



Science and  
Technology  
Facilities Council

# Results and prospects from the Experiment



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**RAL, UK**

On behalf of the T2K Collaboration

XX Lomonosov Conference on Elementary Particle Physics  
20 August 2021, Online

# Neutrino oscillations

- Mixing between distinct **flavour eigenstates** (interacting states) and **mass eigenstates** (propagating states), as governed by the **Pontecorvo-Maki-Nakagawa-Sakata matrix**



Bruno Pontecorvo, 1957

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

**Atmospheric & Accelerator**      **Accelerator & Reactor**      **Reactor & Solar**

$\theta_{23} \approx 48^\circ$        $\theta_{13} \approx 8^\circ$        $\theta_{12} \approx 34^\circ$   
 $|\Delta m_{32}^2| \approx |\Delta m_{31}^2| \approx 2.5 \times 10^{-3} \text{ eV}^2$        $\delta_{CP} = ?$        $\Delta m_{21}^2 \approx 7.5 \times 10^{-5} \text{ eV}^2$

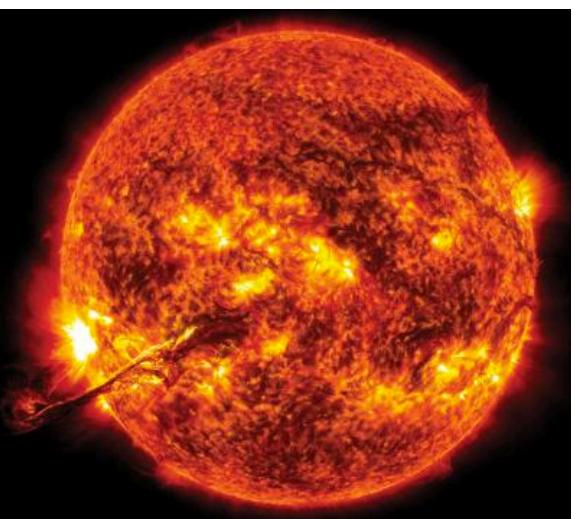
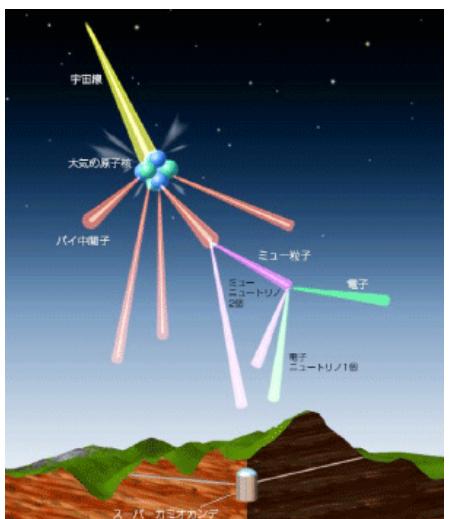
$s_{ij} = \sin \theta_{ij}$   
 $c_{ij} = \cos \theta_{ij}$

B. Pontecorvo, “Mesonium and antimesonium,”

*Zhurnal Eksperimental'noi i Teoreticheskoi Fiziki*, vol. 33, p. 549, 1957

B. Pontecorvo, “Inverse beta processes and nonconservation of lepton charge,”

*Zhurnal Eksperimental'noi i Teoreticheskoi Fiziki*, vol. 34, p. 247, 1957



# Open questions in neutrino oscillation physics

- **Neutrino mass ordering?**
  - Normal ( $\Delta m_{32}^2 > 0$ ) or inverted ( $\Delta m_{32}^2 < 0$ )
- **$\theta_{23}$  octant?**
  - Upper ( $\sin^2 \theta_{23} > 0.5$ ) or lower ( $\sin^2 \theta_{23} < 0.5$ )
  - $\theta_{23} = 45^\circ$  would indicate an underlying  $\mu - \tau$  symmetry of the PMNS matrix
- **CP violation in neutrino sector?**
  - Would like to measure with the same precision as CP violation in hadronic weak interactions
  - $\delta_{CP} \neq 0$ ,  $\pi$  would indicate  $P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$  in vacuum
- T2K is well positioned to answer all three questions!

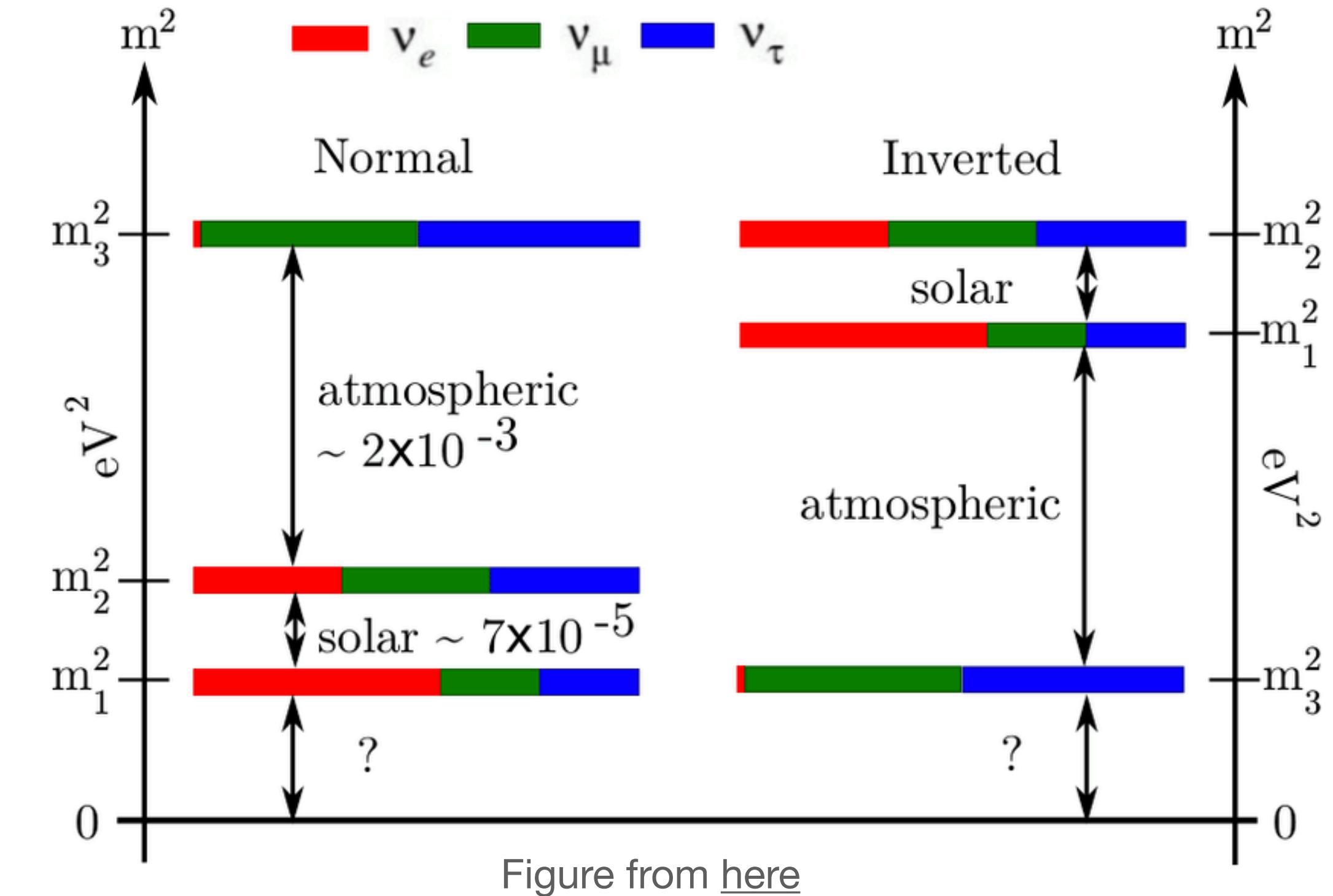
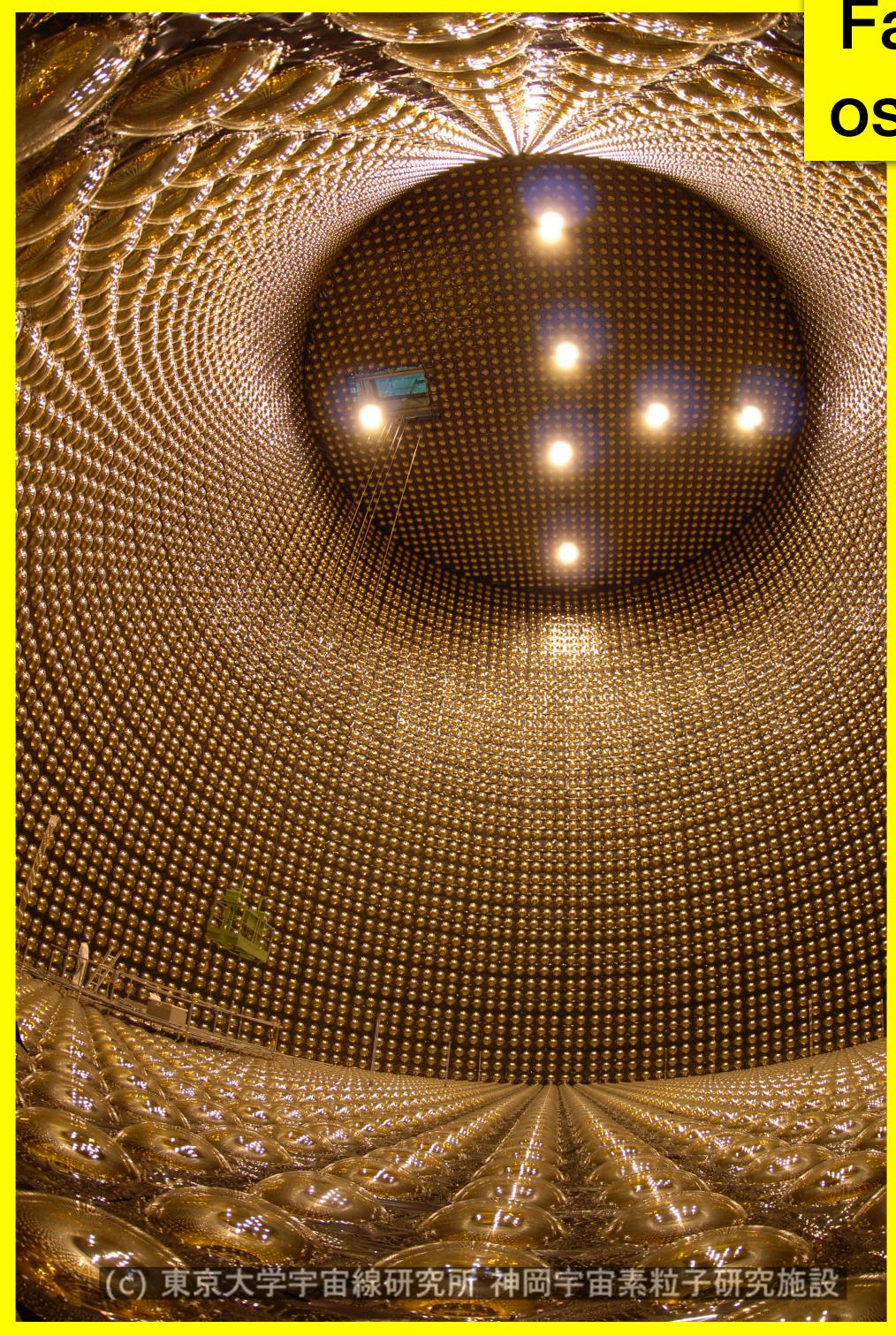


