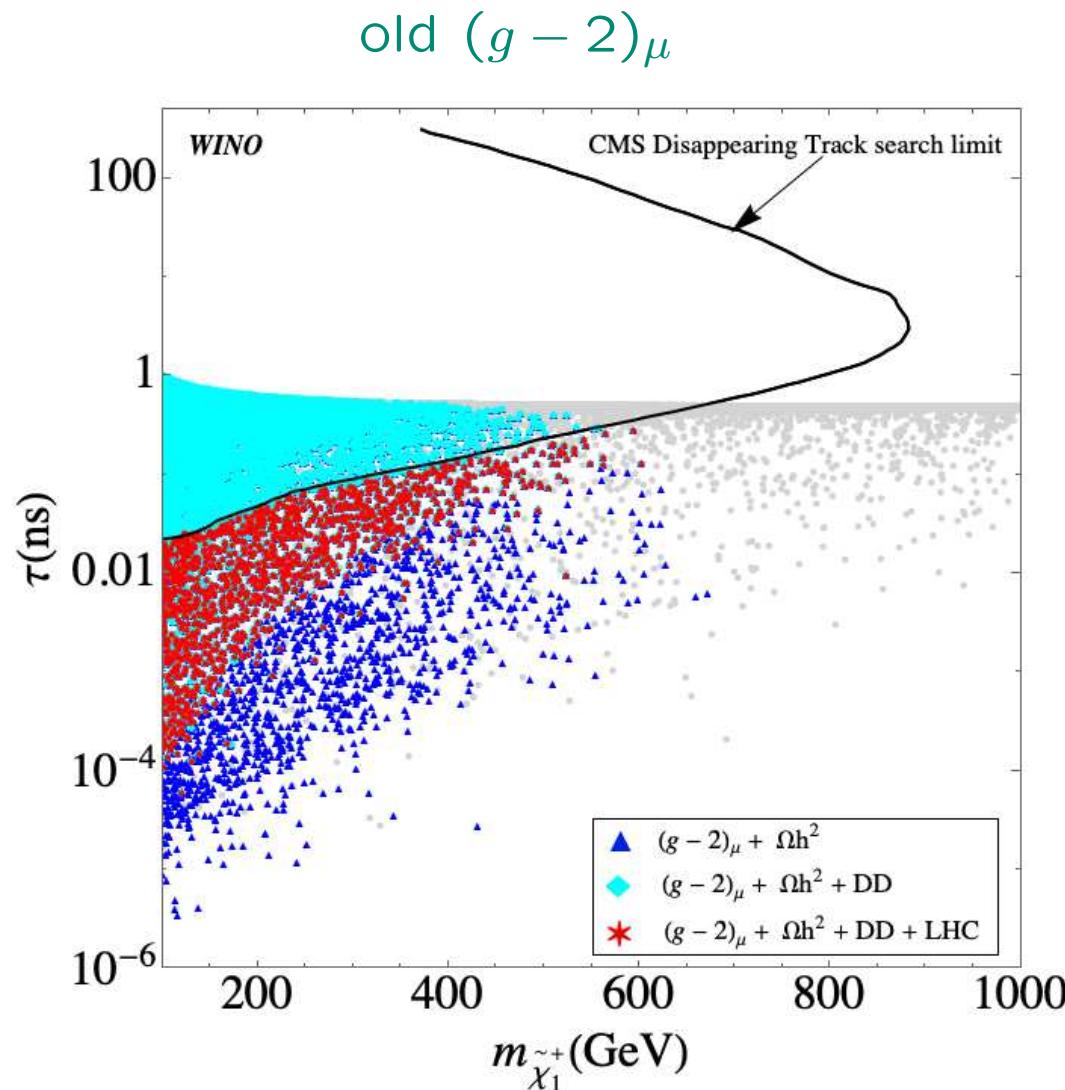
old $(g - 2)_\mu$



⇒ important: compressed spectra searches (11)

⇒ right where the model prediction sits ⇒ very powerful

Example IV: wino DM: $m_{\tilde{\chi}_1^\pm}$ - $\tau_{\tilde{\chi}_1^\pm}$ plane:



→ important: disappearing track limit $\Rightarrow m_{(N)\text{LSP}} \lesssim 600 \text{ GeV}$
 → allowed parameter space squeezed by DD limits and disapp. tracks

Mini summary of the results:

A) bino/wino DM with chargino co-annihilation

relic DM density 100% fulfilled

$\Rightarrow m_{(N)LSP} \lesssim 600(650)$ GeV for new (and old) $(g - 2)_\mu$

B/C) bino DM with slepton co-annihilation

relic DM density 100% fulfilled

$\Rightarrow m_{(N)LSP} \lesssim 550(600)$ GeV for new (and old) $(g - 2)_\mu$

D) higgsino DM: $m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_2^0} \sim m_{\tilde{\chi}_1^\pm} \sim \mu$

relic DM density as upper limit (otherwise $m_{\tilde{\chi}_1^0} \sim 1$ TeV)

$\Rightarrow m_{(N)LSP} \lesssim 500$ GeV

E) wino DM: $m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_1^\pm} \sim M_2$

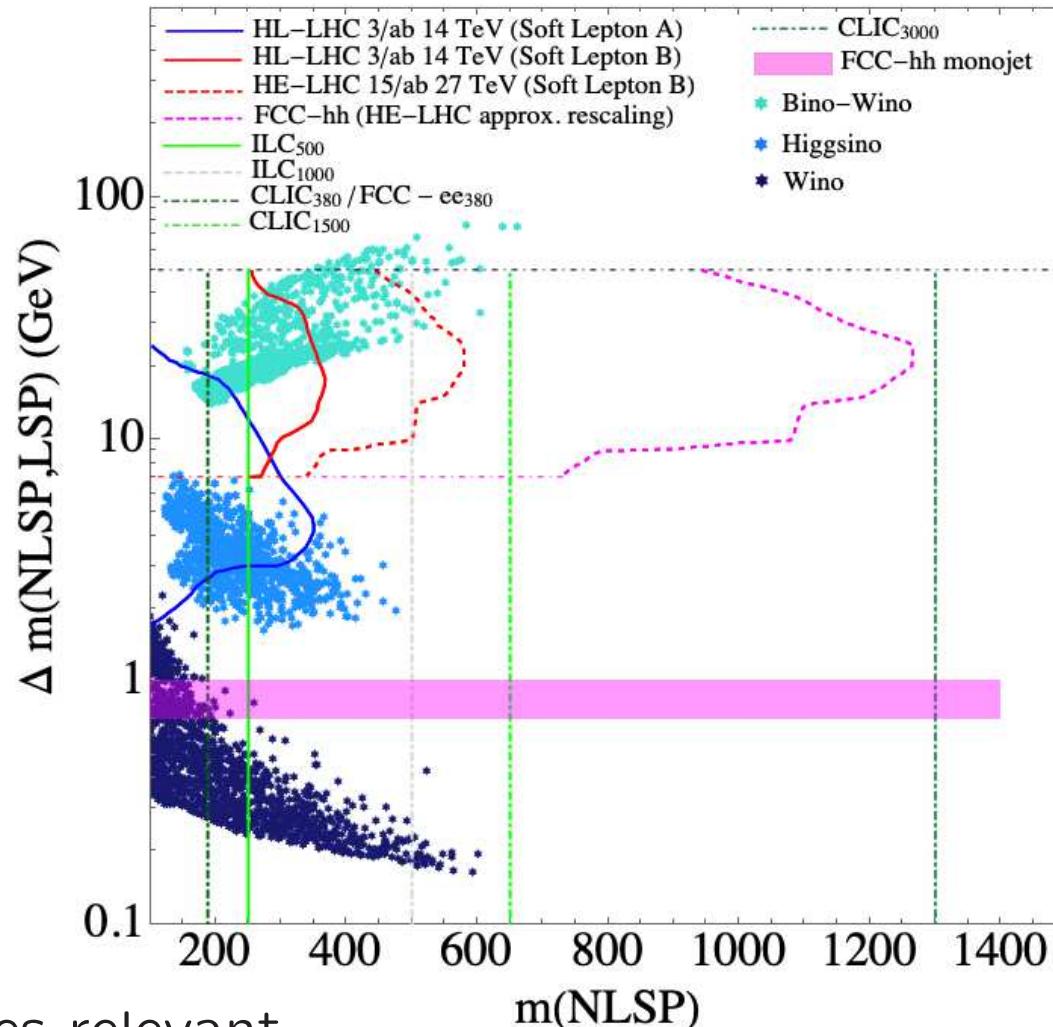
relic DM density as upper limit (otherwise $m_{\tilde{\chi}_1^0} \sim 3$ TeV)

$\Rightarrow m_{(N)LSP} \lesssim 600$ GeV

\Rightarrow predictions for future (e^+e^-) colliders?!

Compressed spectra at current and future colliders

Higgsino, wino and bino/wino DM:



- current searches relevant
- HL-LHC searches can cover some part of the parameter space
- ILC/CLIC needed to cover these scenario

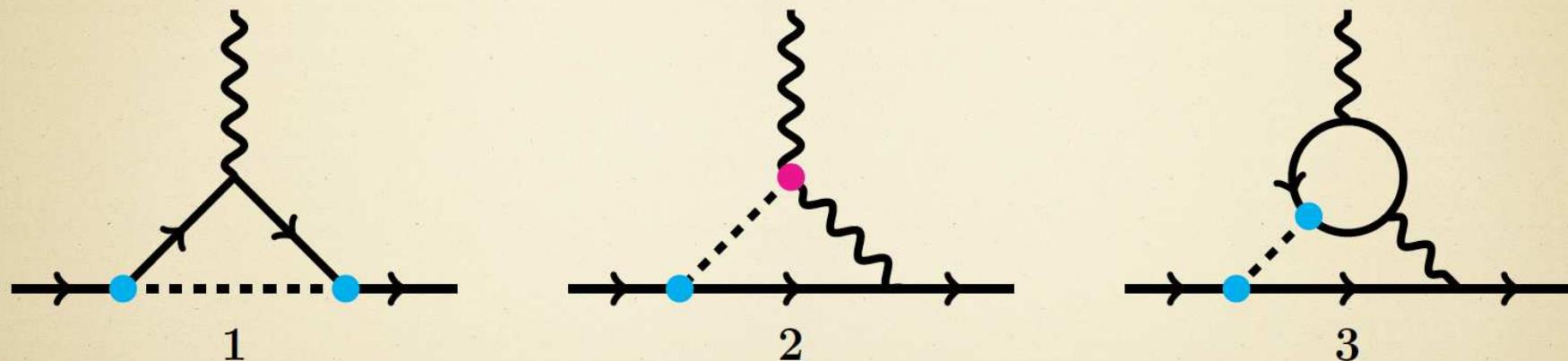
4. “There is always the axion”

→ quote by Howie Baer on WIMP DM limits



Lagrangian and possible contributions (I):

$$\mathcal{L}_{\text{eff}} \supset \frac{c_{ii}}{2} \frac{\partial_\mu a}{f_a} (\bar{\ell}_i \gamma^\mu \gamma^5 \ell_i) + c_{\gamma\gamma} \frac{\alpha}{4\pi} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \dots$$



$$\Delta a_\mu^{(1)} \propto \cancel{O} \frac{c_{\mu\mu}^2}{16\pi^2}, \quad \Delta a_\mu^{(2)} \propto -\frac{c_{\mu\mu} c_{\gamma\gamma} \alpha}{16\pi^3}, \quad \Delta a_\mu^{(3)} \propto -\frac{c_{\mu\mu} c_{ii} \alpha}{16\pi^3},$$

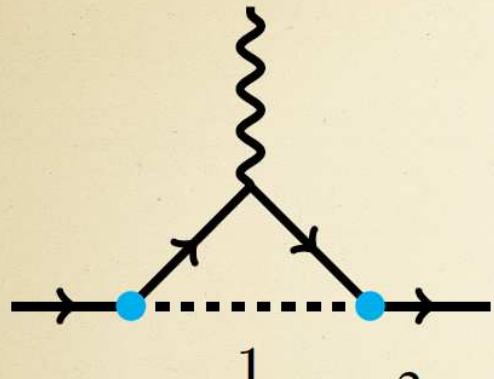
Chang et.al 2001; Marciano et.al. 2016; Bauer et.al. 2017.

Different combinations of the axion couplings: Darme et.al 2020

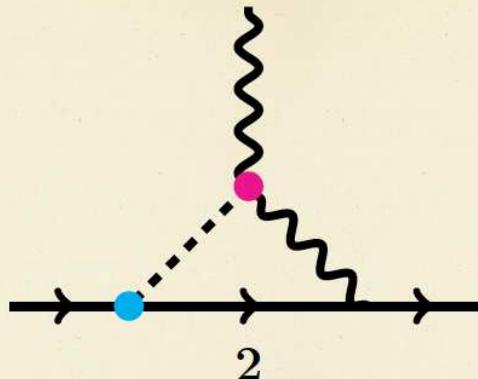
[taken from talk by JJ. Fan, g-2Days21]

Lagrangian and possible contributions (II):

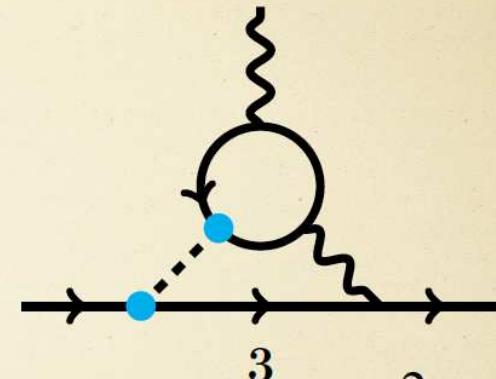
Case 1: only $c_{\mu\mu}$  , $c_{\gamma\gamma}$ 



$$\Delta a_\mu^{(1)} \propto -\frac{c_{\mu\mu}^2}{16\pi^2},$$



$$\Delta a_\mu^{(2)} \propto -\frac{c_{\mu\mu} c_{\gamma\gamma} \alpha}{16\pi^3}$$



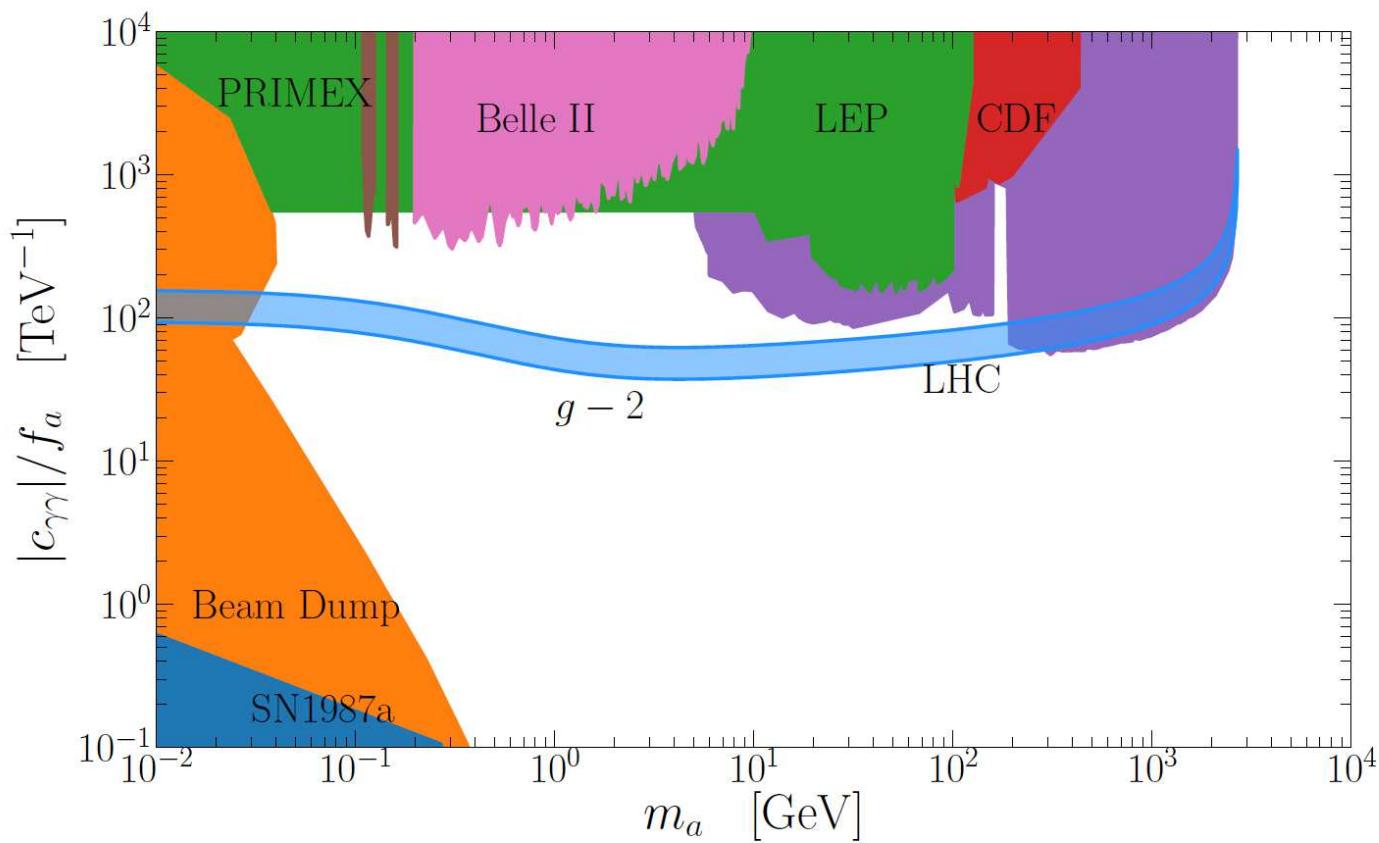
$$\Delta a_\mu^{(3)} \propto -\frac{c_{\mu\mu}^2 \alpha}{16\pi^3}$$

$$\Rightarrow c_{\mu\mu}/c_{\gamma\gamma} < 0$$

[taken from talk by JJ. Fan, g-2Days21]

Allowed parameter space for case 1:

[*M. Buen-Abad, JJ. Fan, M. Reece, C. Sun '21*]

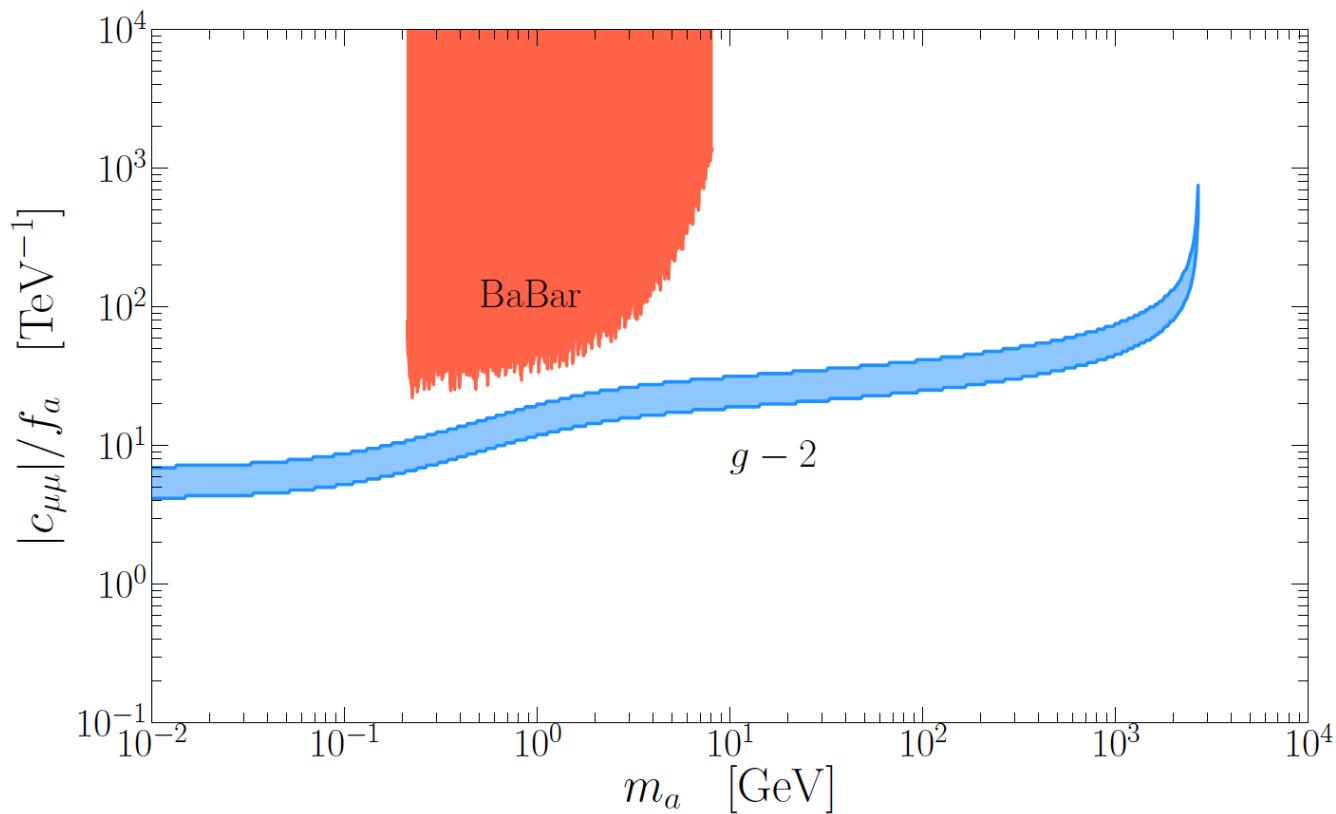


$$c_{\gamma\gamma}/c_{\mu\mu} < 0, \quad m_a \in [40 \text{ MeV}, 200 \text{ GeV}]$$

$$\left| \frac{f_a}{c_{\gamma\gamma}} \right| \lesssim 15 \dots 25 \text{ GeV}, \quad \left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \text{ GeV}$$

Allowed parameter space for case 1:

[*M. Buen-Abad, JJ. Fan, M. Reece, C. Sun '21*]

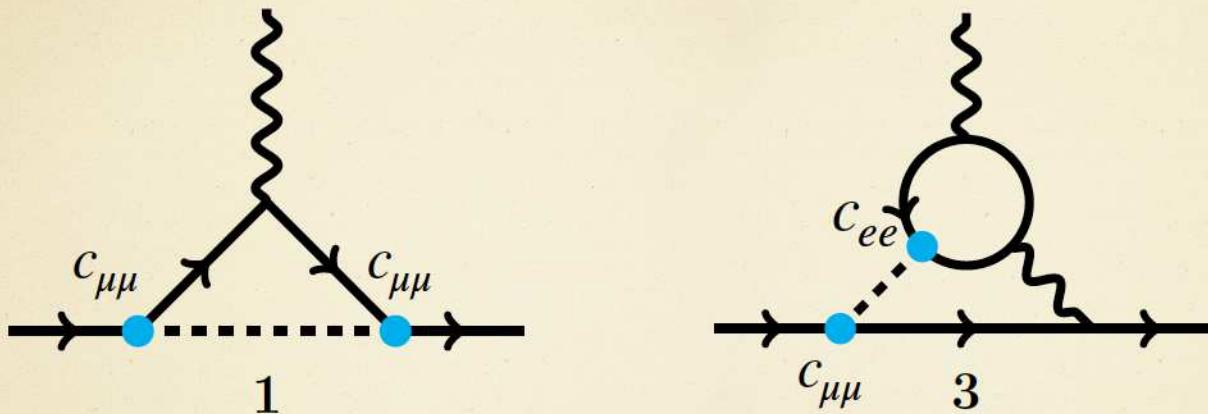


$$c_{\gamma\gamma}/c_{\mu\mu} < 0, \quad m_a \in [40 \text{ MeV}, 200 \text{ GeV}]$$

$$\left| \frac{f_a}{c_{\gamma\gamma}} \right| \lesssim 15 \dots 25 \text{ GeV}, \quad \left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \text{ GeV}$$

Lagrangian and possible contributions (III):

Case 2: only $c_{\mu\mu}$, c_{ee}



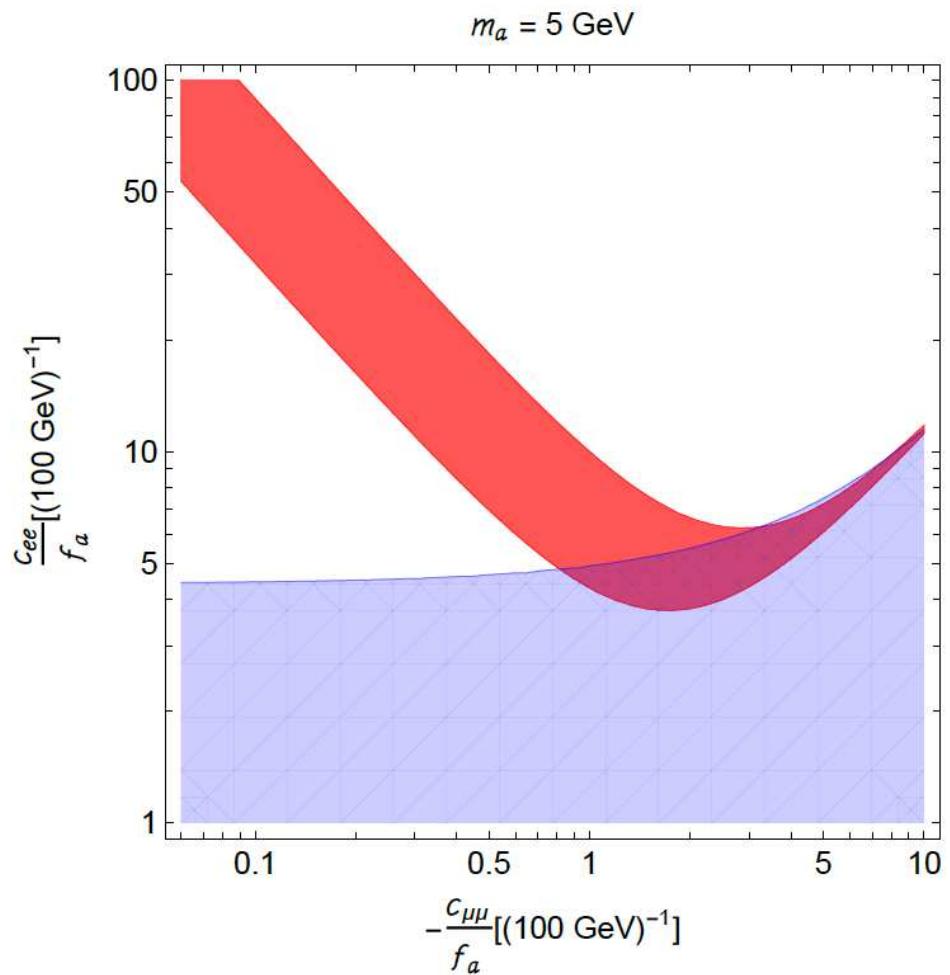
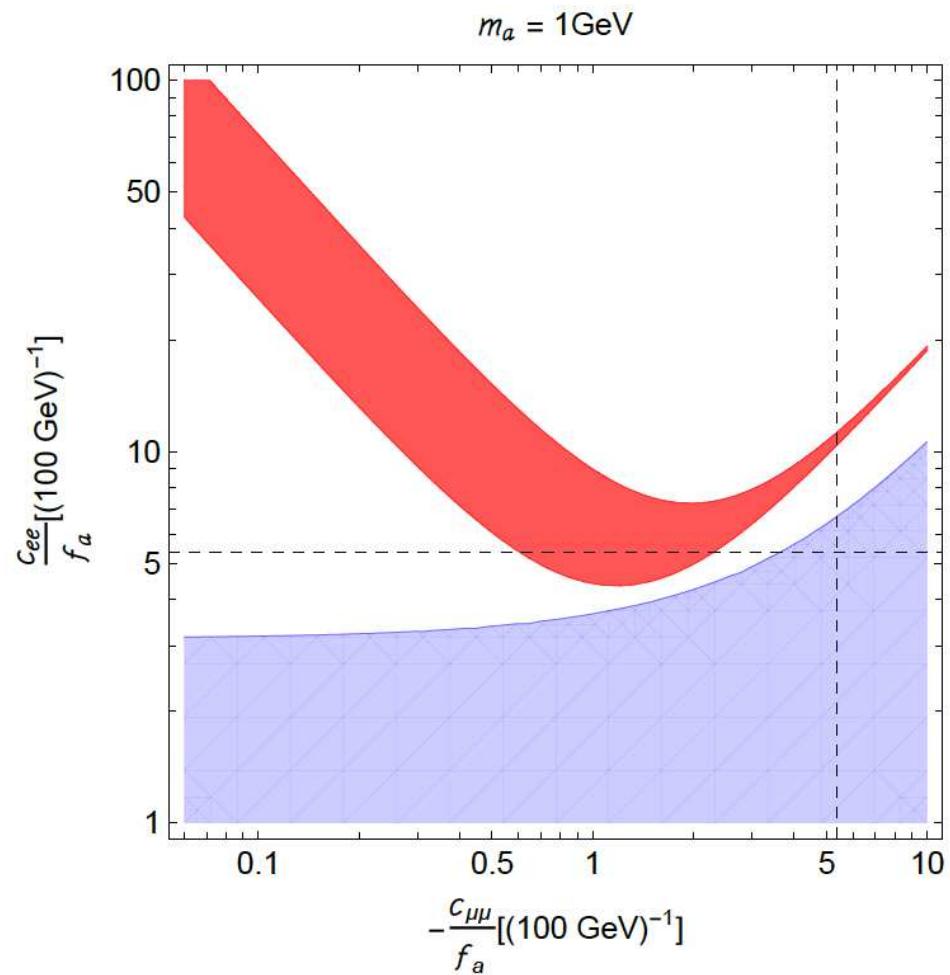
$$\Delta a_\mu^{(1)} \propto -\frac{c_{\mu\mu}^2}{16\pi^2}, \quad \Delta a_\mu^{(3)} \propto -\frac{c_{\mu\mu} c_{ee} \alpha}{16\pi^3}$$

$$\Rightarrow c_{\mu\mu}/c_{ee} < 0$$

[taken from talk by J. Fan, g-2Days21]

Allowed parameter space for case 2:

[M. Buen-Abad, JJ. Fan, M. Reece, C. Sun '21]



$$m_a \gtrsim 2 \text{ GeV}, \quad c_{ee}/c_{\mu\mu} < 0$$

$$\left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \text{ GeV}, \quad \left| \frac{f_a}{c_{ee}} \right| \lesssim 25 \text{ GeV} \quad \text{for } m_a = 5 \text{ GeV}$$

Case 1:

$$c_{\gamma\gamma}/c_{\mu\mu} < 0, \quad m_a \in [40 \text{ MeV}, 200 \text{ GeV}]$$

$$\left| \frac{f_a}{c_{\gamma\gamma}} \right| \lesssim 15 \dots 25 \text{ GeV}, \quad \left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \text{ GeV}$$

Case 2:

$$m_a \gtrsim 2 \text{ GeV}, \quad c_{ee}/c_{\mu\mu} < 0$$

$$\left| \frac{f_a}{c_{\mu\mu}} \right| \lesssim 100 \text{ GeV}, \quad \left| \frac{f_a}{c_{ee}} \right| \lesssim 25 \text{ GeV} \quad \text{for } m_a = 5 \text{ GeV}$$

⇒ axion couplings to SM must be unnaturally large!

⇒ only achieved by integrating out additional matter content

⇒ consider full UV-complete model?!

5. Conclusions

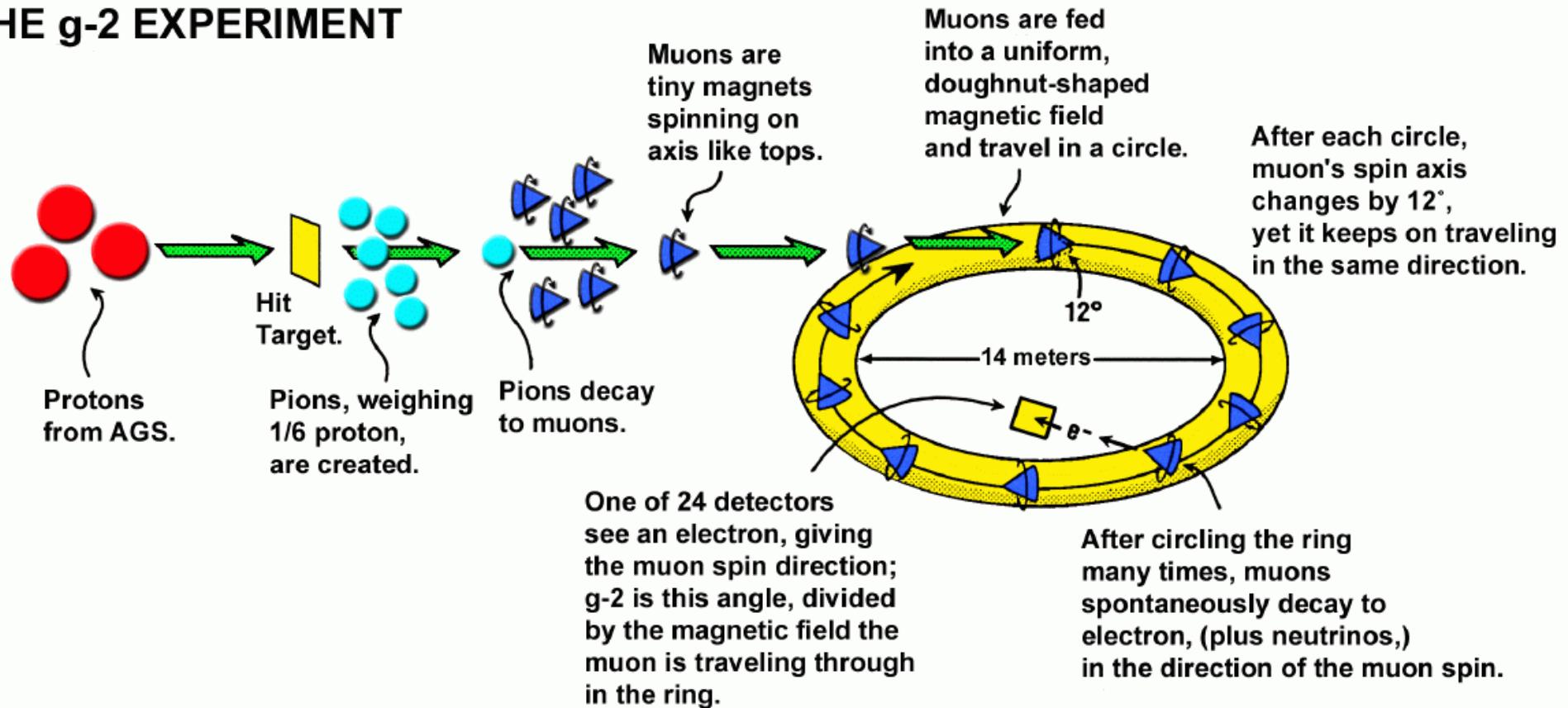
- new $(g - 2)_\mu$ result confirms old result and deviation from the SM
 $(g - 2)_\mu$ is real \Rightarrow (relatively) light EW particles
- 2HDM: only type X can be in agreement with $(g - 2)_\mu$, EWPO, flavor
 $\Rightarrow m_A \lesssim 50$ GeV, $\tan \beta \gtrsim 40$ \Rightarrow not taken into account (yet): $h_{125} \rightarrow AA$
- MSSM:
 - scan the EW sector of the MSSM with all constraints:
 $(g - 2)_\mu$, DM relic density, DM DD, LHC EW searches
 - upper limits on EW masses \Rightarrow evaluate future prospects
LHC searches included via **CheckMate** \Rightarrow crucial!
- A) bino/wino DM with chargino coann. (DM full)
B/C) bino DM with slepton coann. (DM full)
D) higgsino eDM $m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_2^0} \sim m_{\tilde{\chi}_1^\pm} \sim \mu$ (DM upper limit)
E) wino DM $m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_1^\pm} \sim M_2$ (DM upper limit)
 \Rightarrow clear upper limits, $m_{(N)LSP} \lesssim 600$ GeV confirmed
compressed spectra: good HL-LHC prospects, but ILC/CLIC needed
- Axions: in principle possible
 \Rightarrow axion couplings to SM must be **unnaturally large!**
 \Rightarrow integrating out additional matter content, full UV-complete model?