Figure from [here](#)

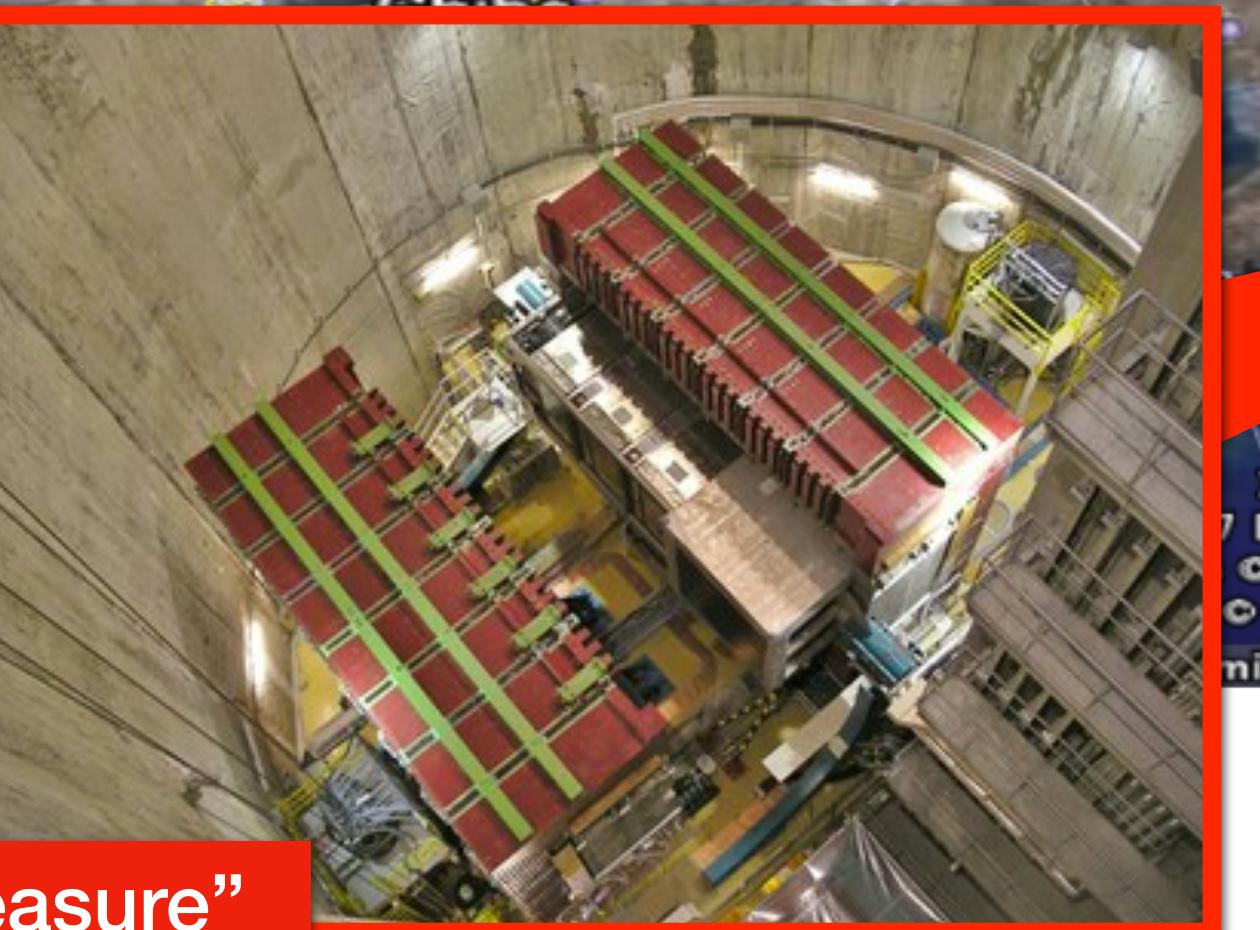
- **PMNS framework validity?**
  - Sterile neutrinos, non standard neutrino interactions etc.

# Tokai-to-Kamioka (T2K) Experiment

Far detector “measures” the oscillated beam composition



Near-detectors “measure”  
the initial beam composition



295 km baseline

E141°

Image NASA  
Europa Technologies  
© 2007 TerraMetrics  
© 2007 ZENRIN  
ning ||||| 100%

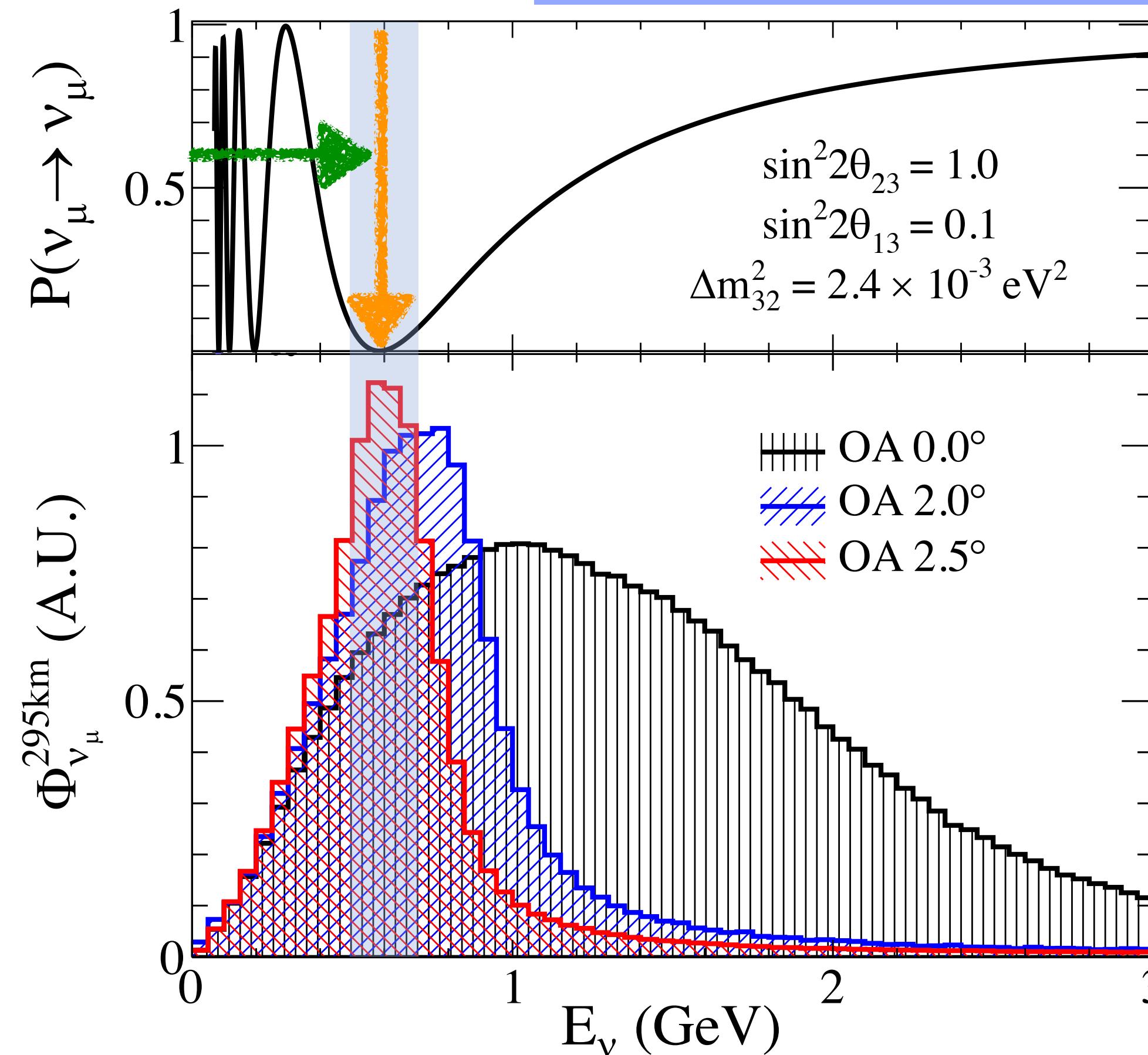


Neutrino beam  
produced at J-PARC

# $\nu_\mu$ ( $\bar{\nu}_\mu$ ) survival probability at T2K

- Muon (anti)neutrino survival probability around oscillation maximum:

$$P_{T2K}(\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu) \approx 1 - 4 \cos^2 \theta_{13} \sin^2 \theta_{23} (1 - \cos^2 \theta_{13} \sin^2 \theta_{23}) \sin^2 \frac{\Delta m_{23}^2 L}{4E}$$



Position of dip in neutrino survival determines  $\Delta m_{32}^2$

Size of dip in neutrino survival determines  $\sin^2 \theta_{23}$

T2K uses a 2.5 degrees off-axis beam for more “monochromatic-like” neutrinos, peaking at the right energy for an oscillation maximum at the position of Super-K

CP-conserving terms (survival probability is CP-conserving even at higher orders!)

# $\nu_e$ ( $\bar{\nu}_e$ ) appearance probability at T2K

- Electron (anti)neutrino appearance probability in vacuum:

$$P_{T2K}(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \frac{\Delta m_{32}^2 L}{4E}$$