Further Questions?

The $(g - 2)_\mu$ experiment:

LIFE OF A MUON: THE g-2 EXPERIMENT

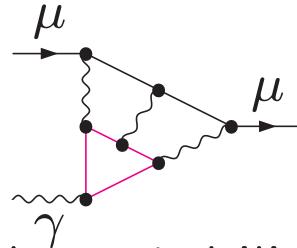


Coupling of muon to magnetic field : $\mu - \mu - \gamma$ coupling

$$\bar{u}(p') \left[\gamma^\mu F_1(q^2) + \frac{i}{2m_\mu} \sigma^{\mu\nu} q_\nu F_2(q^2) \right] u(p) A_\mu \quad F_2(0) = a_\mu$$

Theory of $(g - 2)_\mu$:

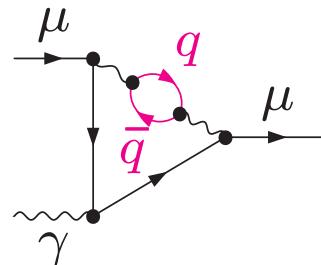
- the light-by-light contribution:



2002: sign error discovered; since then stabilized

2021: confirmed by LQCD

- the hadronic vacuum contribution:



'direct' e^+e^- data:

from **CMD-II, SND, KLOE, BaBar** (radiative return)

⇒ agree relatively well (also with old e^+e^- data)

⇒ tension with LQCD results

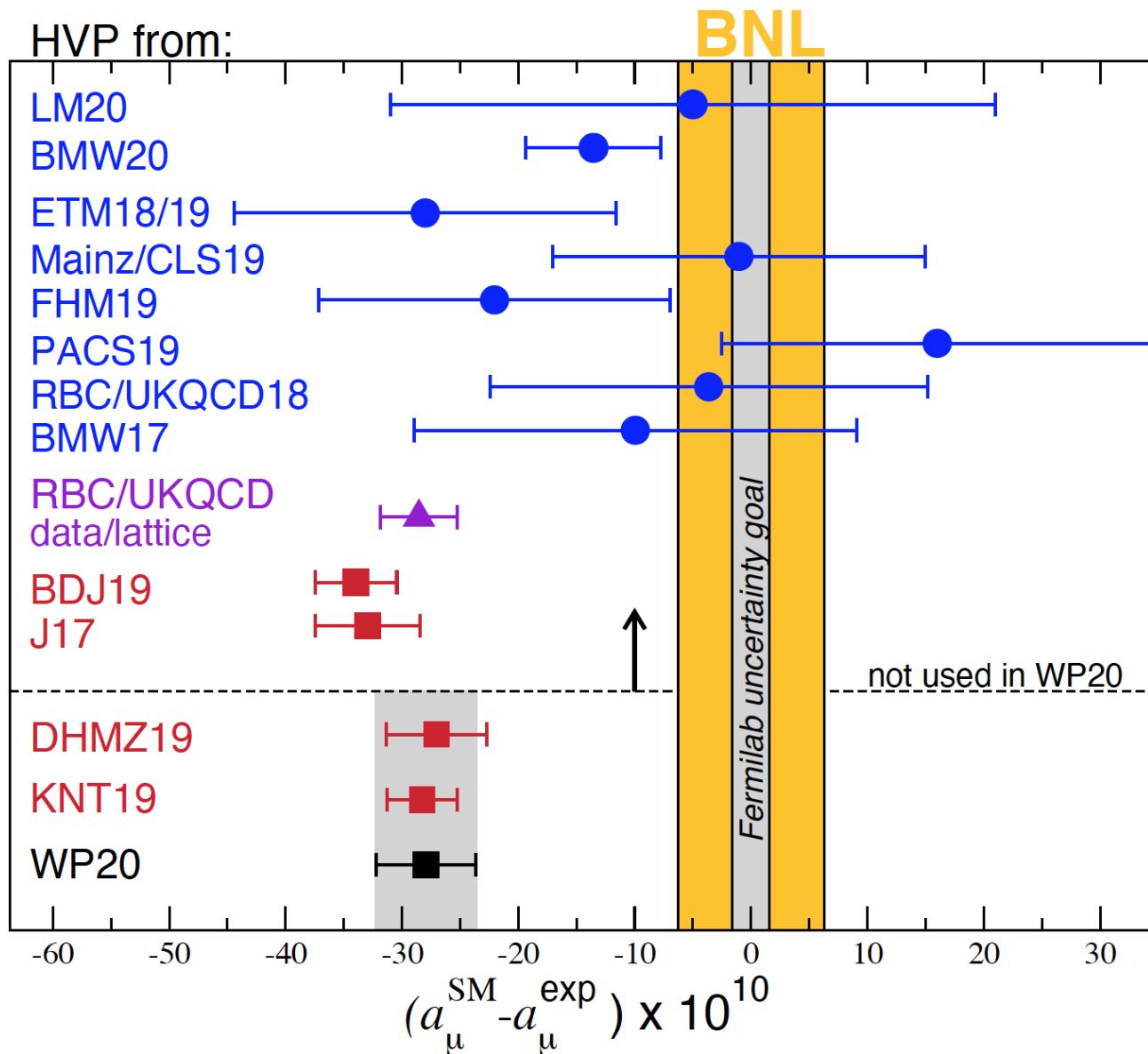
τ data:

tended to be closer to experimental result

inclusion of γ - ρ mixing: agreement with e^+e^- [F. Jegerlehner, R. Szafron '10]

⇒ not used anymore

HVP summary:



⇒ BMW20: difference to experimental data $\sim 1.5\sigma$

LHC searches: (as given for Simplified Model Spectra (SMS))

Decay via sleptons (3l)

$$\begin{aligned}\tilde{\chi}_1^\pm \tilde{\chi}_2^0 &\rightarrow (\tilde{l}^\pm \nu)(\tilde{l}^+ l^-) \rightarrow 3l + \cancel{E}_T , \\ \tilde{\chi}_1^\pm \tilde{\chi}_2^0 &\rightarrow (l^\pm \tilde{\nu})(\tilde{l}^+ l^-) \rightarrow 3l + \cancel{E}_T\end{aligned}\quad (5)$$

Decay via sleptons (2l)

$$\begin{aligned}\tilde{\chi}_1^+ \tilde{\chi}_1^- &\rightarrow (\tilde{l}^+ \nu)(\tilde{l}^- \nu) \rightarrow 2l + \cancel{E}_T , \\ \tilde{\chi}_1^+ \tilde{\chi}_1^- &\rightarrow (l^+ \tilde{\nu})(l^- \tilde{\nu}) \rightarrow 2l + \cancel{E}_T\end{aligned}\quad (6)$$

Decay via gauge bosons

$$\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow (W \tilde{\chi}_1^0)(Z \tilde{\chi}_1^0) \rightarrow 3l + \cancel{E}_T , \quad (7a)$$

$$\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow (W \tilde{\chi}_1^0)(Z \tilde{\chi}_1^0) \rightarrow 2l + \text{jets} + \cancel{E}_T , \quad (7b)$$

$$\tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow (W^+ \tilde{\chi}_1^0)(W^- \tilde{\chi}_1^0) \rightarrow 2l + \cancel{E}_T \quad (8)$$

Decay via Higgs bosons

$$\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow (W \tilde{\chi}_1^0)(h \tilde{\chi}_1^0) \rightarrow l + b\bar{b} + \cancel{E}_T \quad (9)$$

\tilde{l} -pair production (2l)

$$\tilde{l}^+ \tilde{l}^- \rightarrow (l^+ \tilde{\chi}_1^0)(l^- \tilde{\chi}_1^0) \rightarrow 2l + \cancel{E}_T \quad (10)$$

Compressed spectra

$$\tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow (W^* \tilde{\chi}_1^0)(Z^* \tilde{\chi}_1^0) \rightarrow 2l + \cancel{E}_T + \text{ISR} , \quad (11)$$

$$\tilde{l}^+ \tilde{l}^- \rightarrow (l^+ \tilde{\chi}_1^0)(l^- \tilde{\chi}_1^0) \rightarrow 2l + \cancel{E}_T + \text{ISR} \quad (12)$$

Searches involving Staus

→ all newly included into CheckMate [M.C & I.S.]

Exception: compressed spectra ⇒ direct application

A) Bino/wino DM with chargino co-annihilation

[*M. Chakraborti, S.H., I. Saha '20, '21*]

Parameter scan:

$$100 \text{ GeV} \leq M_1 \leq 1 \text{ TeV},$$

$$M_1 \leq M_2 \leq 1.1M_1,$$

$$1.1M_1 \leq \mu \leq 10M_1,$$

$$5 \leq \tan \beta \leq 60,$$

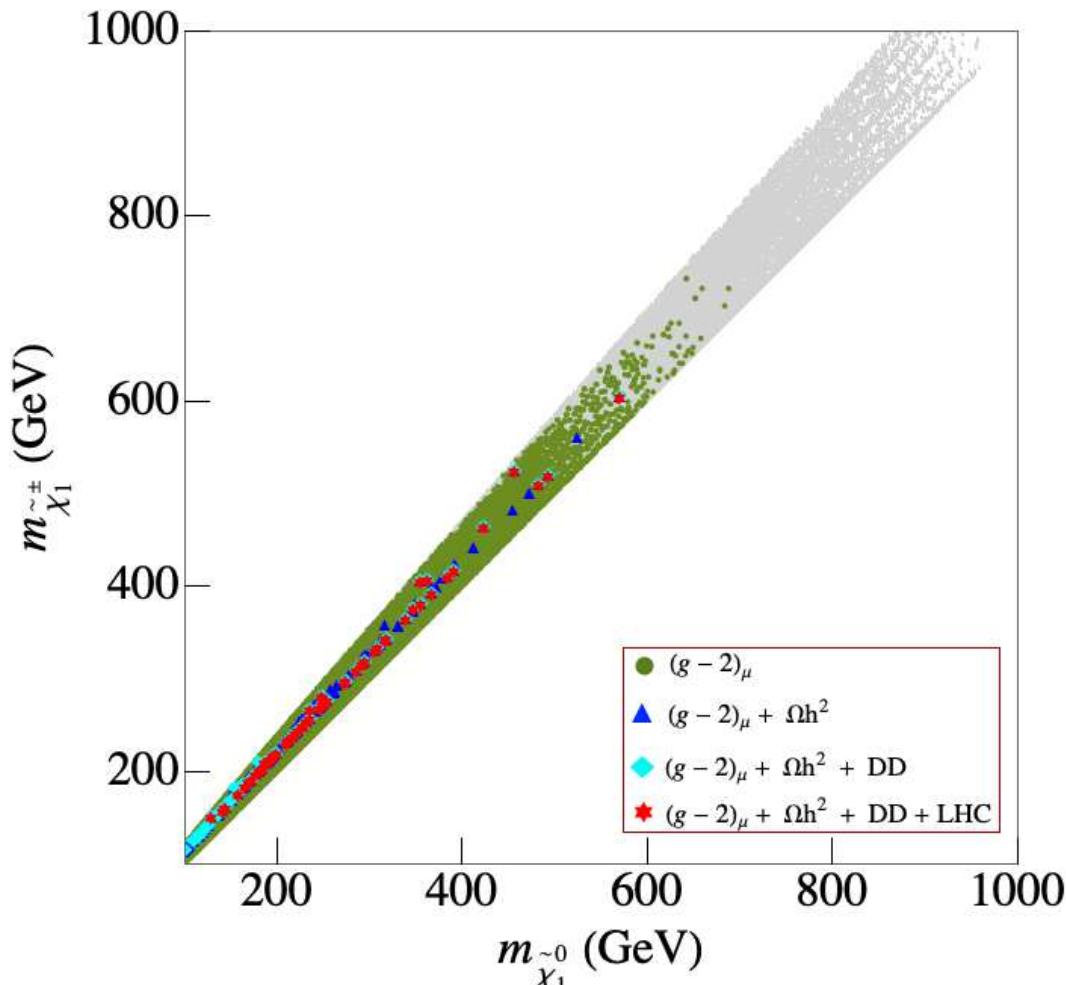
$$100 \text{ GeV} \leq m_{\tilde{L}} \leq 1 \text{ TeV},$$

$$m_{\tilde{R}} = m_{\tilde{L}}.$$

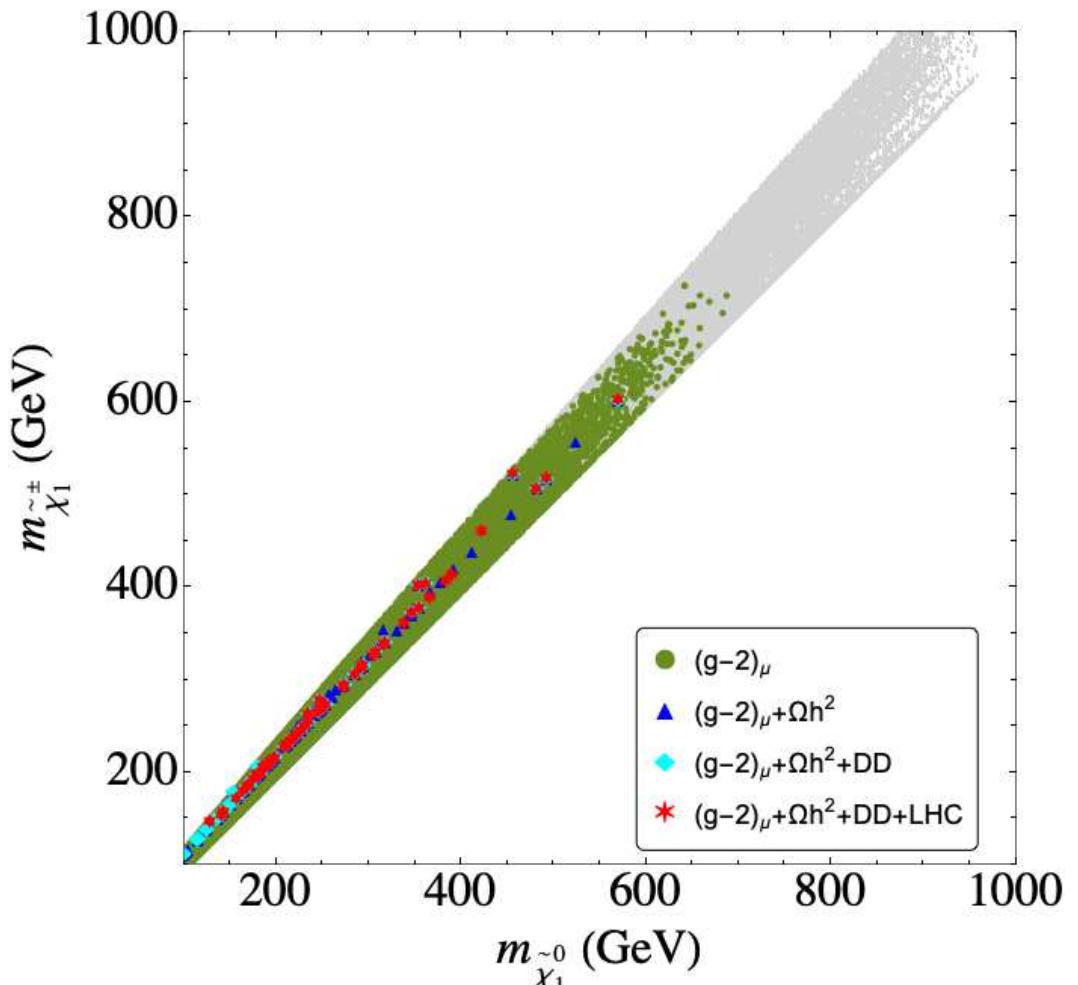
(latter condition only to make the analysis simpler, no relevant effect)

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{\chi}_1^\pm}$ plane:

old $(g - 2)_\mu$



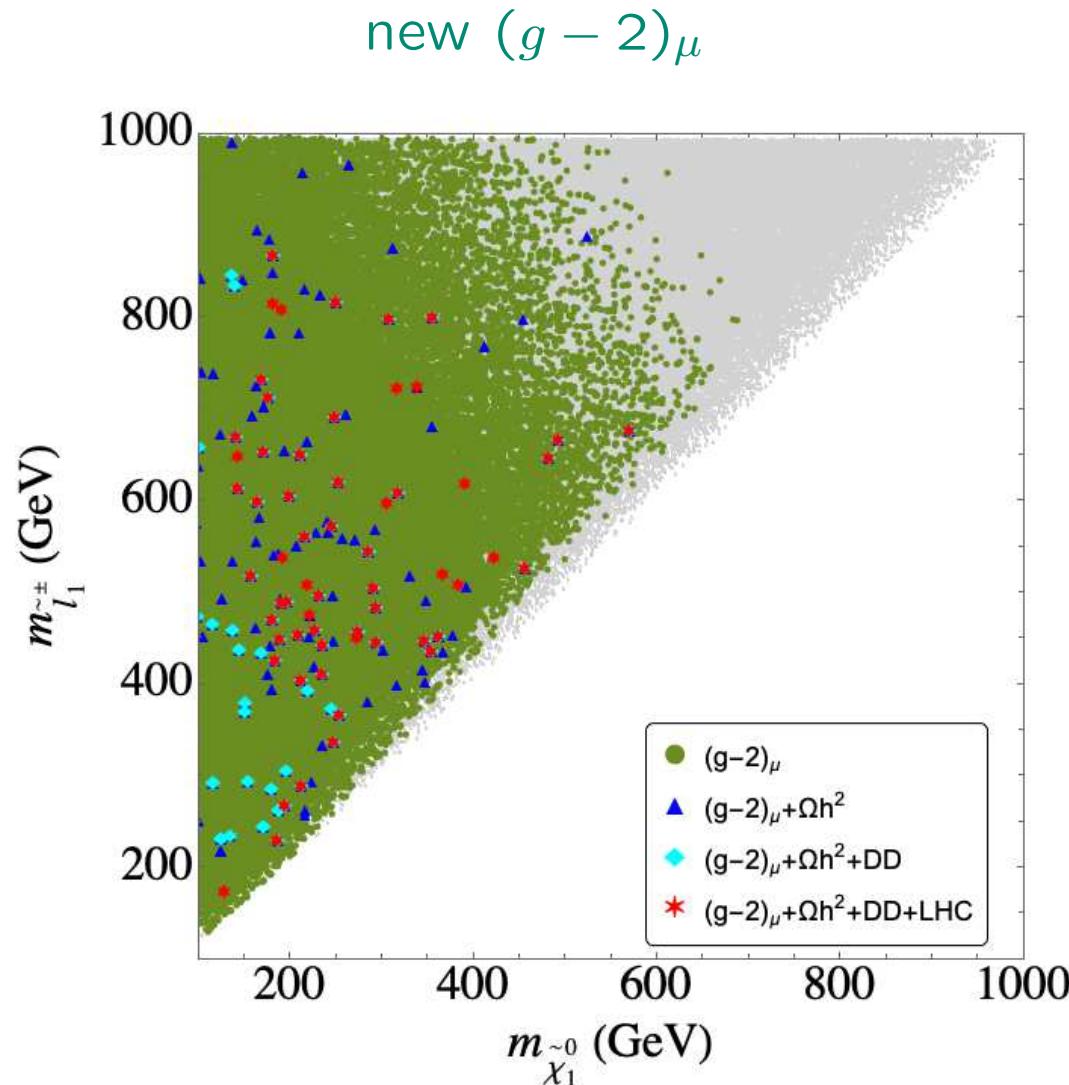
new $(g - 2)_\mu$



⇒ compressed spectrum as expected

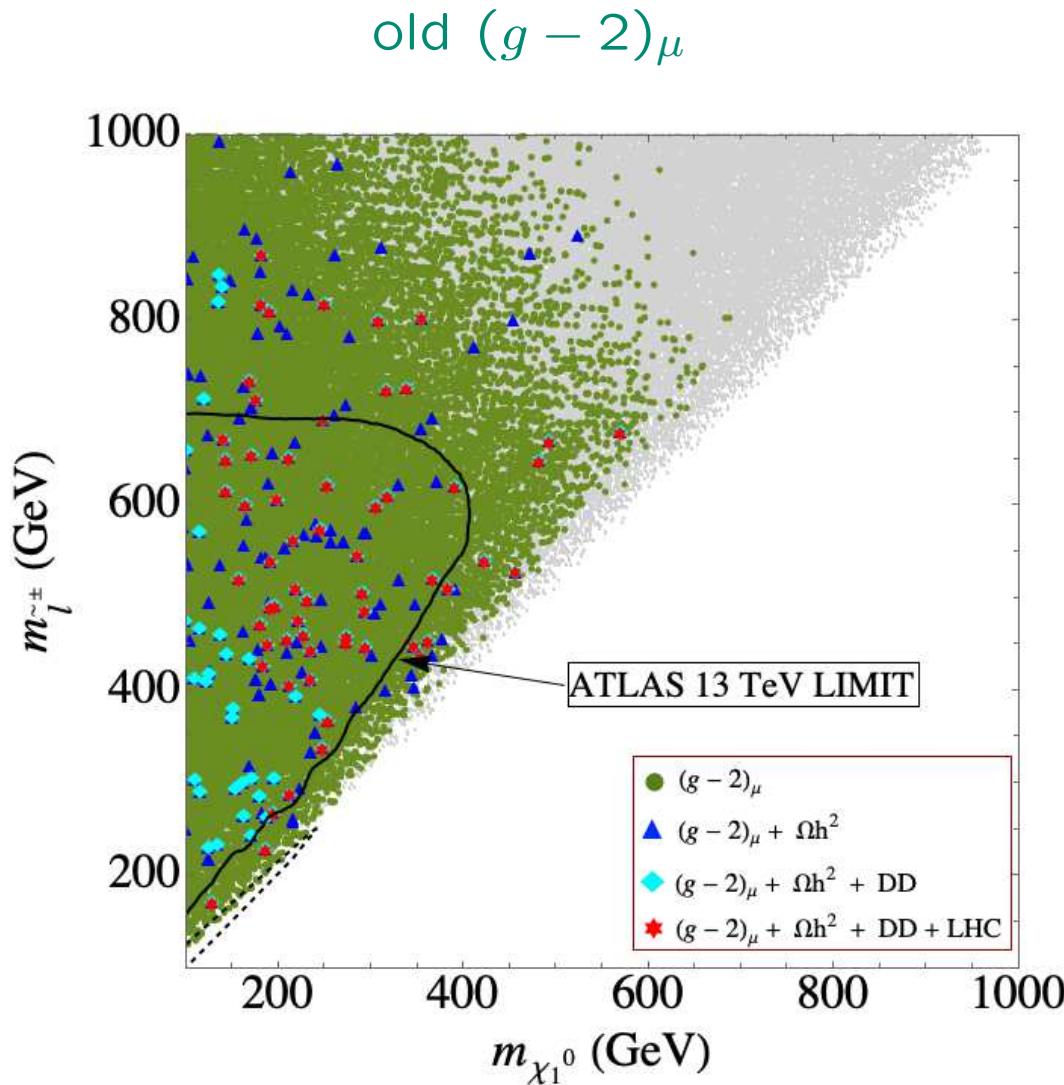
⇒ clear upper limits, $m_{(N)LSP} \lesssim 600(650)$ GeV confirmed

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:



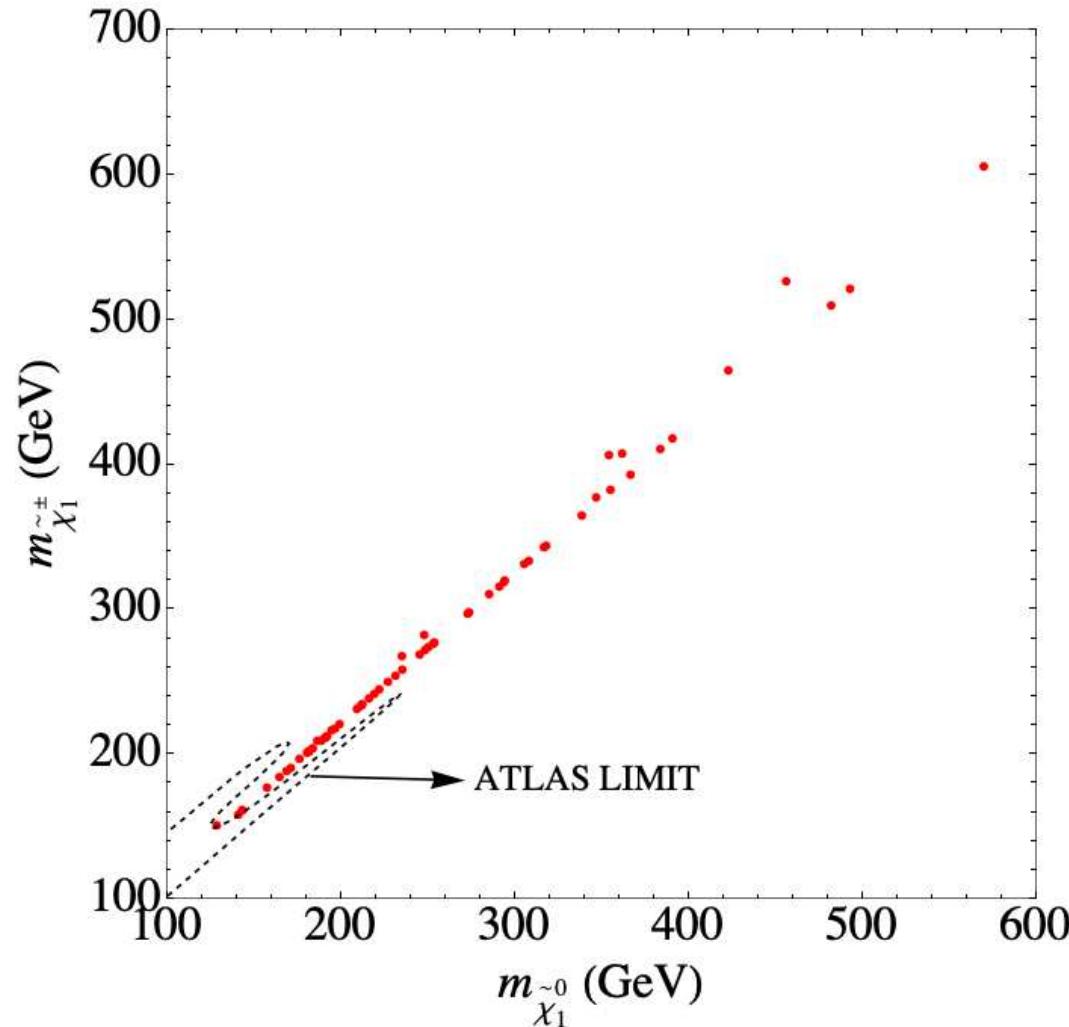
⇒ important: \tilde{l} -pair production searches (10)

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{t}_1}$ plane:



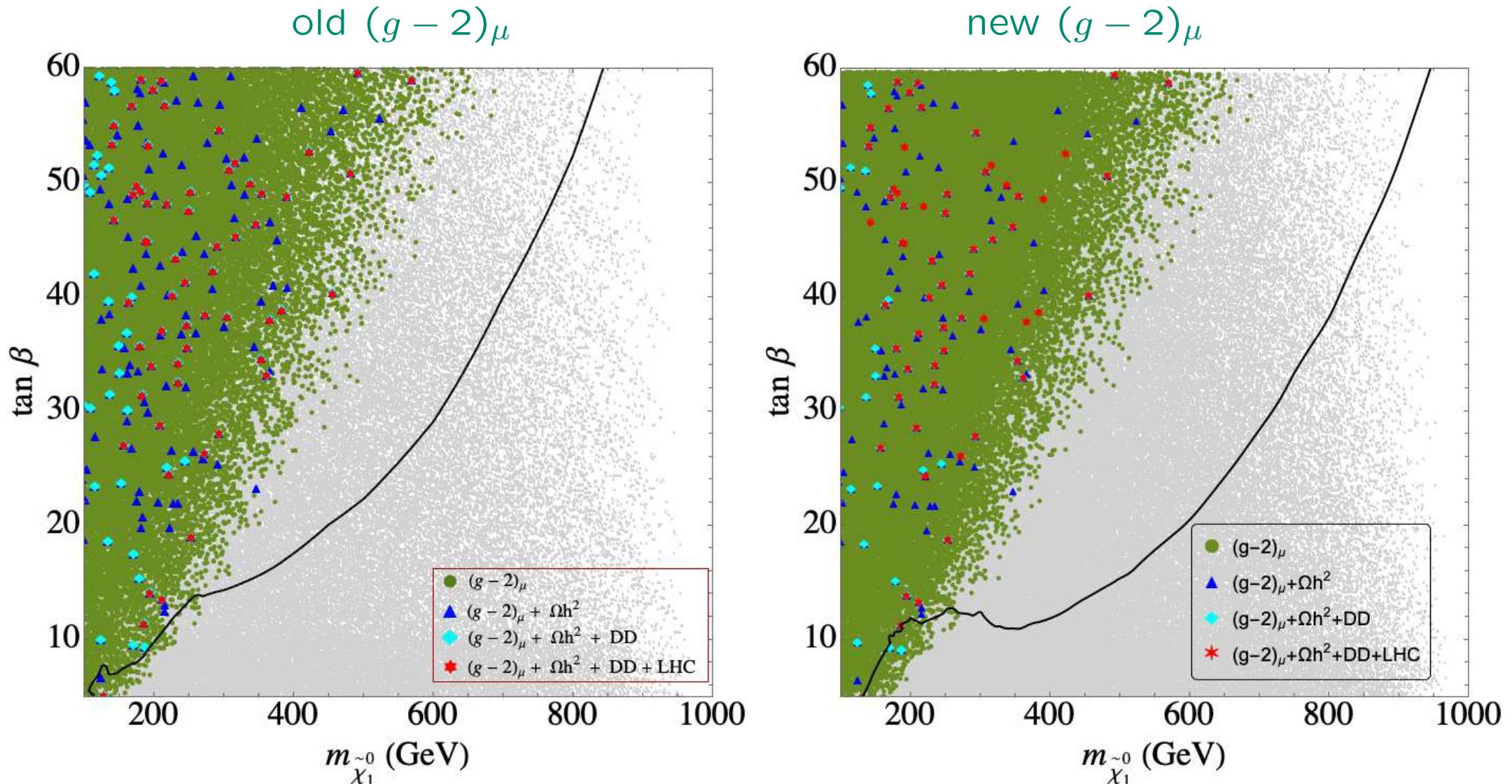
⇒ important: \tilde{t} -pair production searches (10)
⇒ naive application of LHC bounds fails

Comparison with the compressed spectra searches:



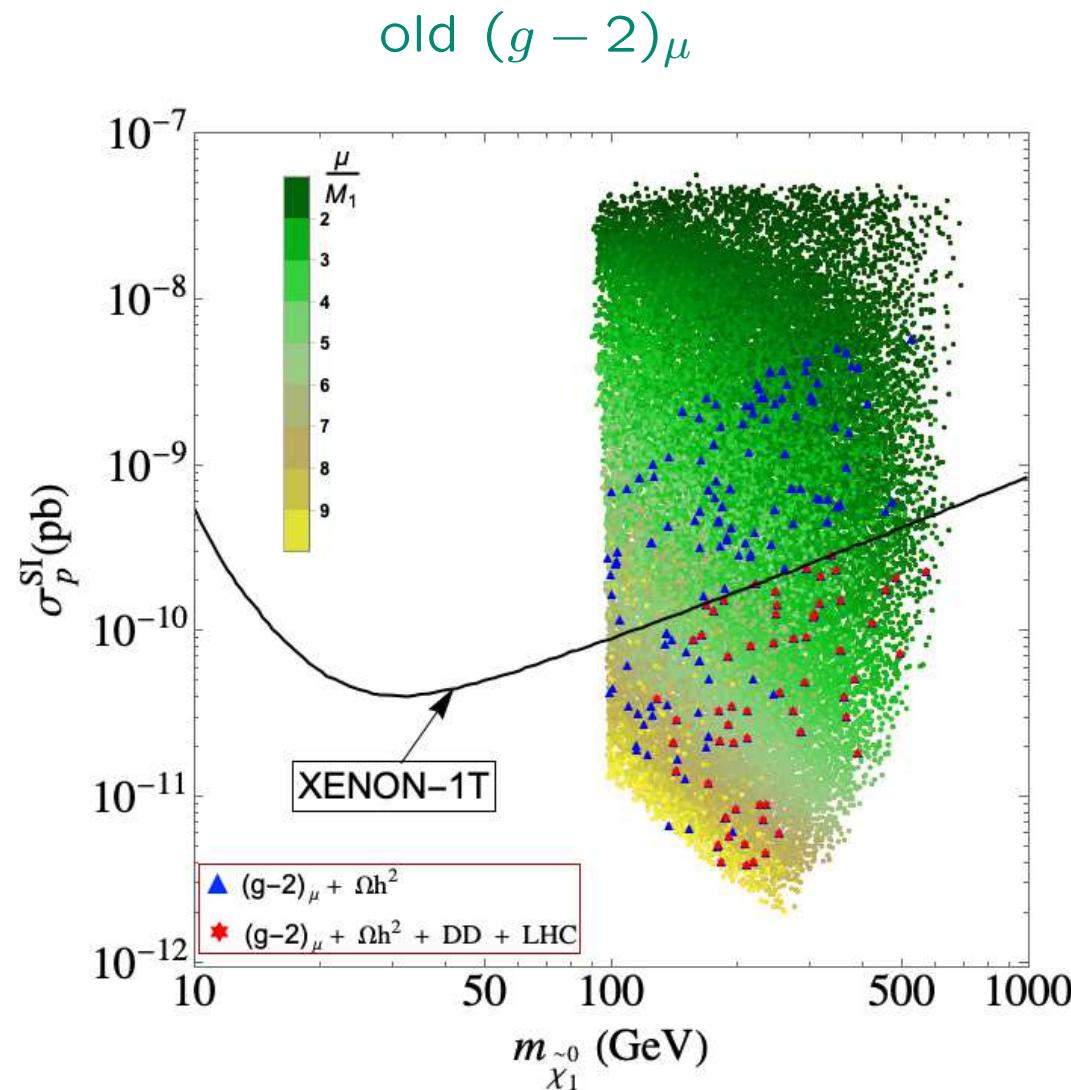
→ compressed spectrum avoids current bounds!