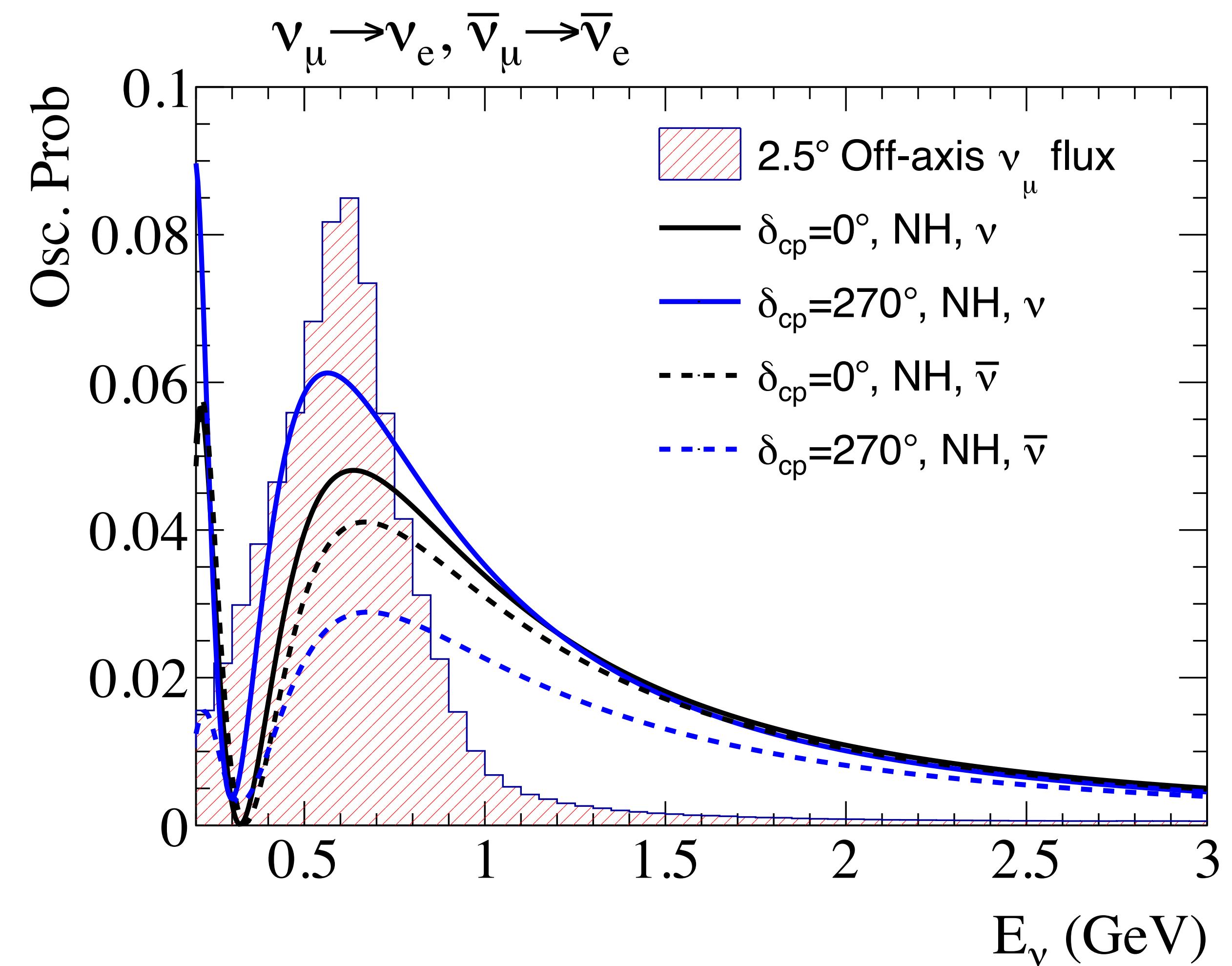
*CP-conserving term*

$$+ J \cos \left( \frac{\Delta m_{32}^2 L}{4E} \right) \oplus \delta_{CP}$$

Minus sign for  $\bar{\nu}$

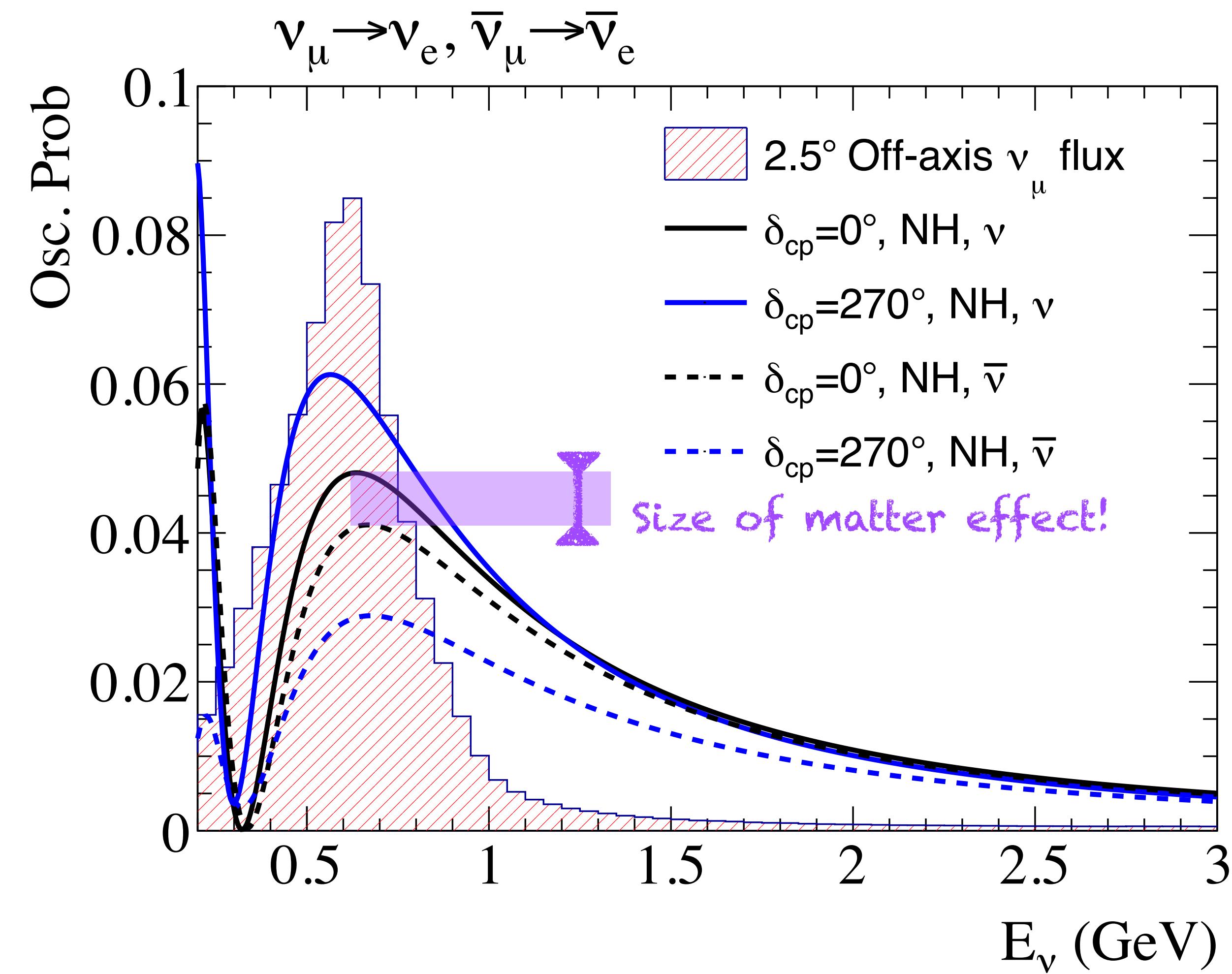
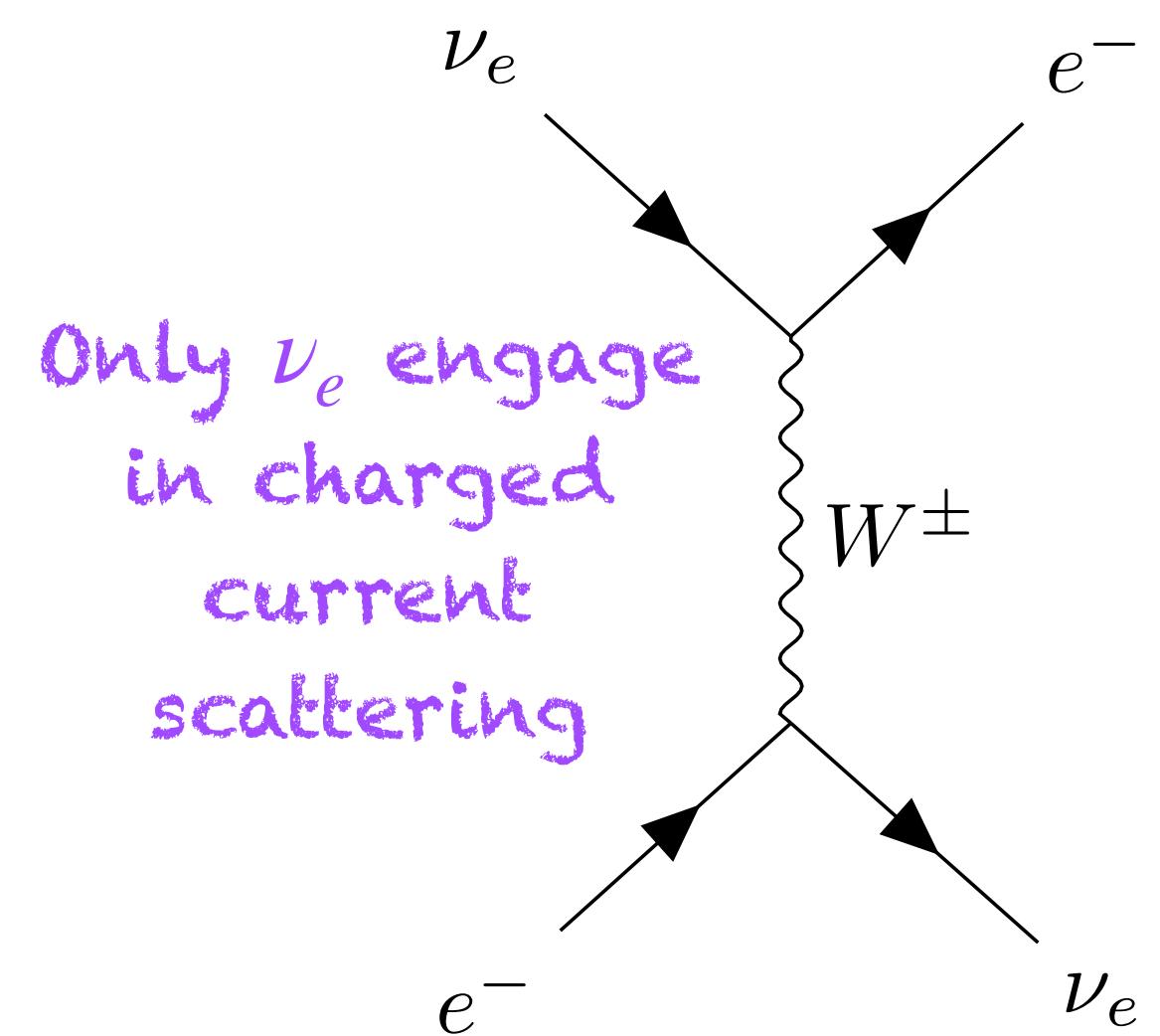
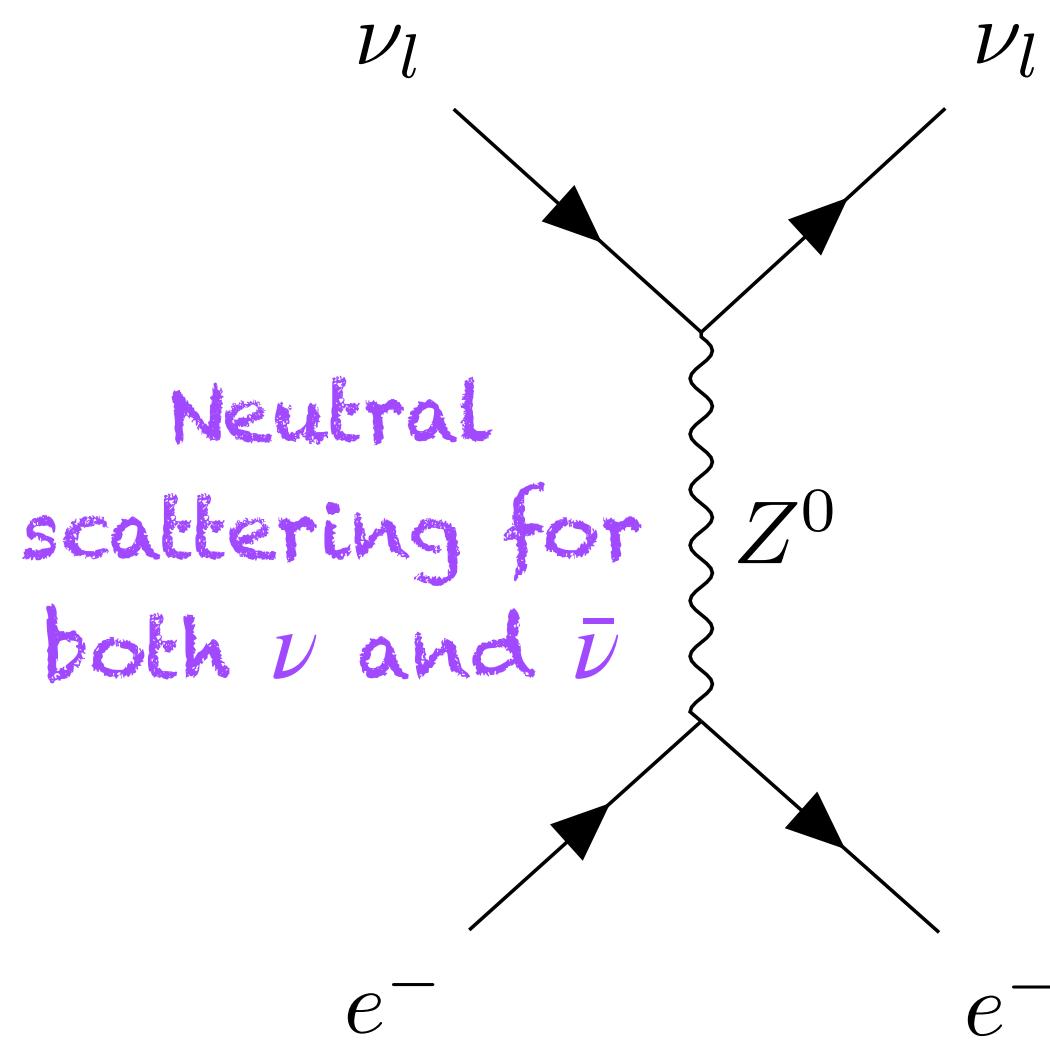
CP-violating term for  $\delta_{CP} \neq 0, \pi$