Results in the $m_{\tilde{\chi}_1^0}$ - $\tan \beta$ plane:



black contour: (simplified) application of $H/A \rightarrow \tau^+ \tau^-$
 $\rightarrow A\text{-pole annihilation effectively excluded}$

Results in the $m_{\tilde{\chi}_1^0}$ - σ_p^{SI} plane:



⇒ larger μ values favored

B/C) Bino DM with slepton co-annihilation

[*M. Chakraborti, S.H., I. Saha '20, '21*]

Parameter scan:

$$100 \text{ GeV} \leq M_1 \leq 1 \text{ TeV},$$

$$M_1 \leq M_2 \leq 10M_1,$$

$$1.1M_1 \leq \mu \leq 10M_1,$$

$$5 \leq \tan \beta \leq 60,$$

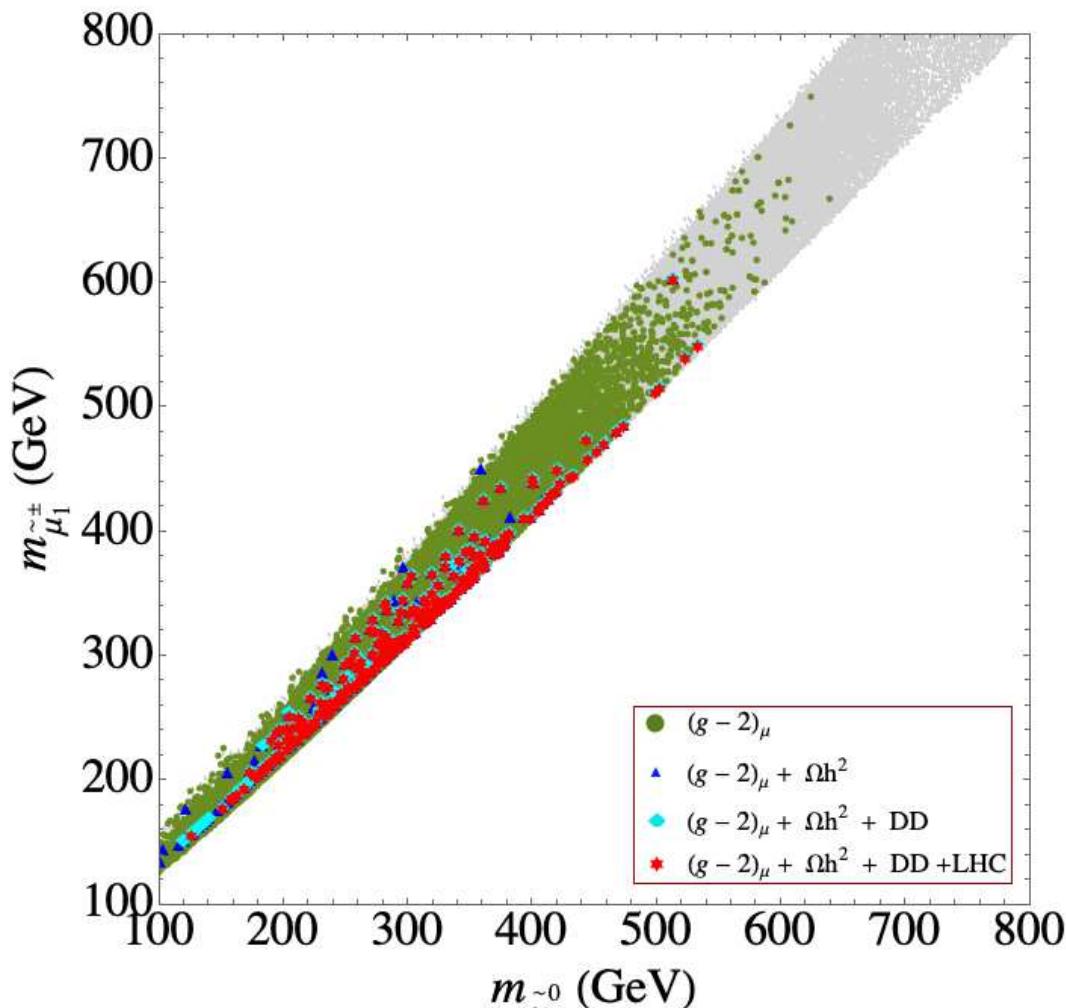
Case-L: $M_1 \leq m_{\tilde{L}} \leq 1.2M_1, \quad M_1 \leq m_{\tilde{R}} \leq 10M_1.$

Case-R: $M_1 \leq m_{\tilde{R}} \leq 1.2M_1, \quad M_1 \leq m_{\tilde{L}} \leq 10M_1.$

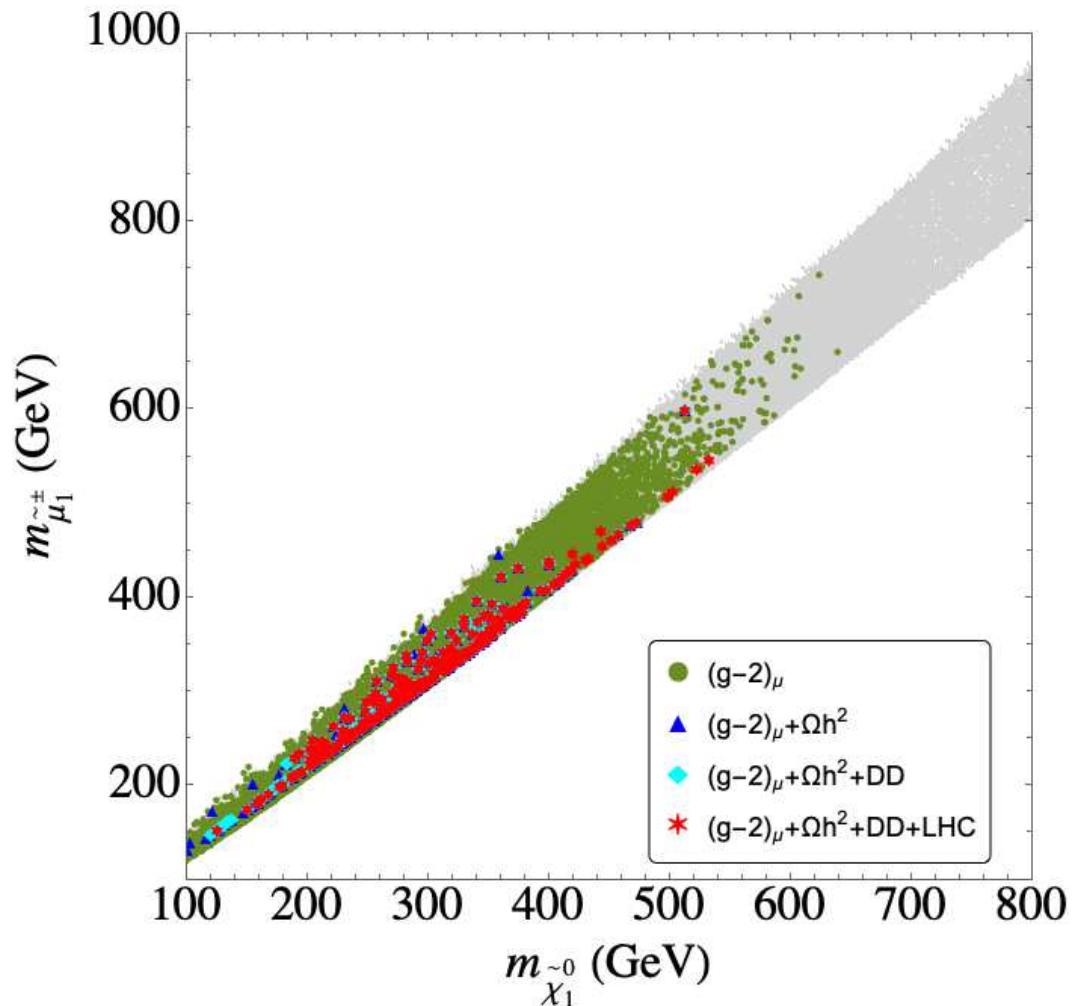
⇒ here focus on Case-L

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:

old $(g - 2)_\mu$



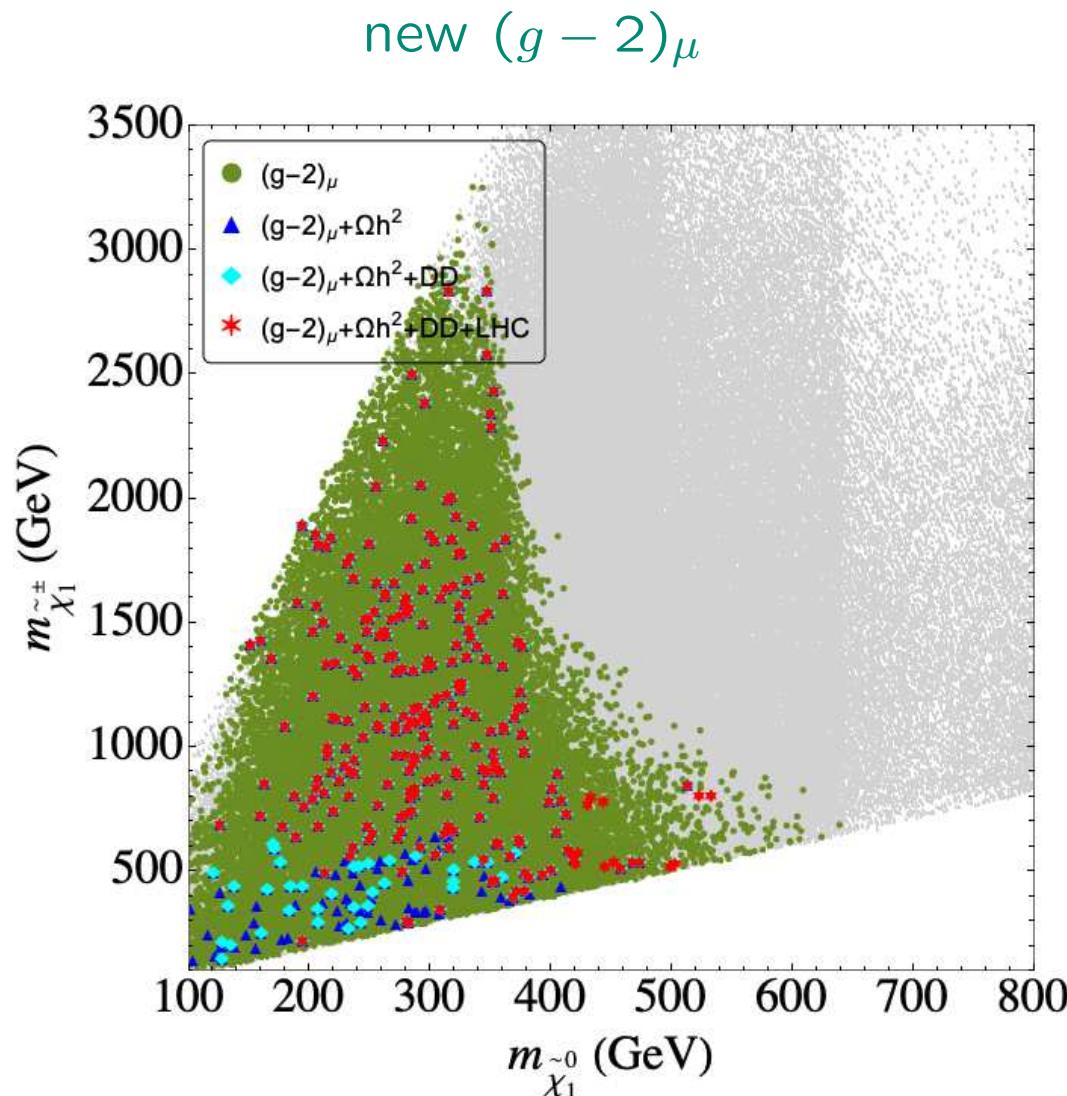
new $(g - 2)_\mu$



⇒ compressed spectrum as expected

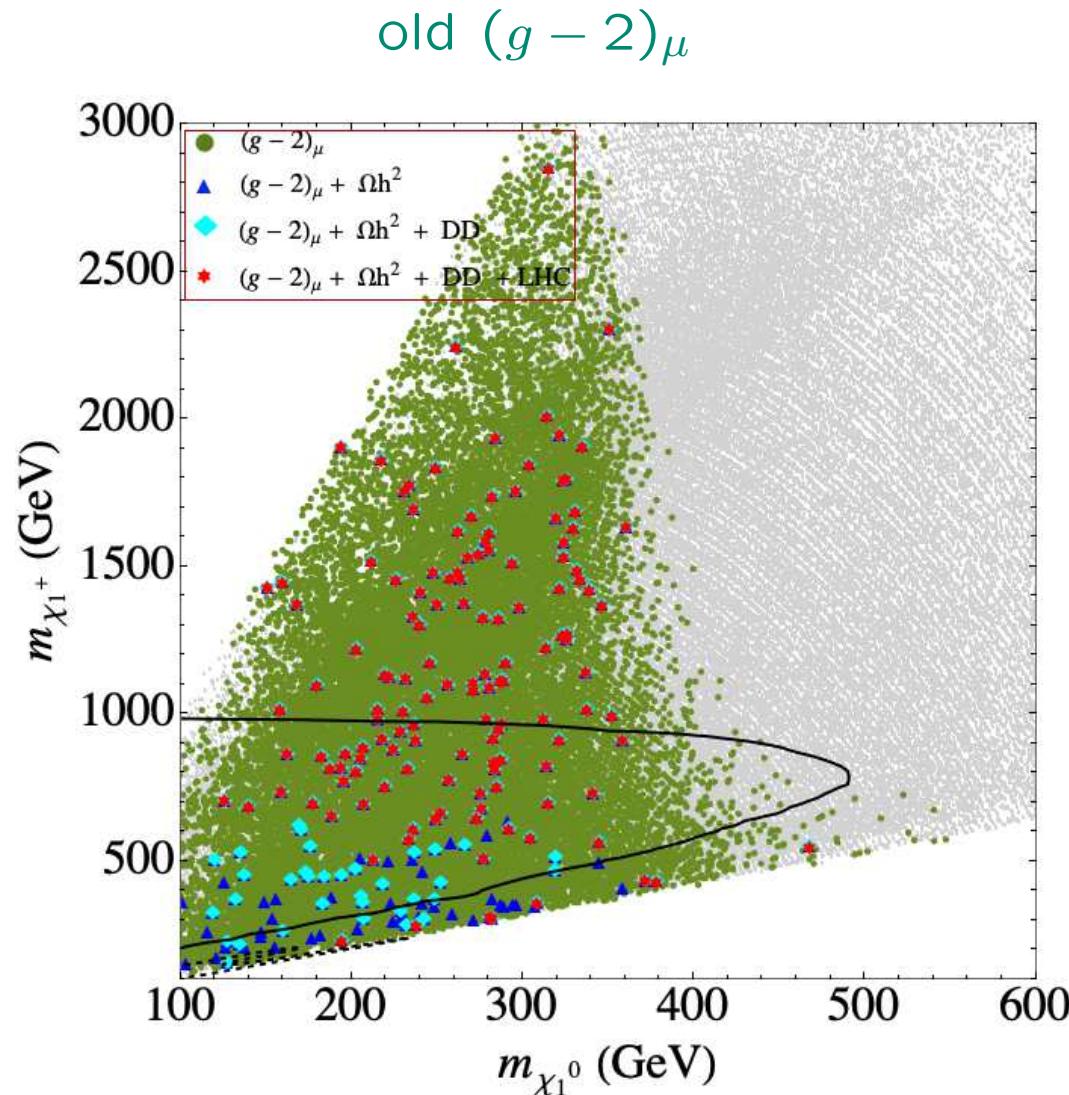
⇒ clear upper limits, $m_{(N)LSP} \lesssim 550(600)$ GeV confirmed

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{\chi}_1^\pm}$ plane:



⇒ important: $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production searches (5)

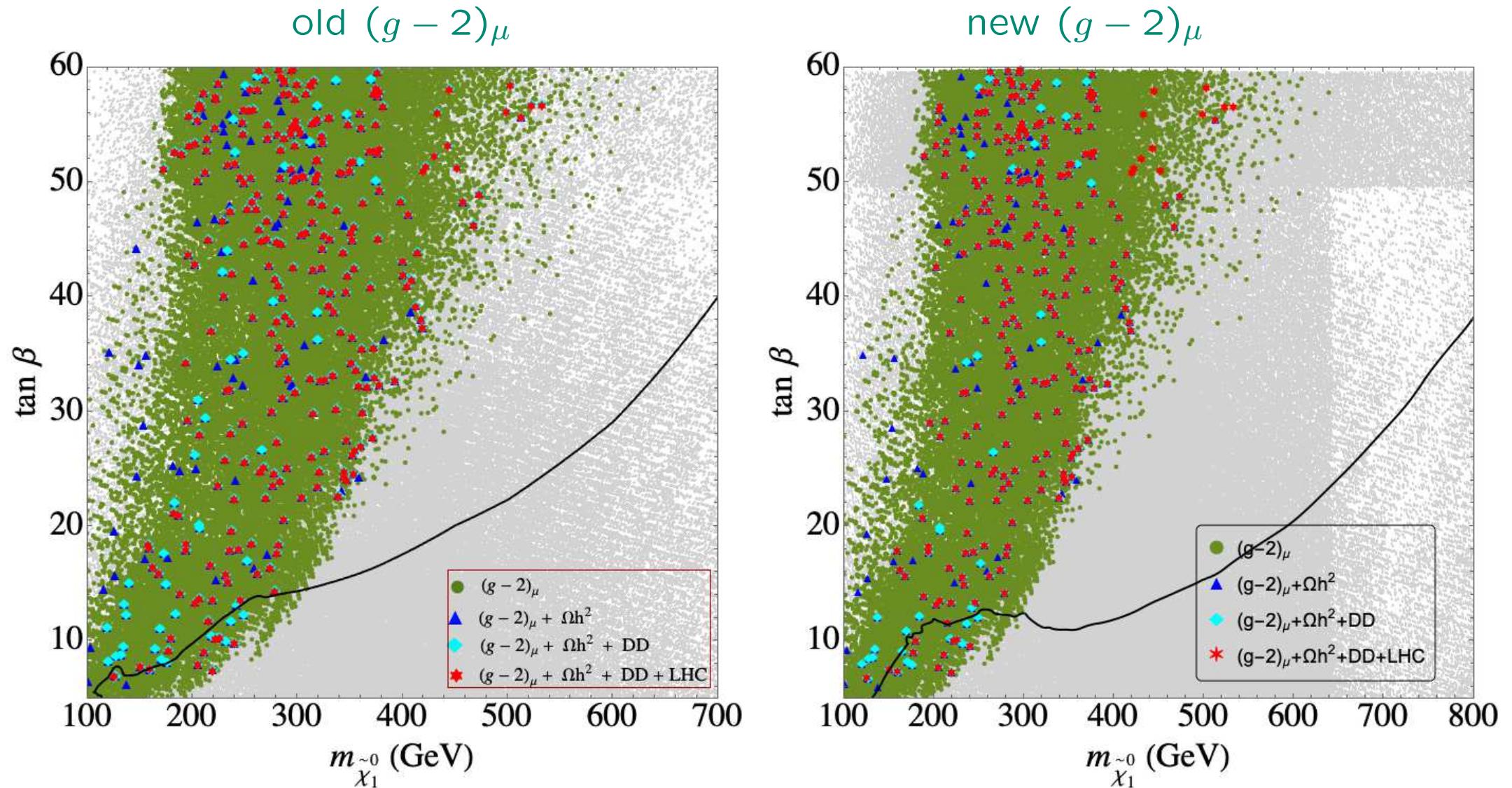
Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:



⇒ important: $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production searches (5)

⇒ naive application of LHC bounds fails

Results in the $m_{\tilde{\chi}_1^0}$ - $\tan \beta$ plane:



D) Higgsino DM

[*M. Chakraborti, S.H., I. Saha '21*]

Parameter scan:

$$100 \text{ GeV} \leq \mu \leq 1.2 \text{ TeV},$$

$$1.1\mu \leq M_1 \leq 10\mu,$$

$$1.1M_2 \leq \mu \leq 10\mu,$$

$$5 \leq \tan \beta \leq 60,$$

$$100 \text{ GeV} \leq m_{\tilde{L}}, m_{\tilde{R}} \leq 2 \text{ TeV},$$

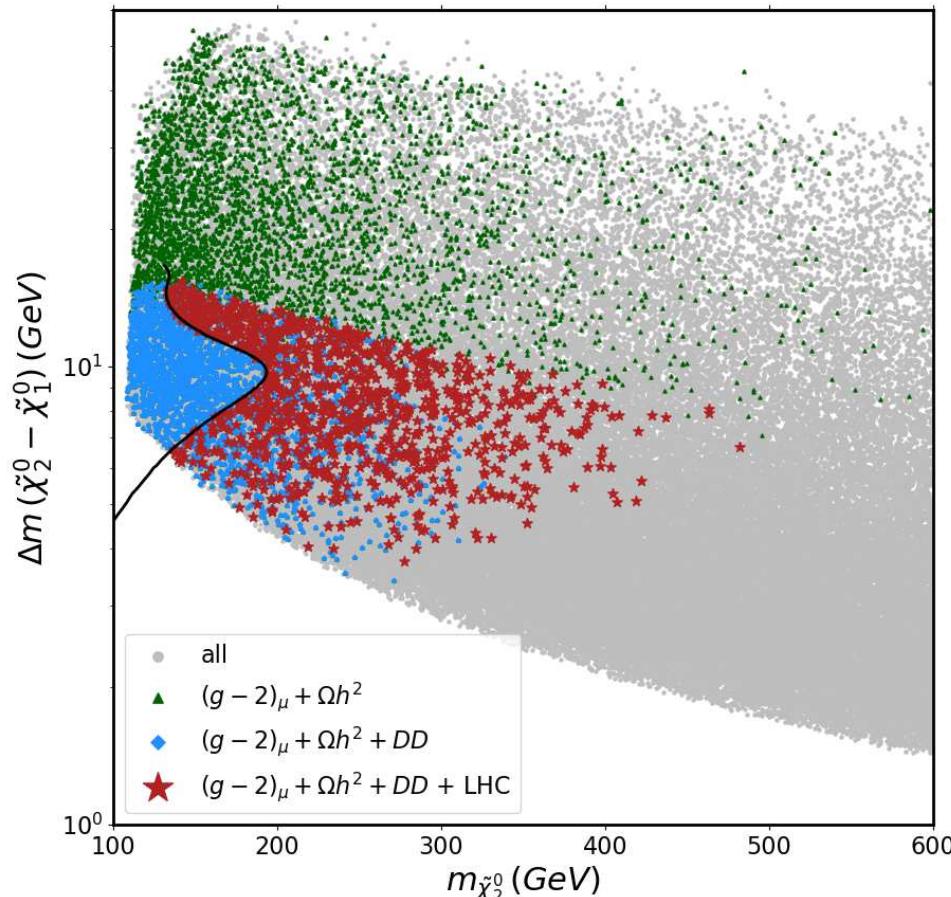
$$\Rightarrow m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_2^0} \sim m_{\tilde{\chi}_1^\pm} \sim \mu$$

Full DM relic density reached only for $m_{\tilde{\chi}_1^0} \sim 1 \text{ TeV}$

\Rightarrow incompatible with $(g - 2)_\mu$

Results in the $m_{\tilde{\chi}_2^0}$ - Δm plane:

old $(g - 2)_\mu$

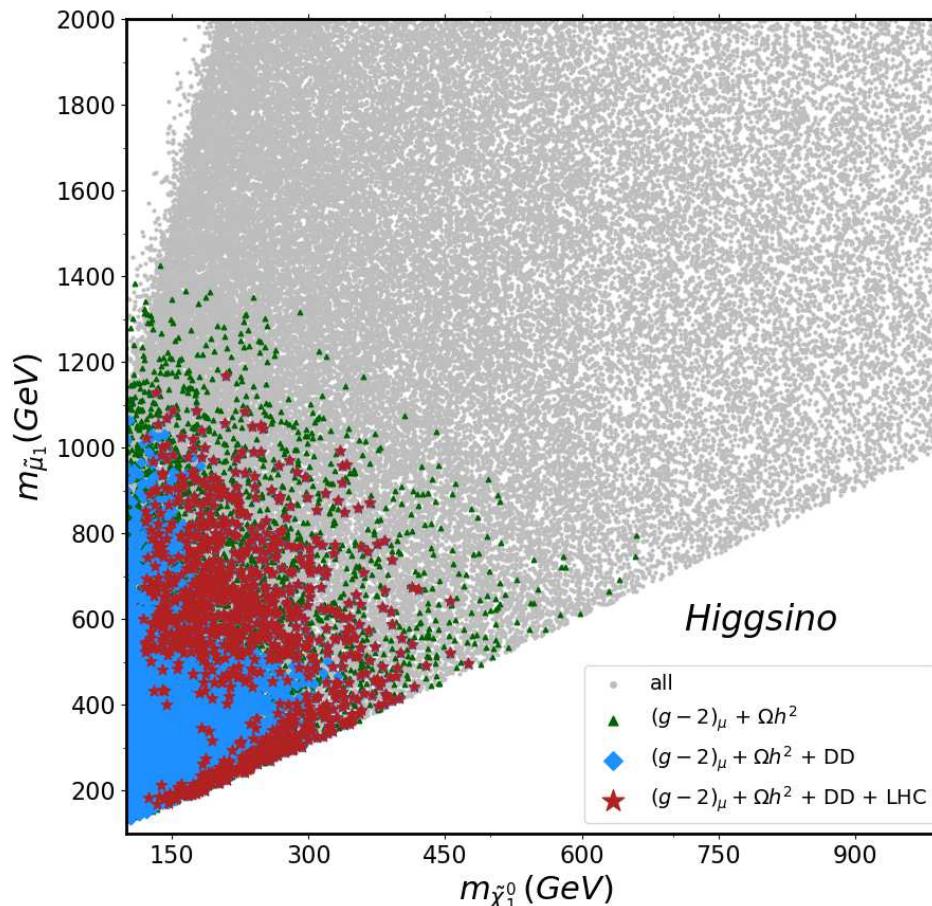


⇒ important: compressed spectra searches (11)

⇒ right where the model prediction sits ⇒ very powerful

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:

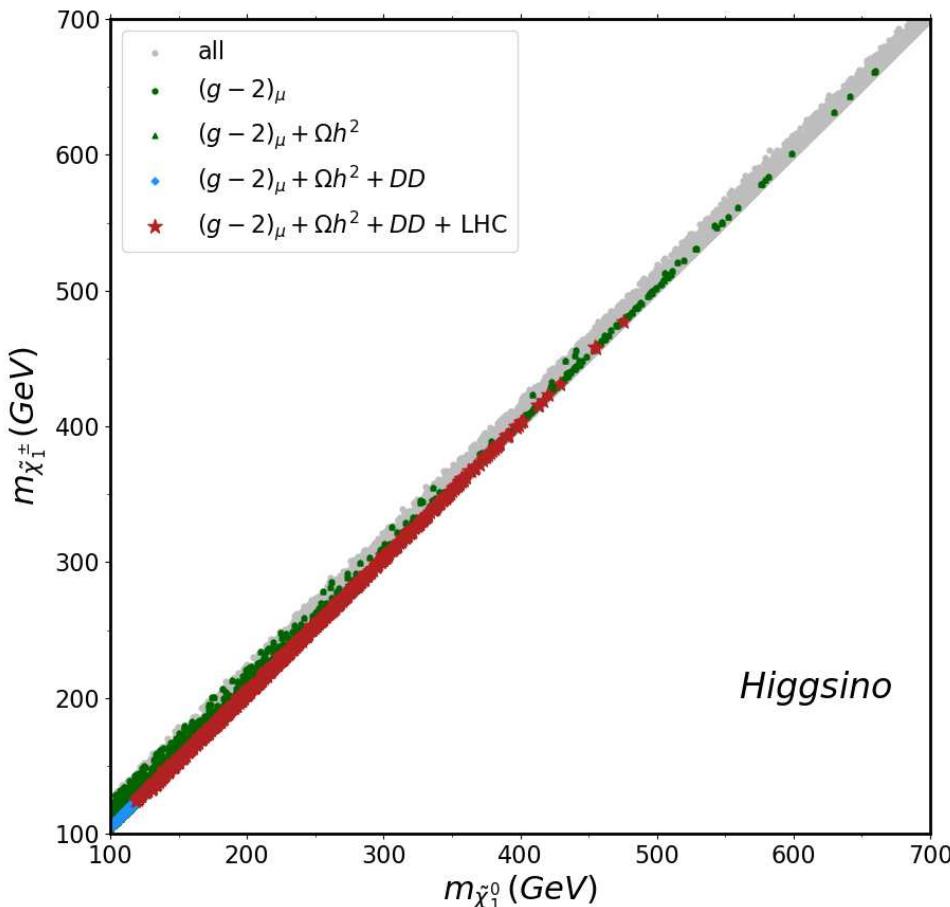
old $(g - 2)_\mu$



\Rightarrow upper limit on slepton masses: $m_{\tilde{l}_1} \lesssim 1.1$ TeV

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{\chi}_1^\pm}$ plane:

old $(g - 2)_\mu$

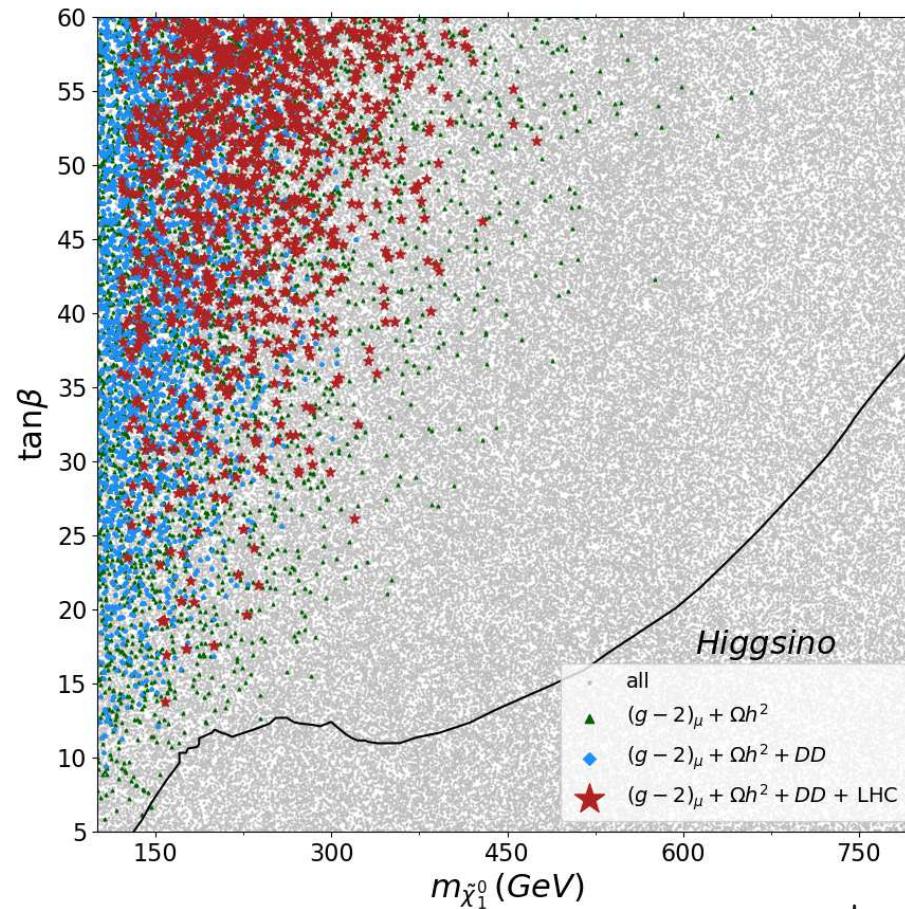


⇒ important: compressed spectra searches (11)

$\Rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0 \rightarrow (W^* \tilde{\chi}_1^0)(Z^* \tilde{\chi}_1^0) \rightarrow 2l + \cancel{E}_T + \text{ISR}$ $\Rightarrow m_{(\text{N})\text{LSP}} \lesssim 500 \text{ GeV}$

Results in the $m_{\tilde{\chi}_1^0}$ - $\tan\beta$ plane:

old $(g - 2)_\mu$



black contour: (simplified) application of $H/A \rightarrow \tau^+ \tau^-$
 $\Rightarrow A$ -pole annihilation fully excluded

E) Wino DM

[*M. Chakraborti, S.H., I. Saha '20, '21*]