- The appearance probability gets modified because of charged current  $\nu_e$  coherent scattering from electrons in the Earth's crust

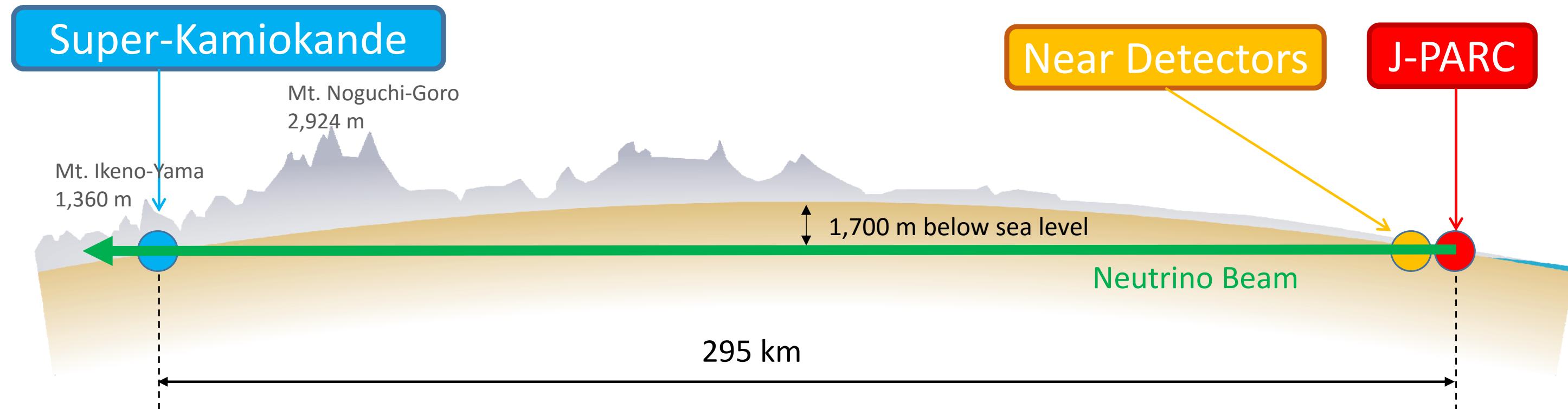
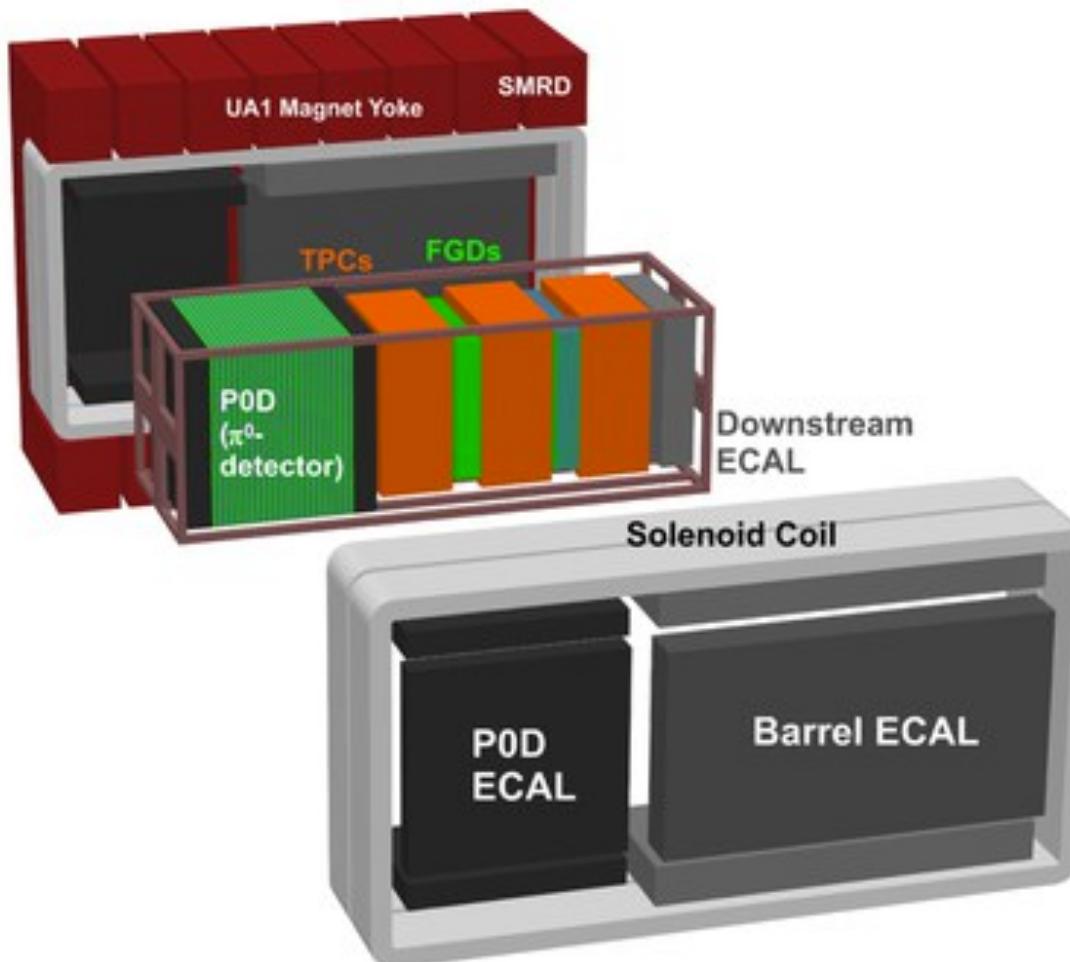


# $\nu_e$ ( $\bar{\nu}_e$ ) appearance probability at T2K

- The appearance probability gets modified because of charged current  $\nu_e$  coherent scattering from electrons in the Earth's crust
- Mikheyev-Smirnov-Wolfenstein effect



# Neutrino-induced event counting at T2K



$\nu_\mu$ -induced event count at ND280:

$$N_\mu^{ND}(E_\nu^{rec}) = \sigma^{ND}(E_\nu^{true}) \times \Phi_\nu^{ND}(E_\nu^{true}) \times \epsilon^{ND}(E_\nu^{true}) \times S^{ND}(E_\nu^{true}, E_\nu^{rec})$$

Cross section on  
 $CH$  and  $H_2O$

Unoscillated  
flux  
(detector specific)

Efficiency  
(detector specific)

Smearing  
matrix  
(detector specific)

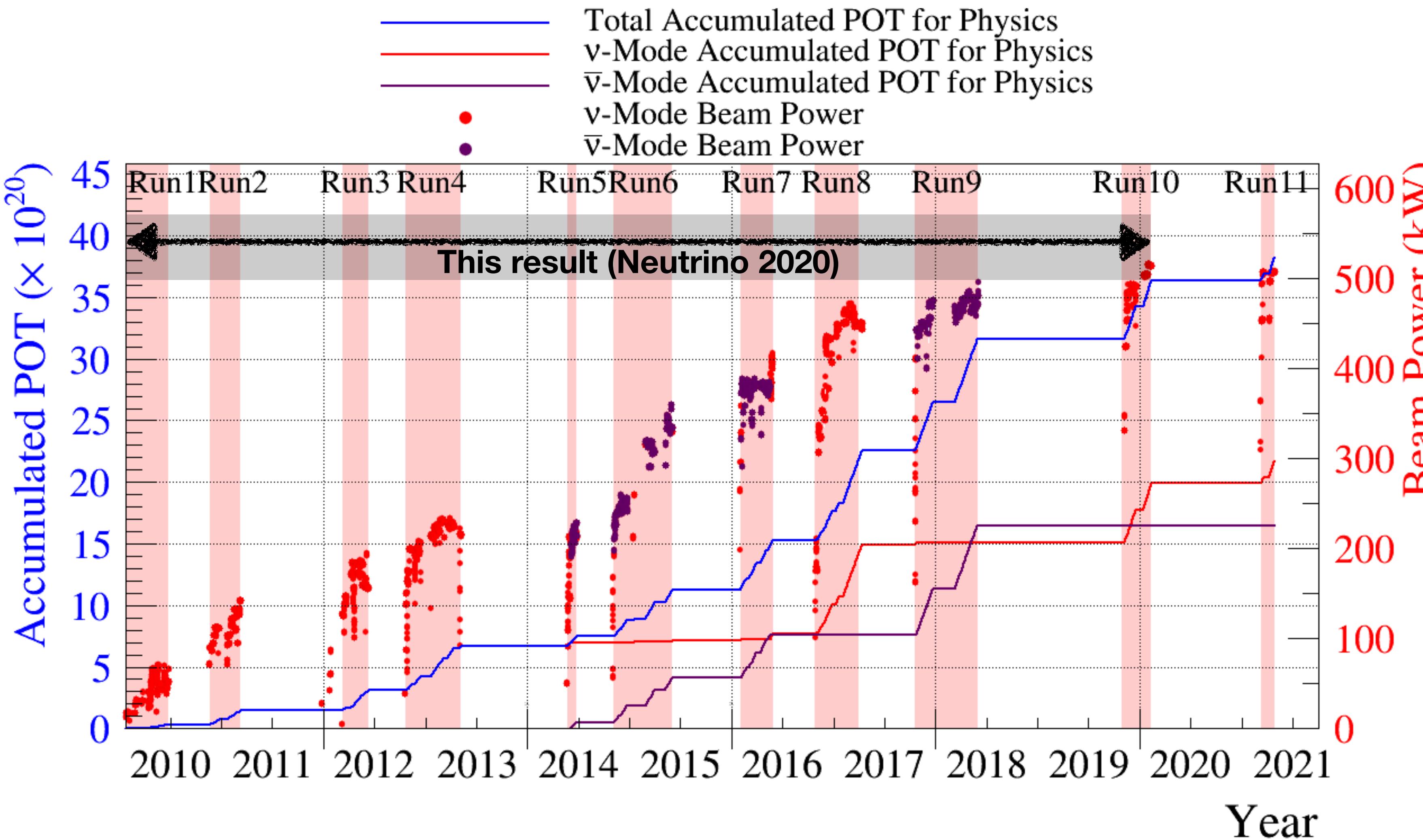
$\nu_\mu$  ( $\nu_e$ )-induced event count at Super-K:

$$\frac{N_\mu^{SK}(E_\nu^{rec})}{N_e^{SK}(E_\nu^{rec})} = \frac{P(\nu_\mu \rightarrow \nu_\mu)}{P(\nu_\mu \rightarrow \nu_e)} \times \sigma^{SK}(E_\nu^{true}) \times \Phi_\nu^{SK}(E_\nu^{true}) \times \epsilon^{SK}(E_\nu^{true}) \times S^{SK}(E_\nu^{true}, E_\nu^{rec})$$

Oscillation  
probability

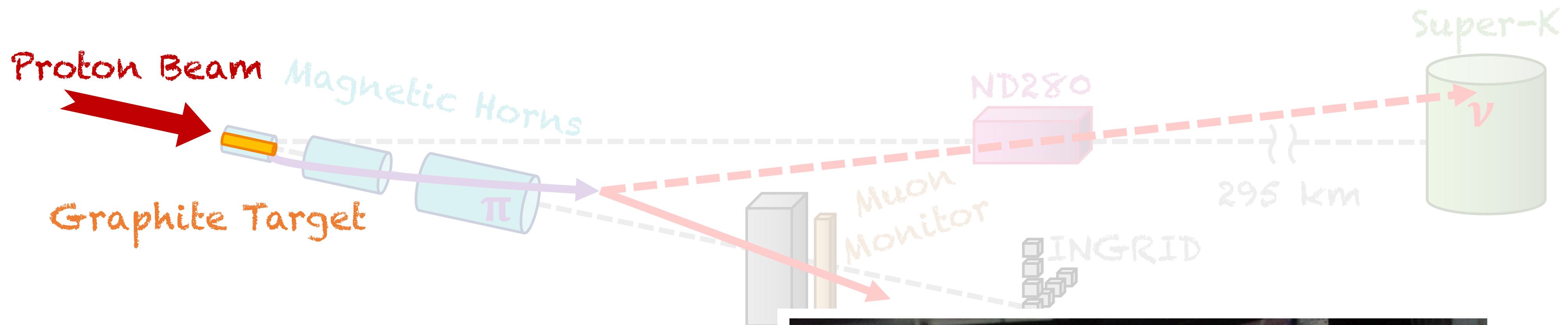
Cross section on  $H_2O$

# T2K datasets



- Reached maximum beam power of 522.6 kW
- For today's result (Runs 1-10), collected  $3.6 \times 10^{21}$  POT, split into  $1.96 \times 10^{21}$  POT (54.61%) in  $\nu$ -mode and  $1.63 \times 10^{21}$  POT (45.39%) in  $\bar{\nu}$ -mode
- **Today's result includes ~33% more POT in  $\nu$ -mode compared to our previous measurement**
- Stay tuned for new results with the inclusion of Run 11, with  $1.78 \times 10^{21}$  POT in  $\nu$ -mode, the very first T2K data with a Gd-doped Super-K

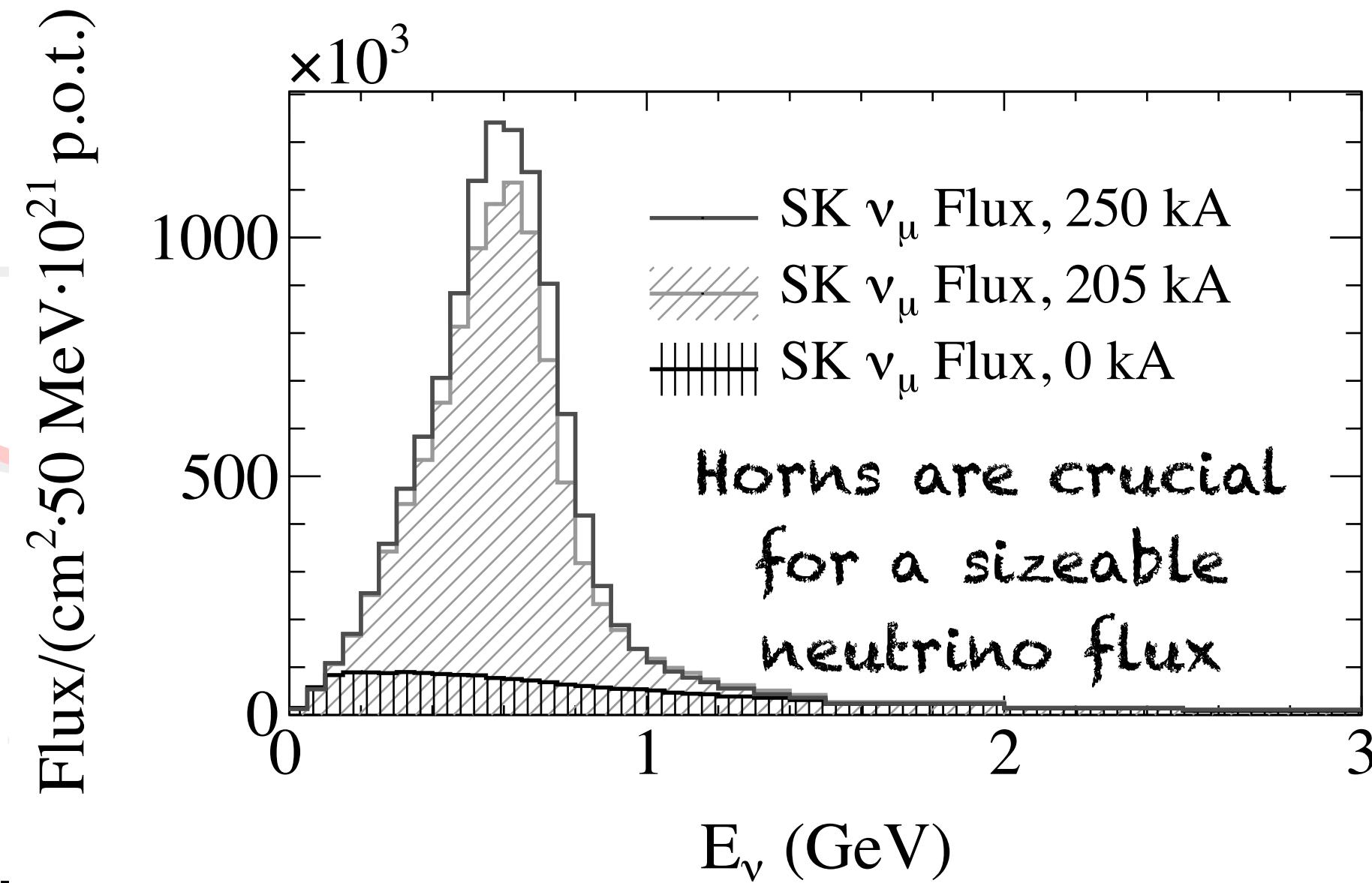
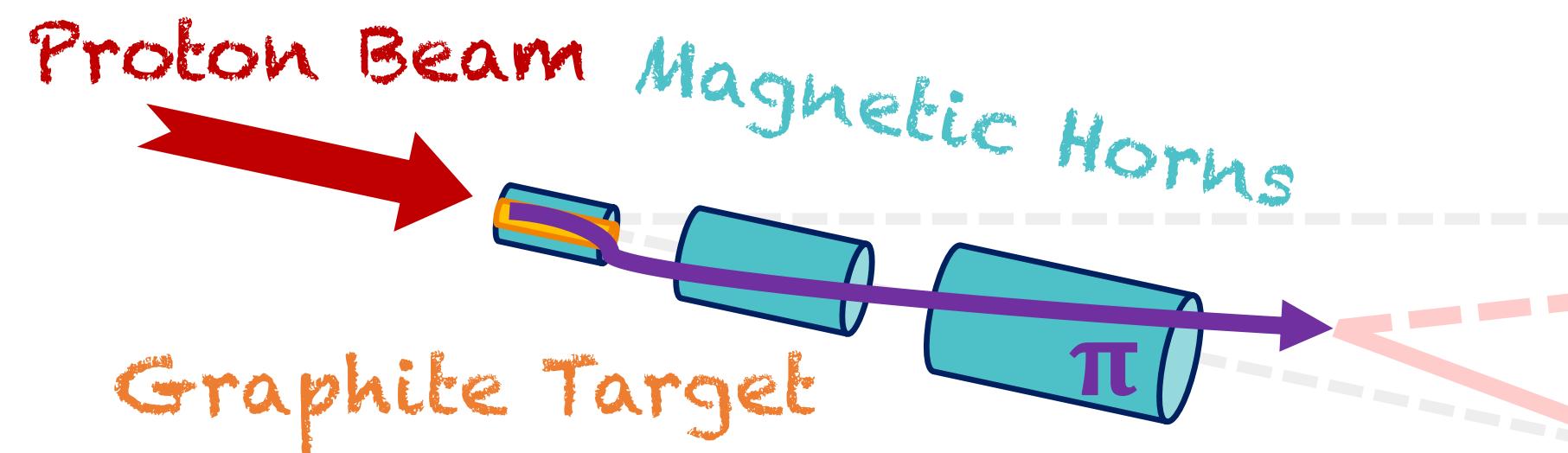
# T2K beamline



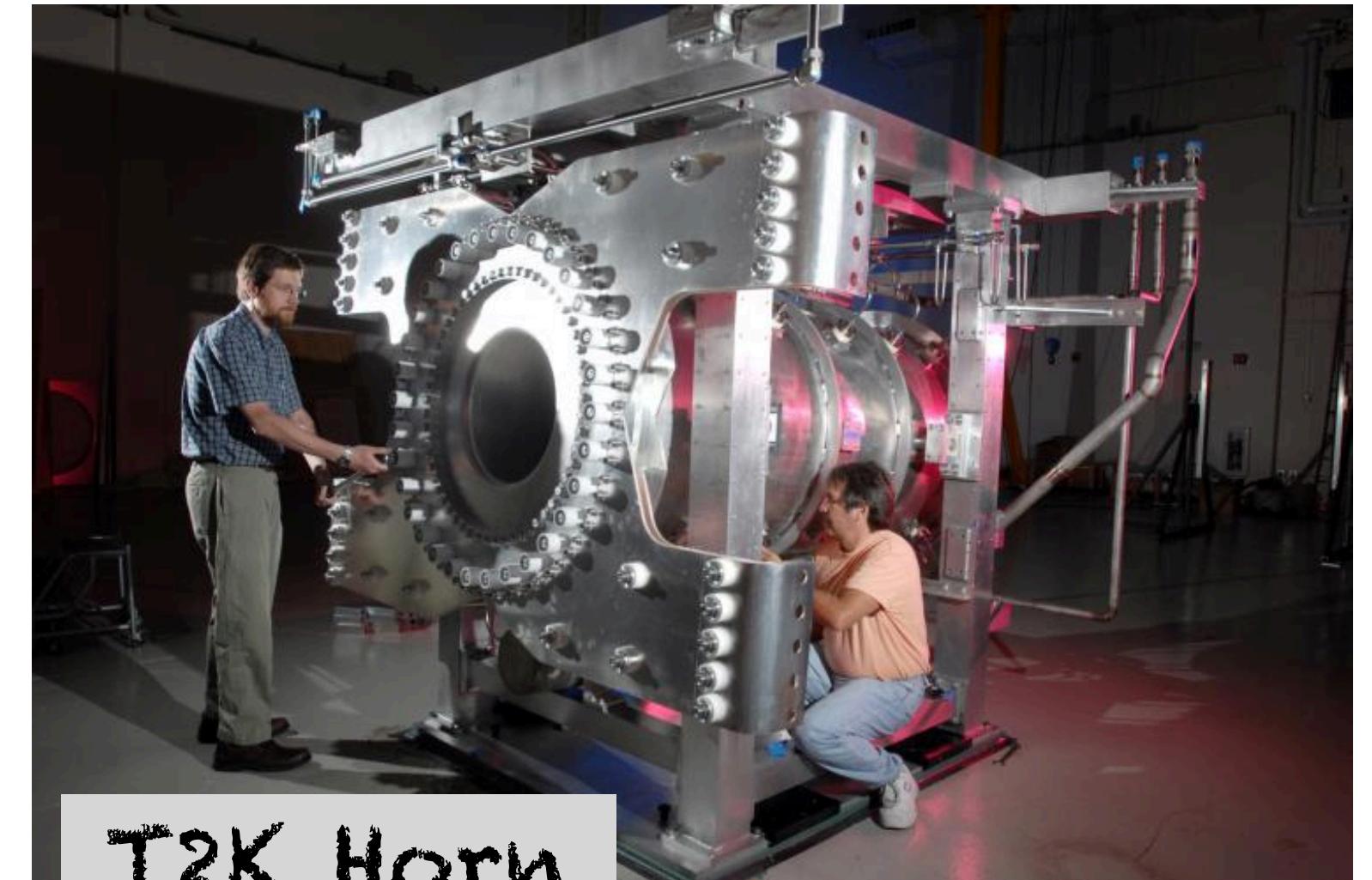
- 30 GeV protons strike the 90 cm long T2K graphite target
- QCD theory is non-perturbative at these energies, thus not easily calculable