Parameter scan:

$$100 \text{ GeV} \leq M_2 \leq 1.5 \text{ TeV},$$

$$1.1M_2 \leq M_1 \leq 10M_2,$$

$$1.1M_2 \leq \mu \leq 10M_2,$$

$$5 \leq \tan \beta \leq 60,$$

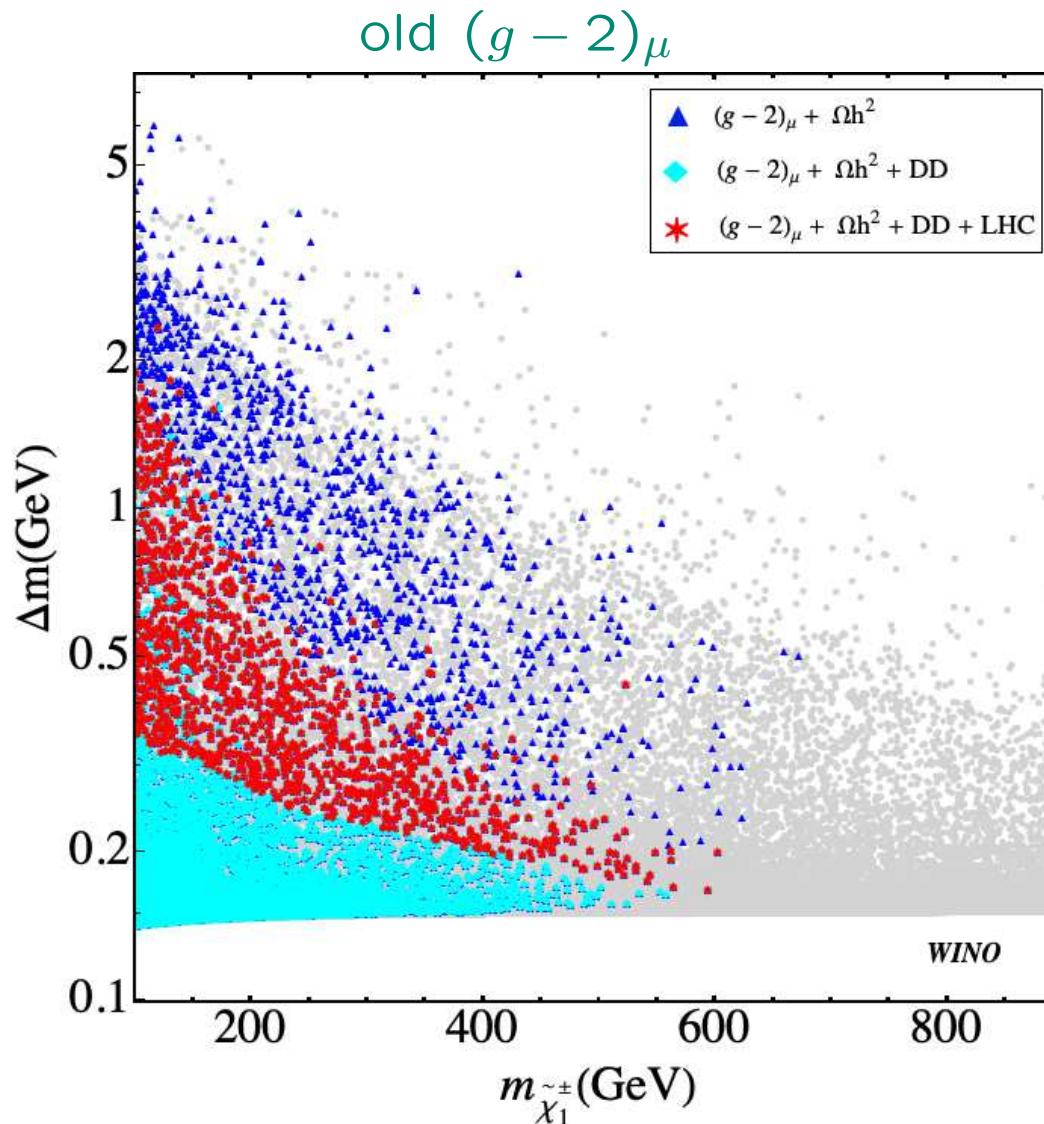
$$100 \text{ GeV} \leq m_{\tilde{L}}, m_{\tilde{R}} \leq 2 \text{ TeV},$$

$$\Rightarrow m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_1^\pm} \sim M_2$$

Full DM relic density reached only for $m_{\tilde{\chi}_1^0} \sim 3 \text{ TeV}$

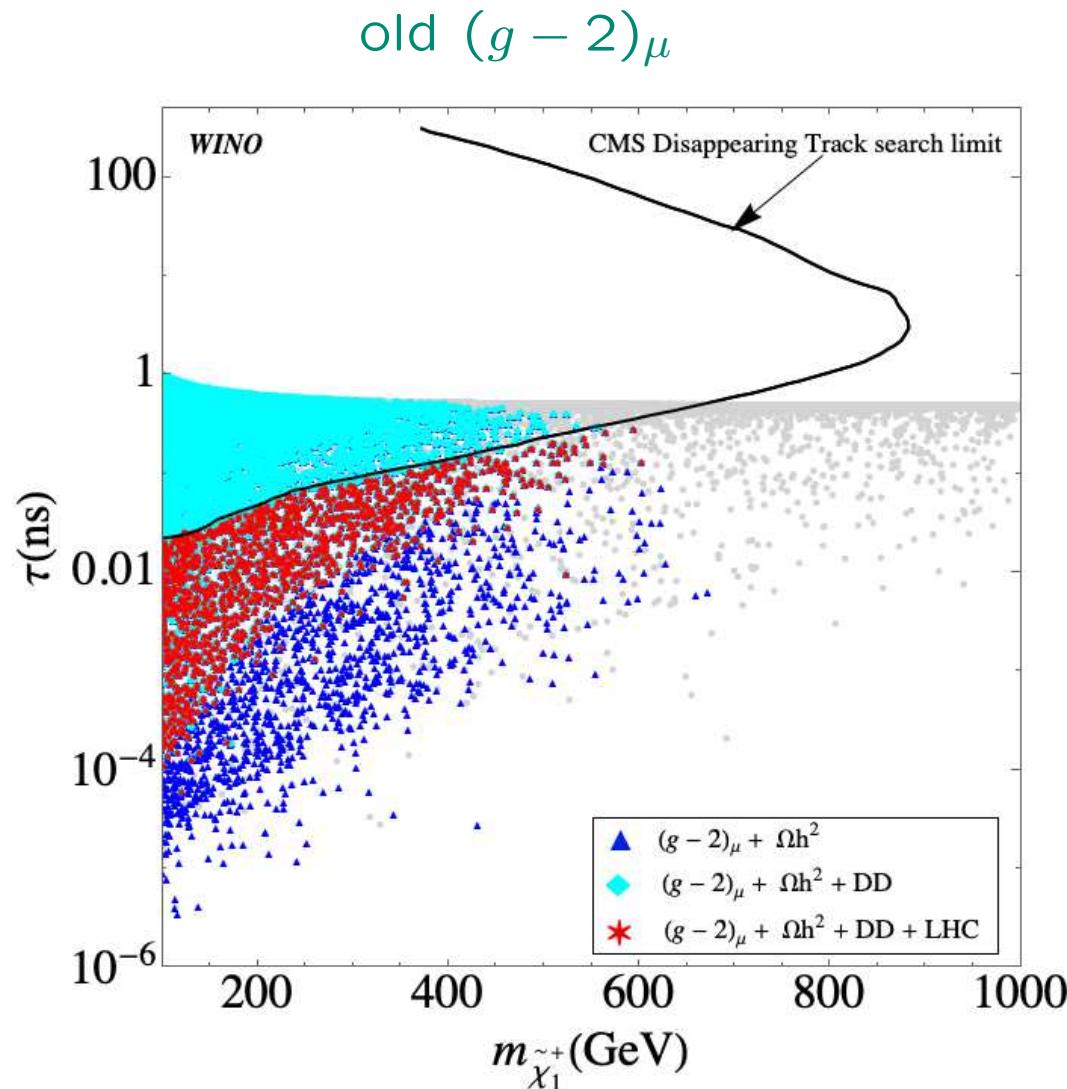
\Rightarrow incompatible with $(g - 2)_\mu$

Results in the $m_{\tilde{\chi}_1^\pm}$ - Δm plane:



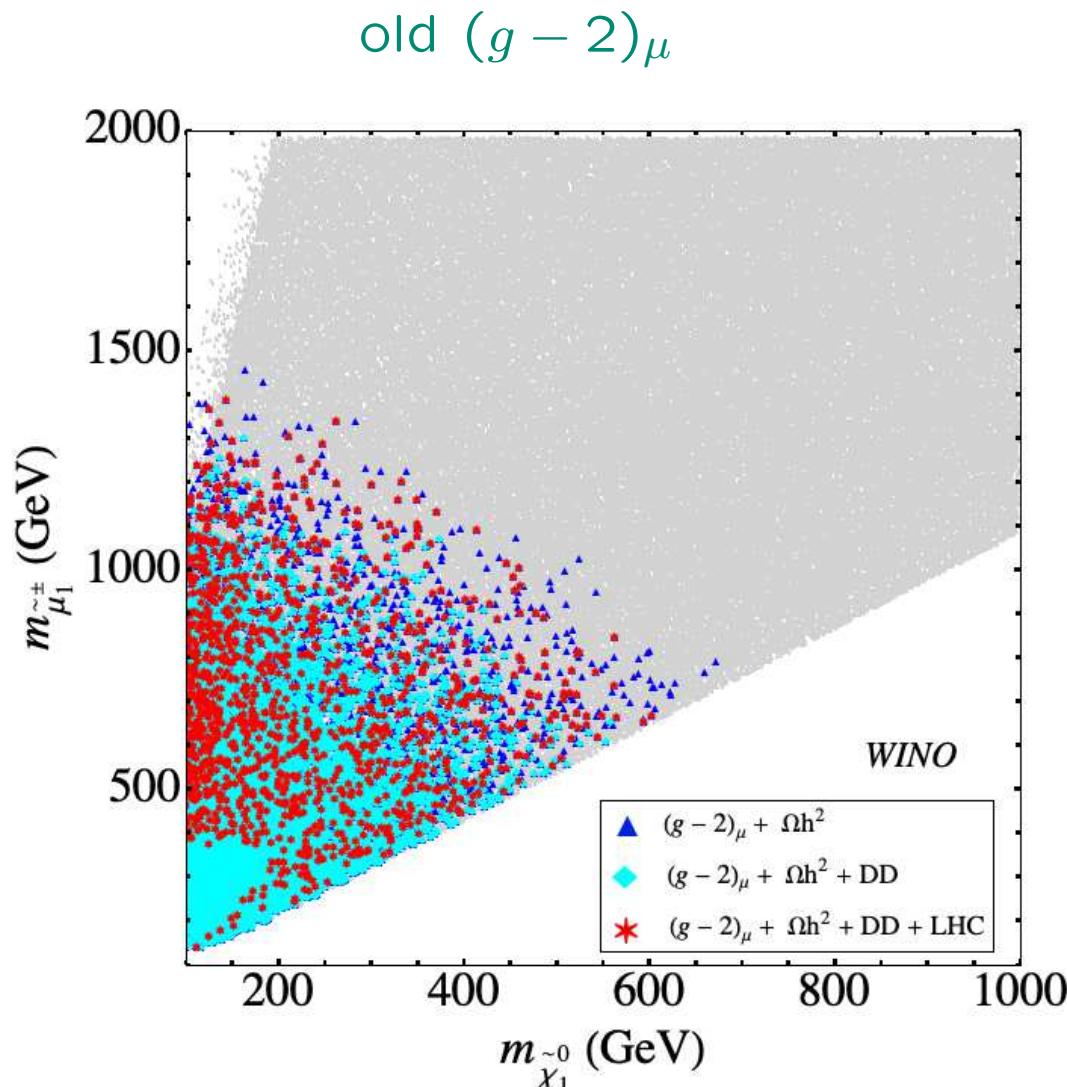
→ important: disappearing track limit $\Rightarrow m_{(N)LSP} \lesssim 600$ GeV
 → allowed parameter space squeezed by DD limits and disapp. tracks

Results in the $m_{\tilde{\chi}_1^\pm}$ - $\tau_{\tilde{\chi}_1^\pm}$ plane:



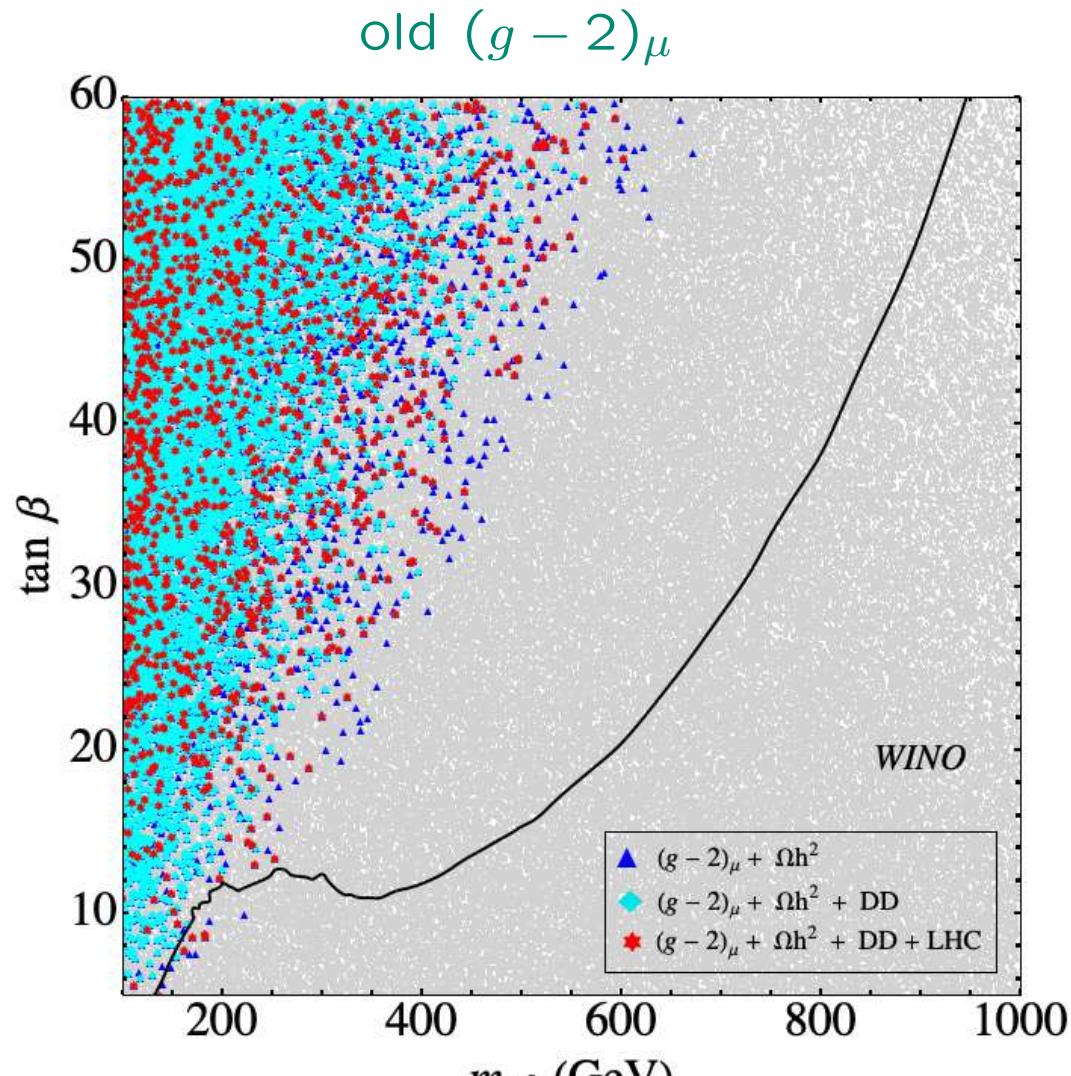
→ important: disappearing track limit $\Rightarrow m_{(N)LSP} \lesssim 600$ GeV
 → allowed parameter space squeezed by DD limits and disapp. tracks

Results in the $m_{\tilde{\chi}_1^0}$ - $m_{\tilde{l}_1}$ plane:



$\Rightarrow m_{\tilde{l}_1} \lesssim 1400(1200)$ GeV

Results in the $m_{\tilde{\chi}_1^0}$ - $\tan \beta$ plane:

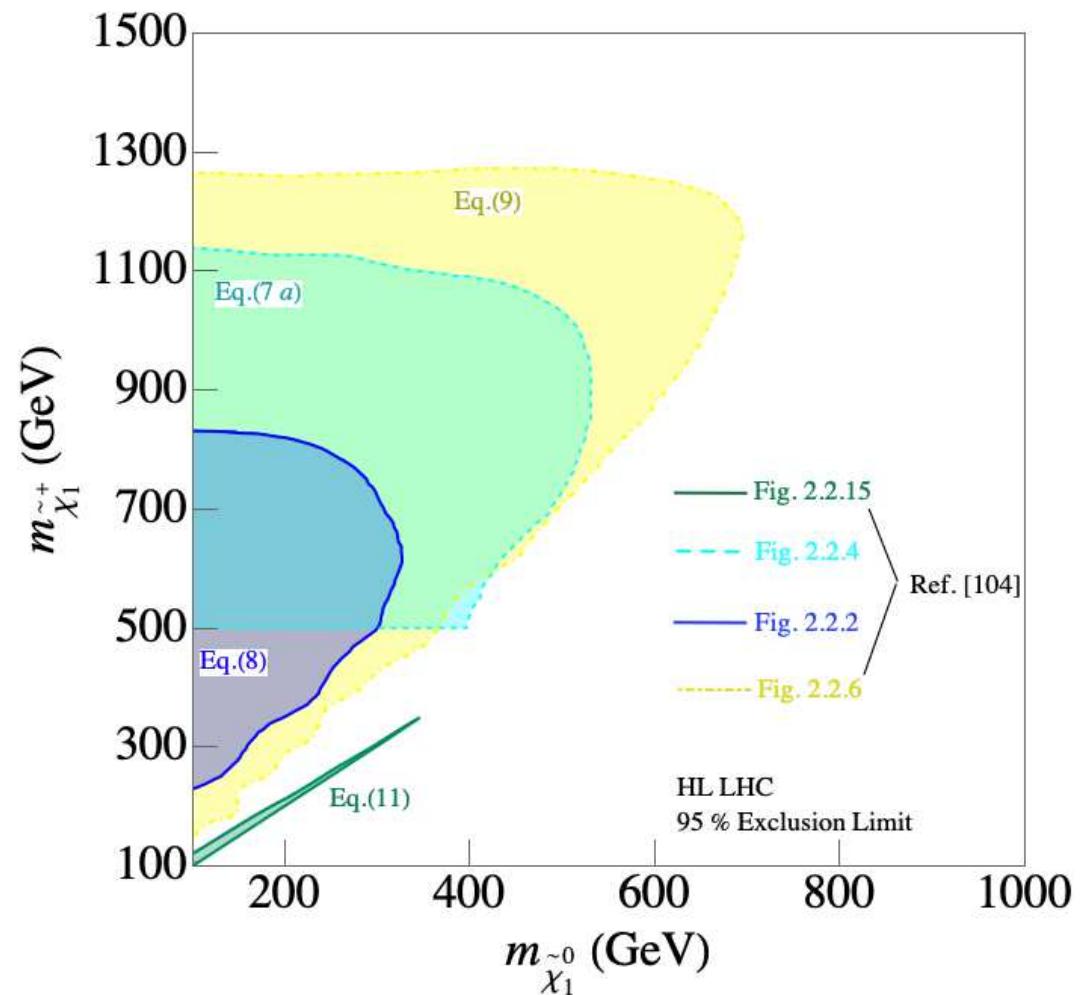
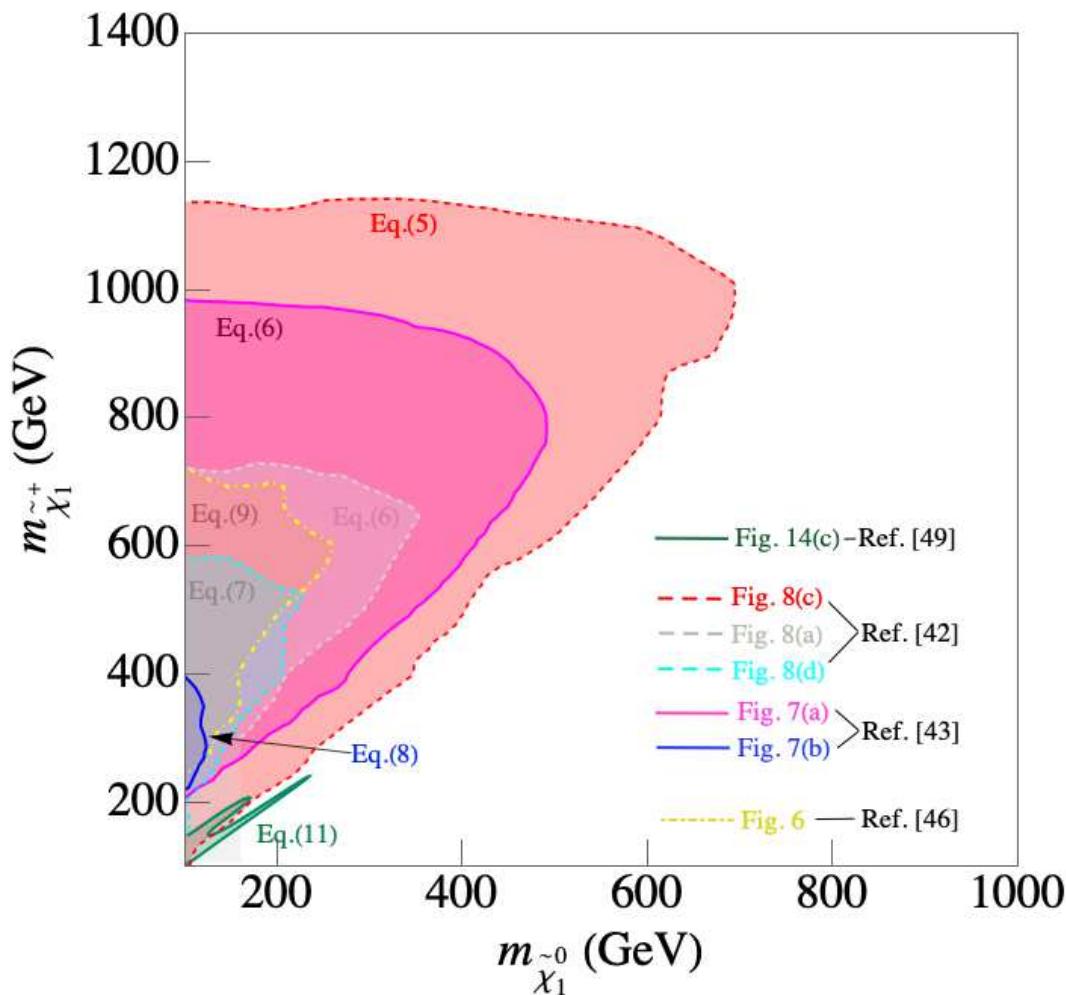


black contour: (simplified) application of $H/A \rightarrow \tau^+ \tau^-$
 $\Rightarrow A$ -pole annihilation largely excluded

LHC exclusion bounds vs. HL-LHC exclusion bounds

not all channels available

[YR18]



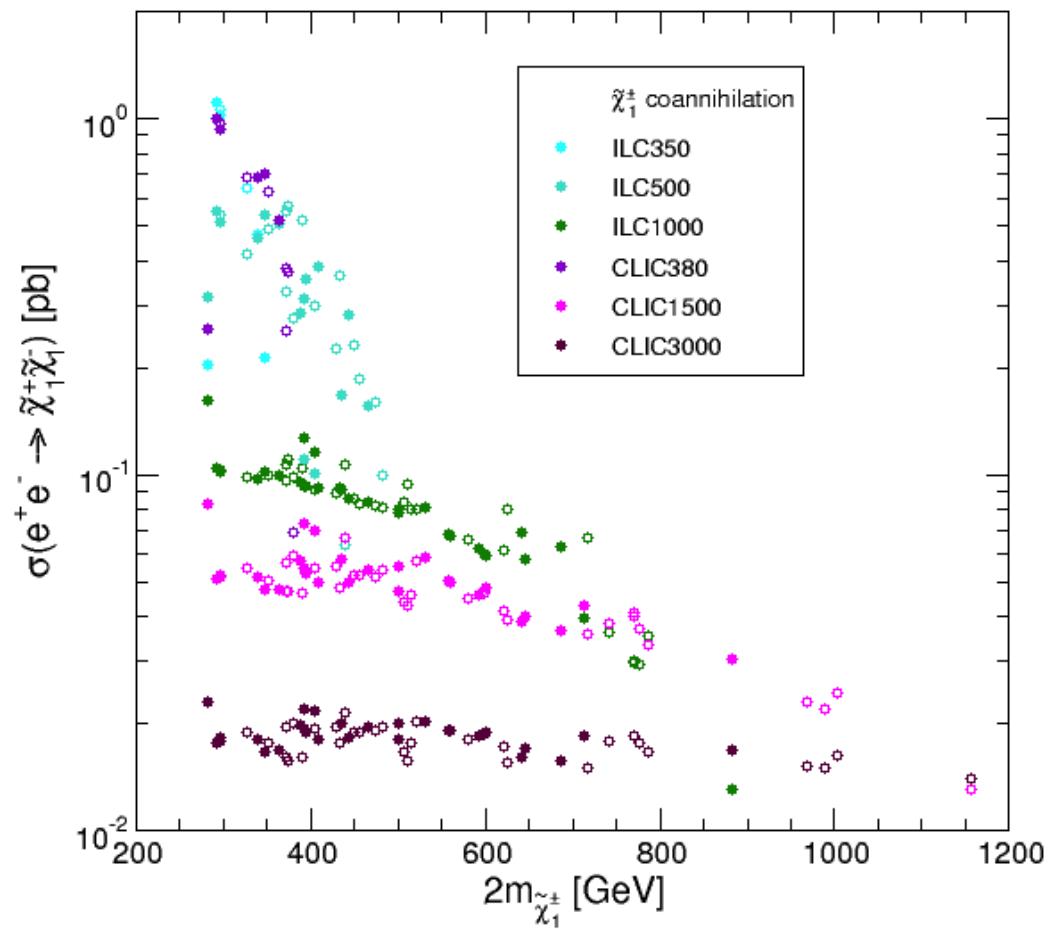
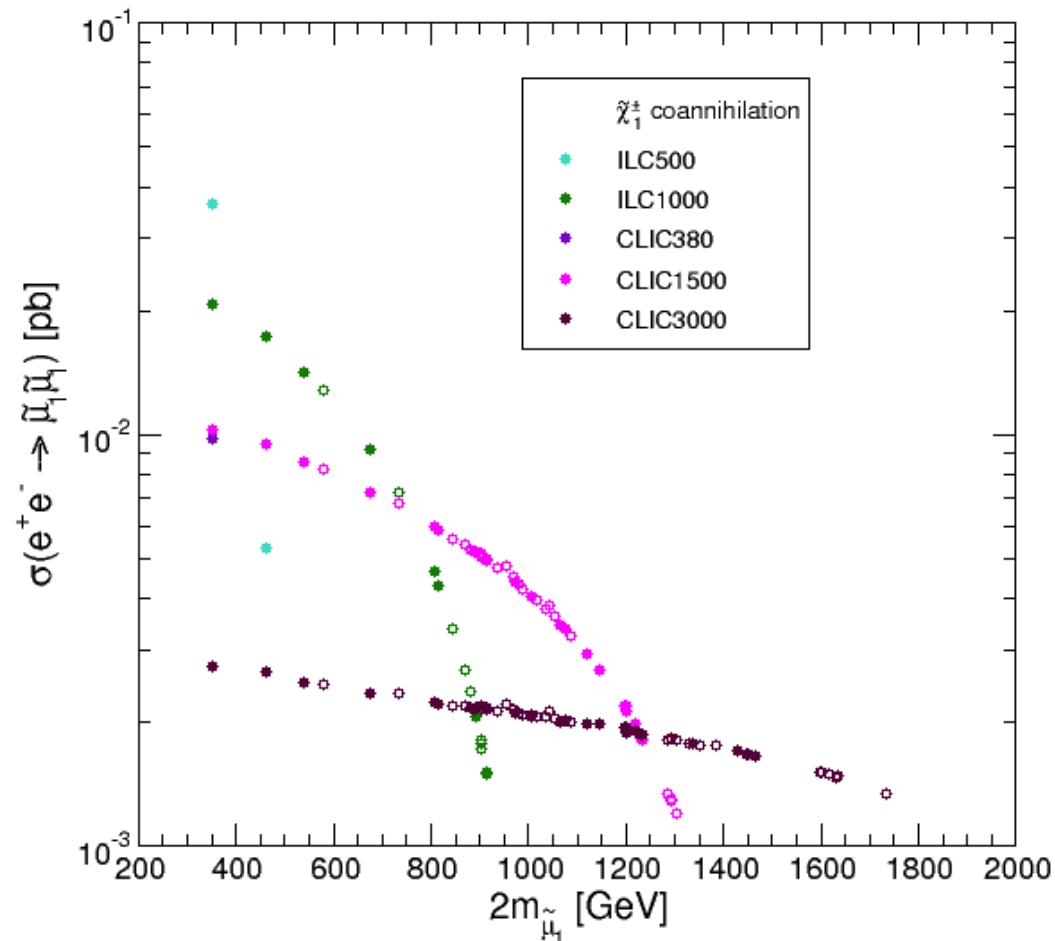
⇒ exclusion reach can be important

⇒ no CheckMate inclusion available ...

Direct production at e^+e^- colliders (ILC/CLIC)

wino/bino DM with chargino co-ann.

(open/full: "old" $(g - 2)_\mu$)



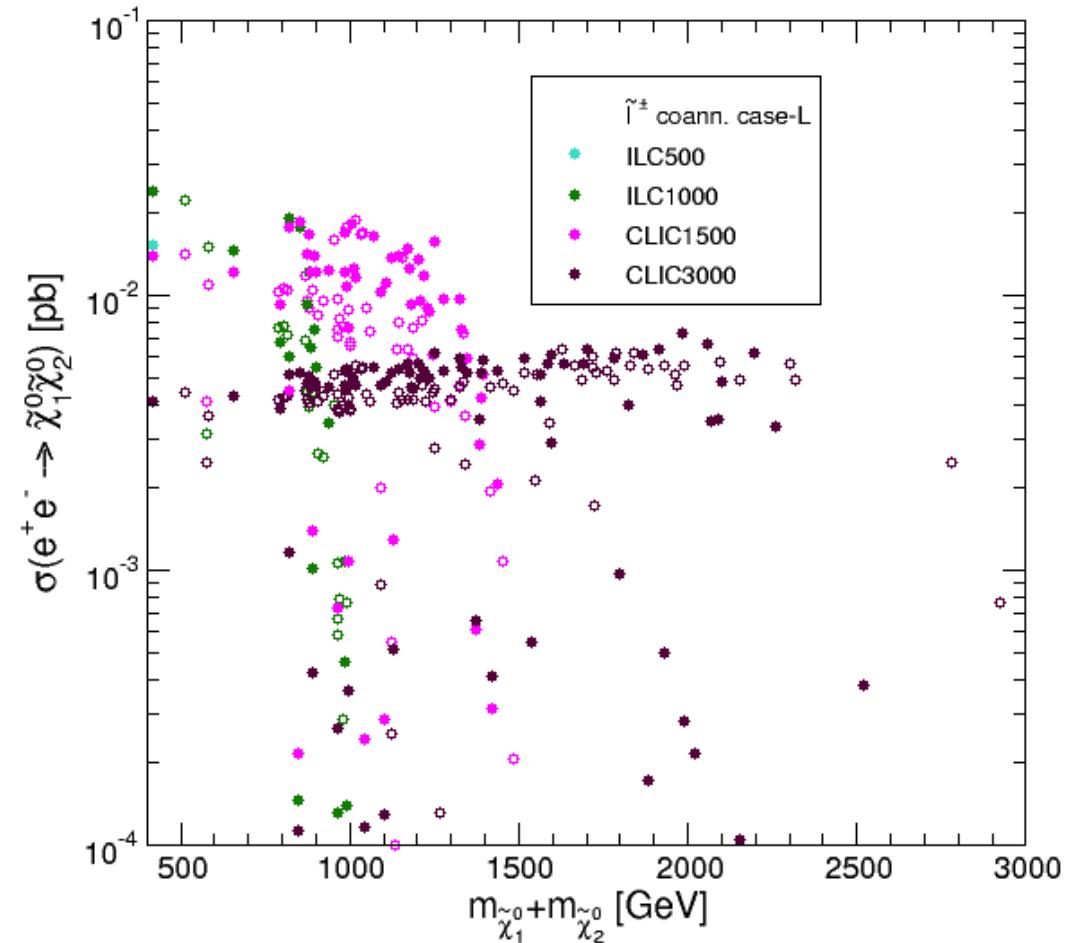
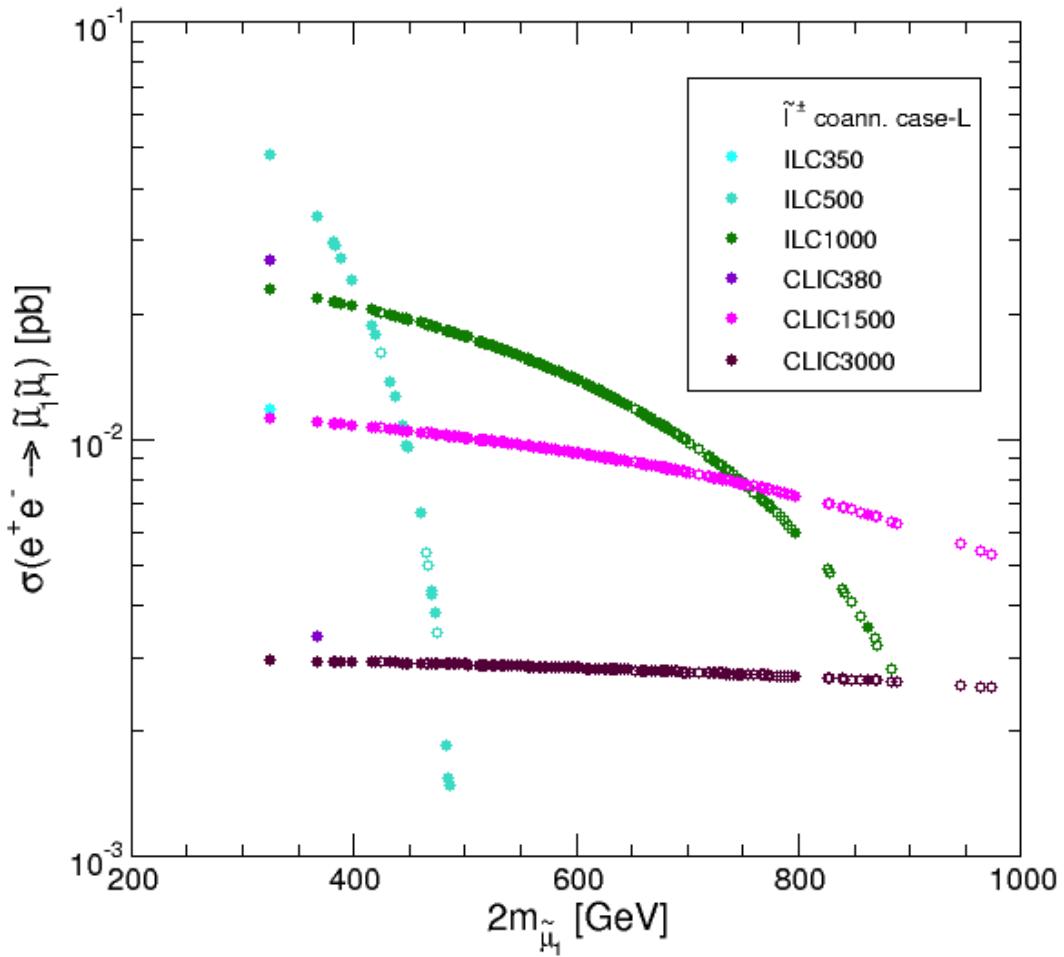
→ ILC has good prospects (particularly for $\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$)

→ CLIC can cover everything

Direct production at e^+e^- colliders (ILC/CLIC)

bino DM with slepton co-ann.

(open/full: "old" $(g - 2)_\mu$)



→ ILC can nearly covers full smuon channel (but not $\tilde{\chi}_1^0 \tilde{\chi}_2^0$)

→ CLIC can cover everything