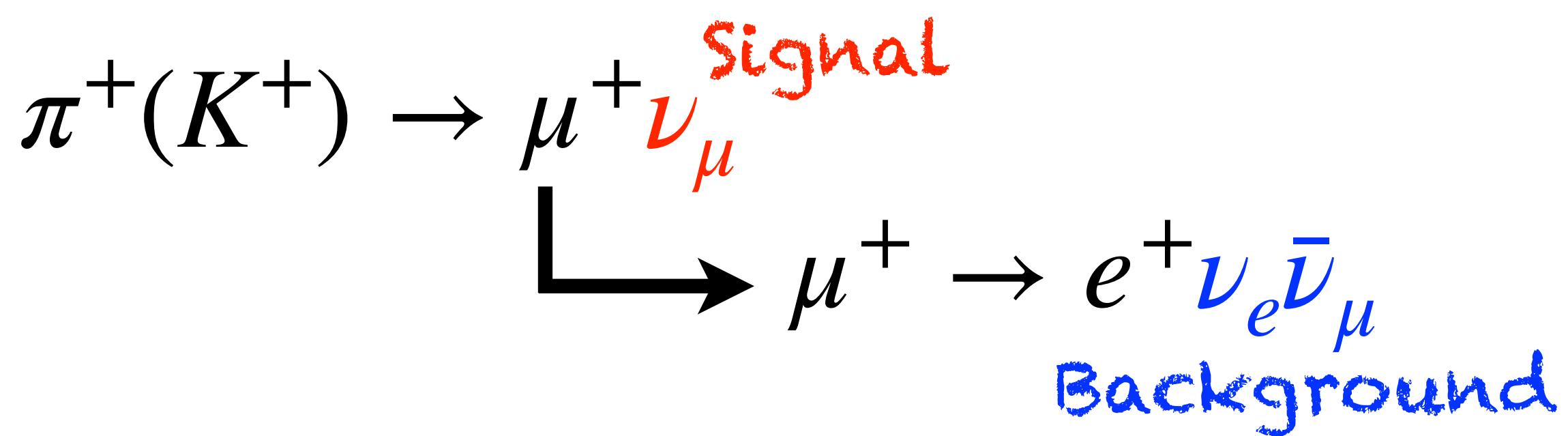
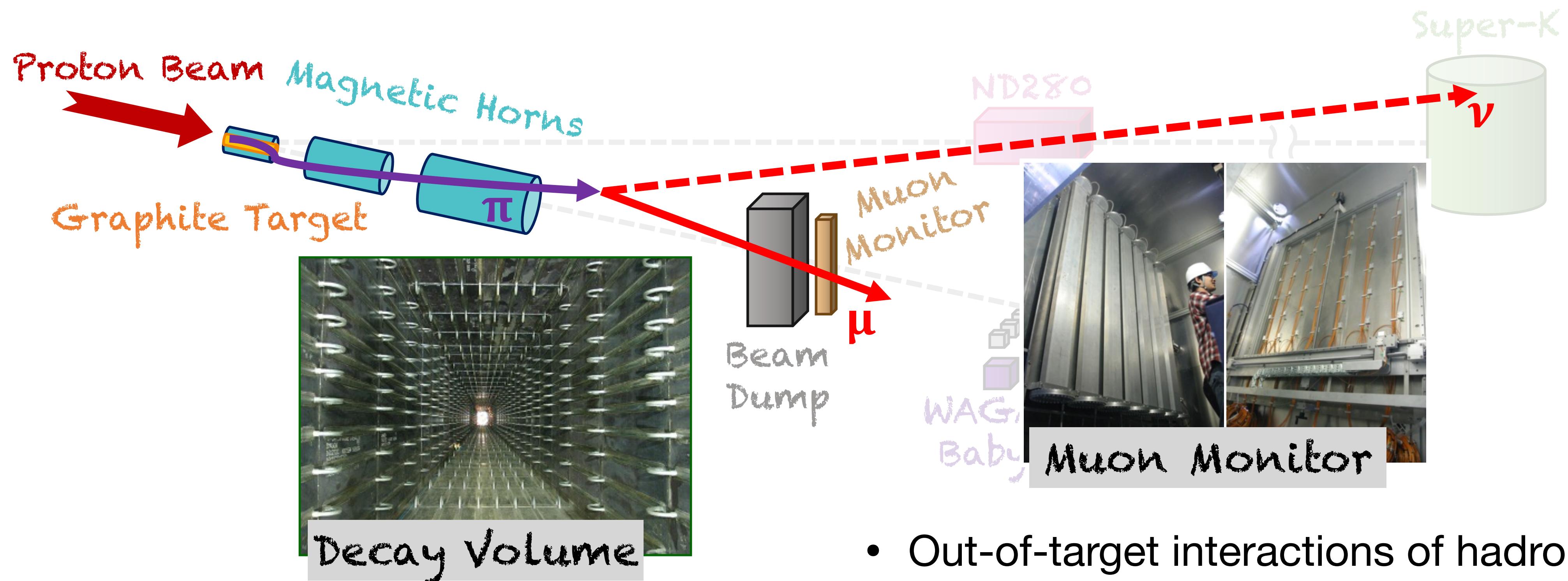
# T2K beamline



- Hadronic cascade is initiated within the target
- 3 magnetic horns used to bend and (de)focus the pions and kaons leaving the target

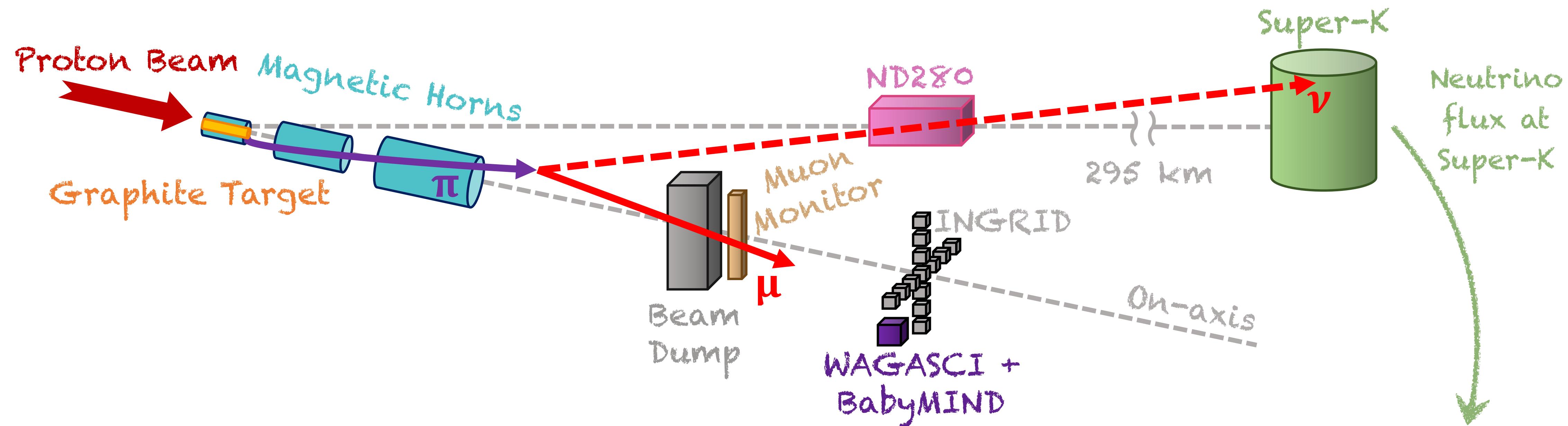


# T2K beamline

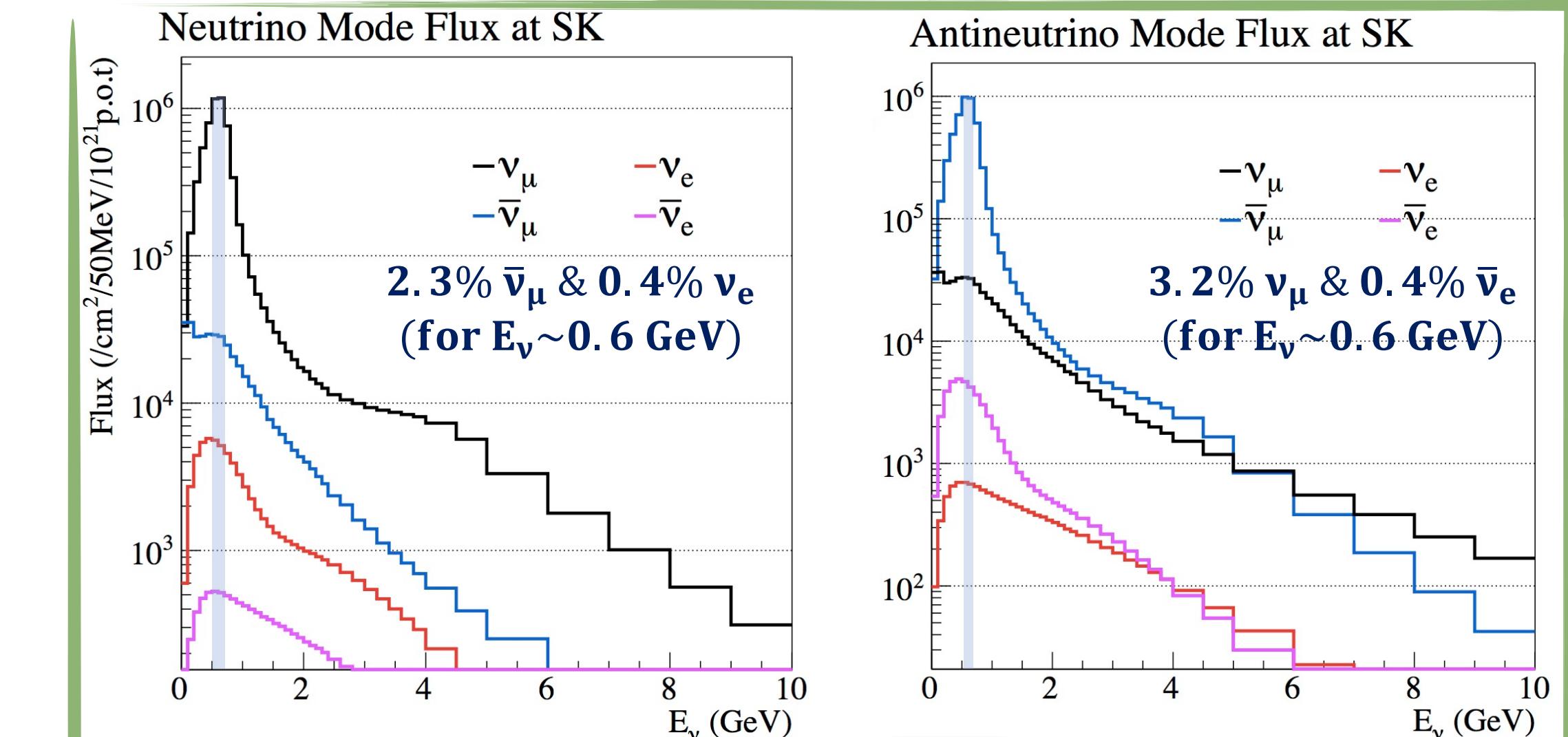


- Out-of-target interactions of hadrons with secondary beamline components affect the neutrino flux
- In-flight pion and kaon decays inside the decay volume (96 m) produce neutrinos

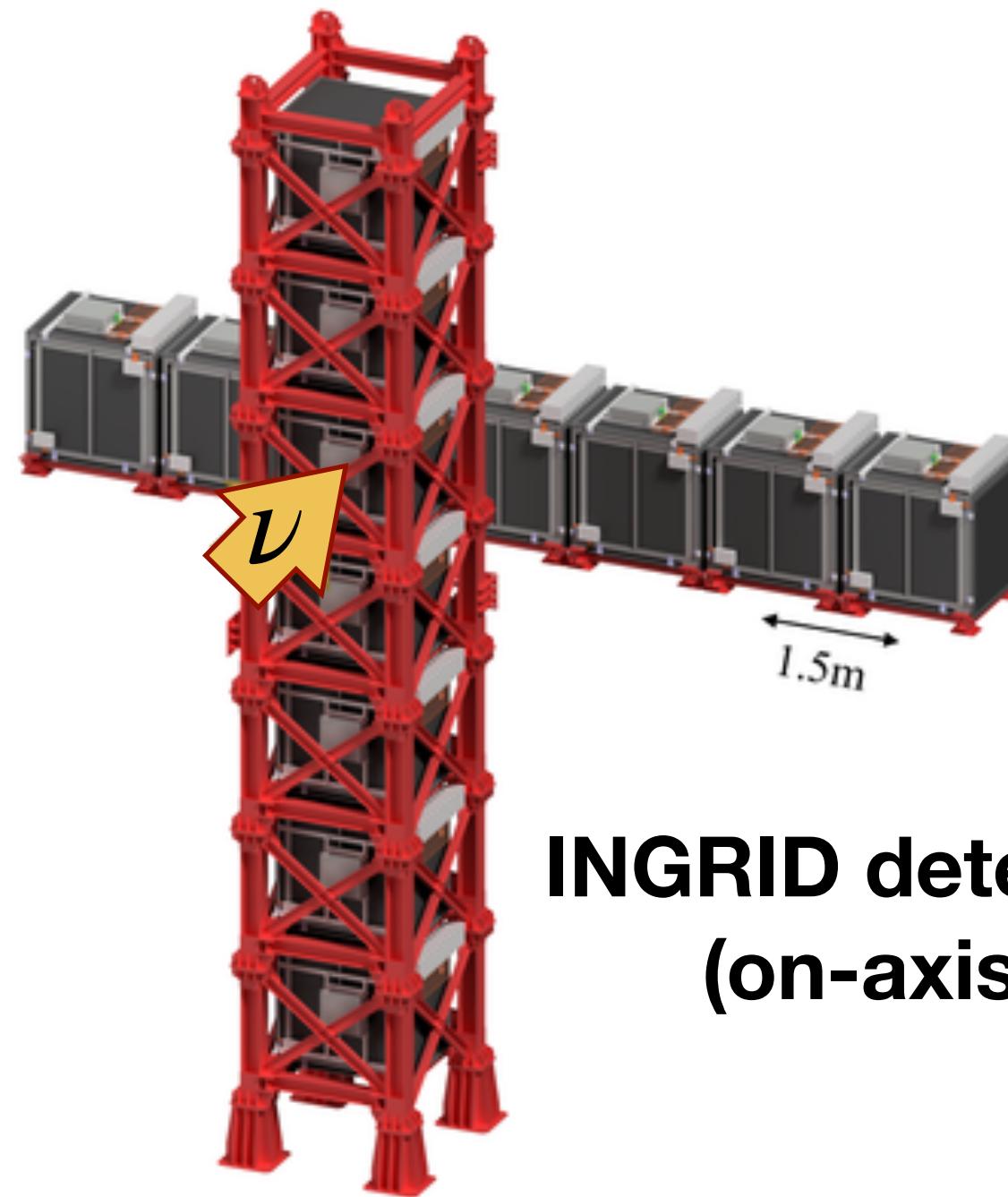
# T2K beamline



- By changing the horn current polarity, we can focus either a neutrino or an anti-neutrino enhanced beam towards the T2K detectors

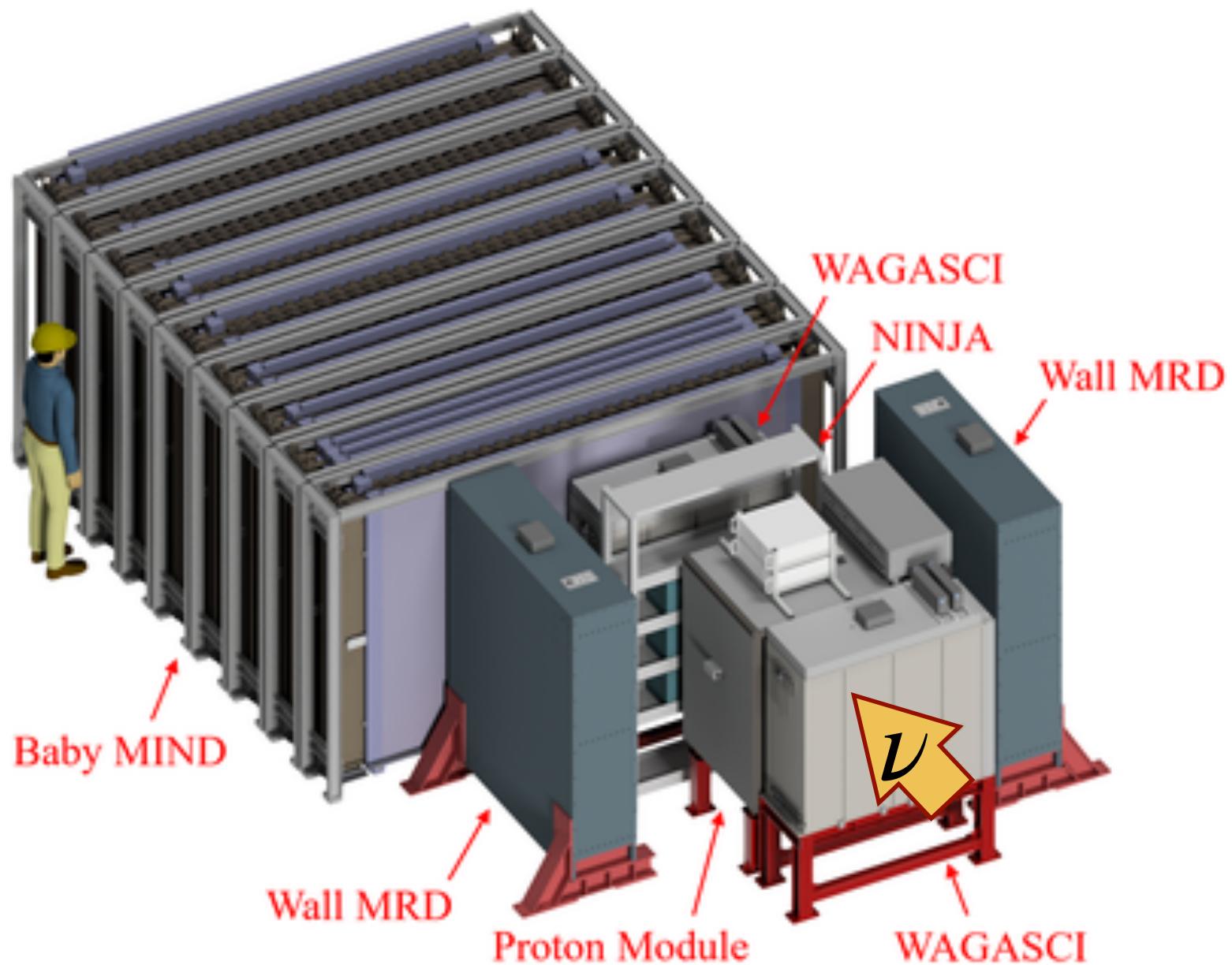


# T2K near detectors



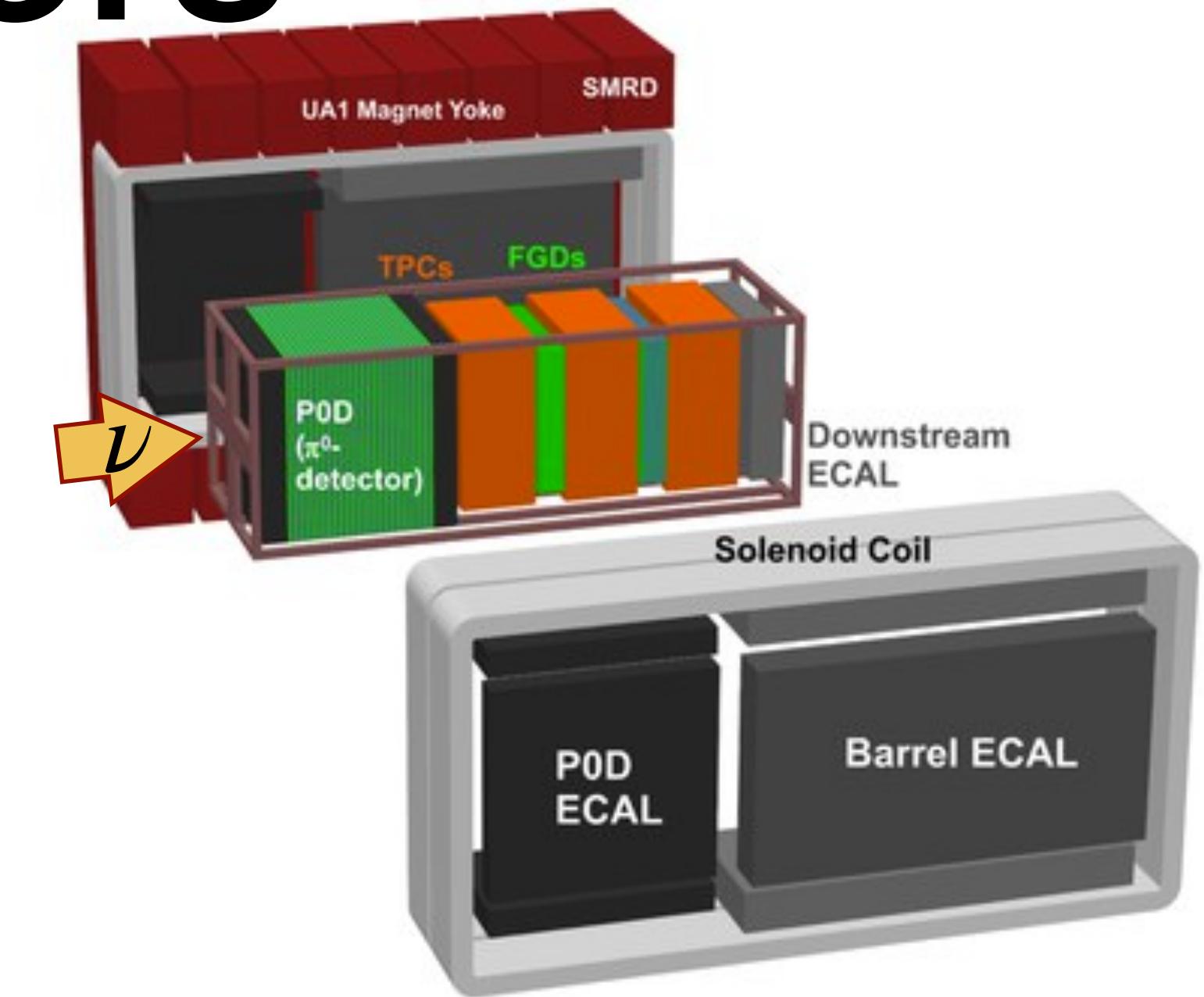
**INGRID detector  
(on-axis)**

- Consists of 16 identical modules, each made of alternating iron and plastic scintillator layers
- In the oscillation analysis, **used to constrain neutrino beam systematics** (e.g. off-axis angle)
- Also measures cross sections for neutrino interaction models



**WAGASCI and BabyMIND detectors  
(1.5 degrees off-axis) *New!***

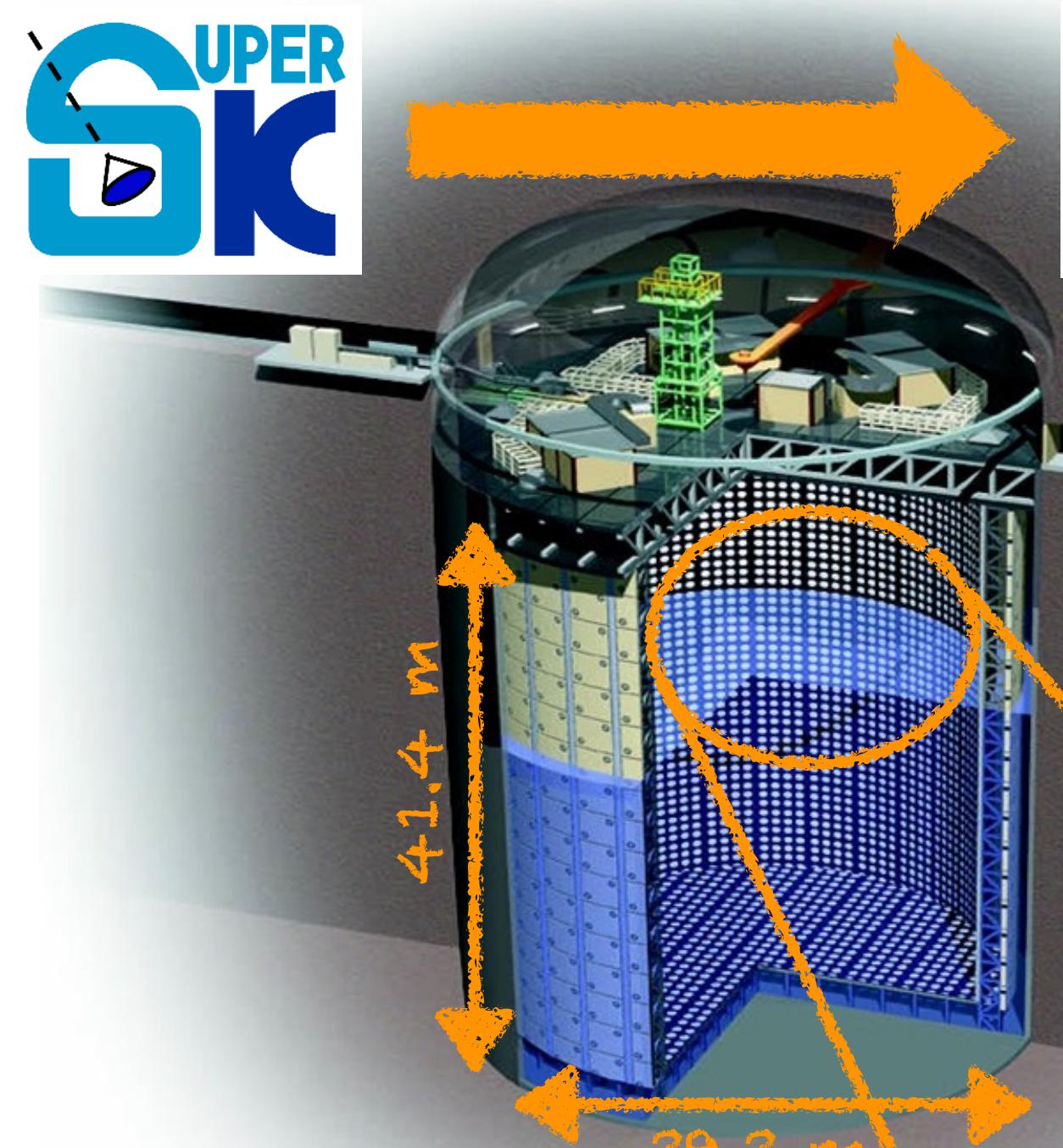
- Water filled plastic scintillator lattice (WAGASCI) and magnetised tracking (BabyMIND) detectors
- Plans to use in the oscillation analysis
- First WAGASCI cross section paper:  
[Phys. Rev. D 97, 012001](#)



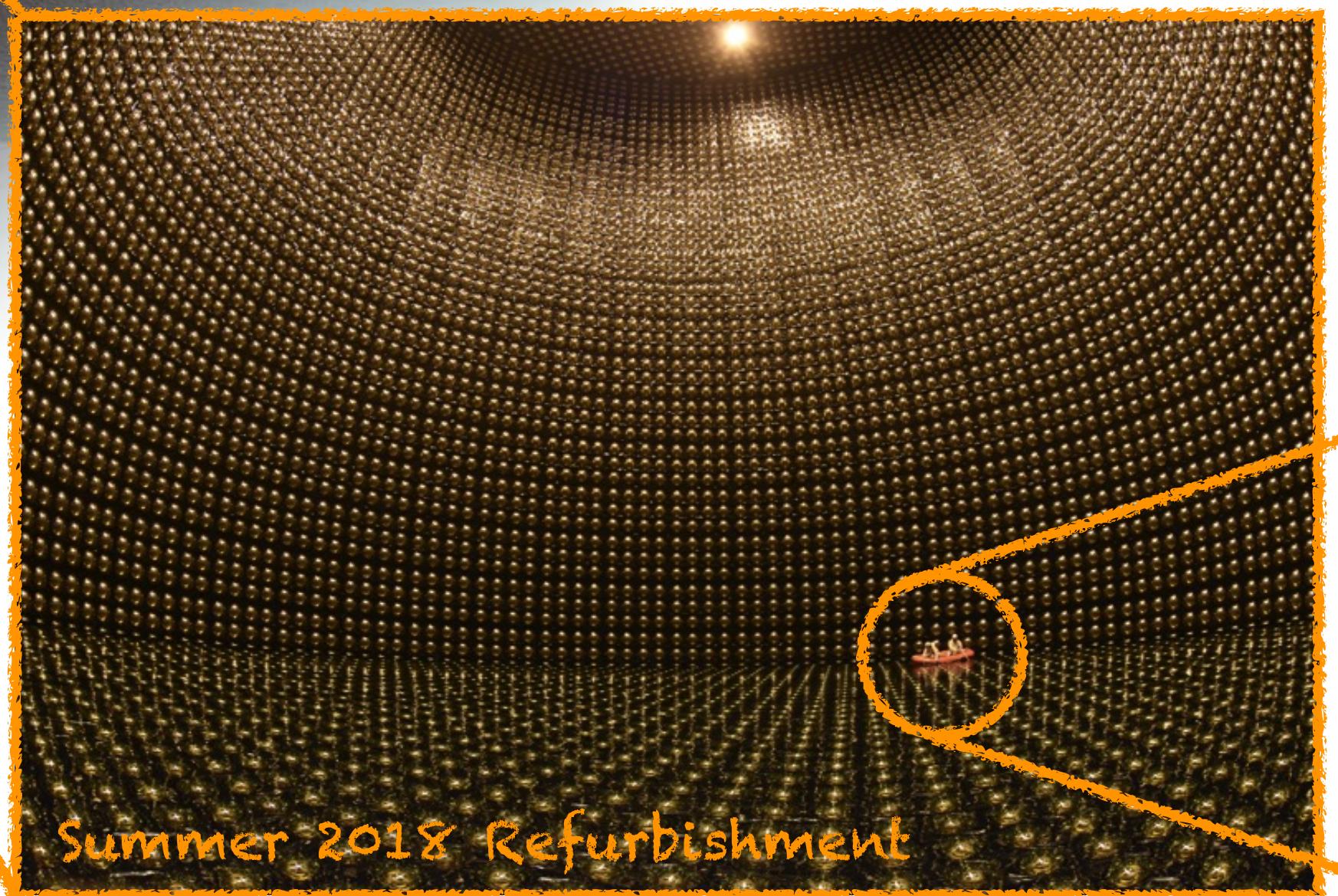
**ND280 detector  
(2.5 degrees off-axis)**

- **Main constraint of flux and cross section model parameters in the oscillation analysis**
- Takes sophisticated cross section measurements for neutrino interaction models (e.g. recently)  
[Phys. Rev. D 103, 112009](#)

# The Super-Kamiokande detector



For more details on Gd loading,  
see the talk from Neutrino 2020  
(click [here](#))



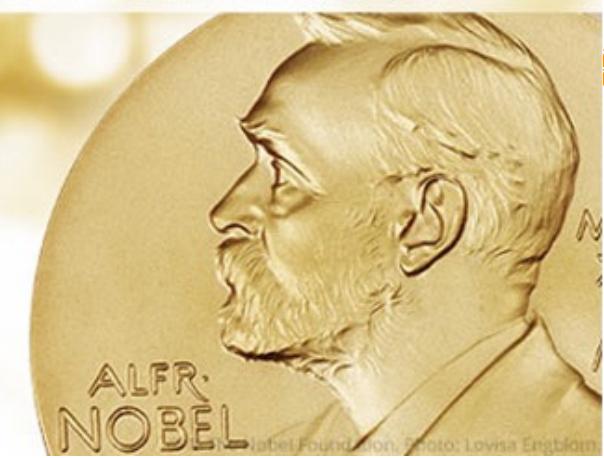
- 50 kt water Cherenkov detector, sensitive to beam, atmospheric, solar, supernova neutrinos
- 11k 20" inner and 1.8k 8" outer (veto) tank PMTs
- Now Gd loaded for enhanced neutron detection, with the potential to separate between  $\nu$  and  $\bar{\nu}$  interactions



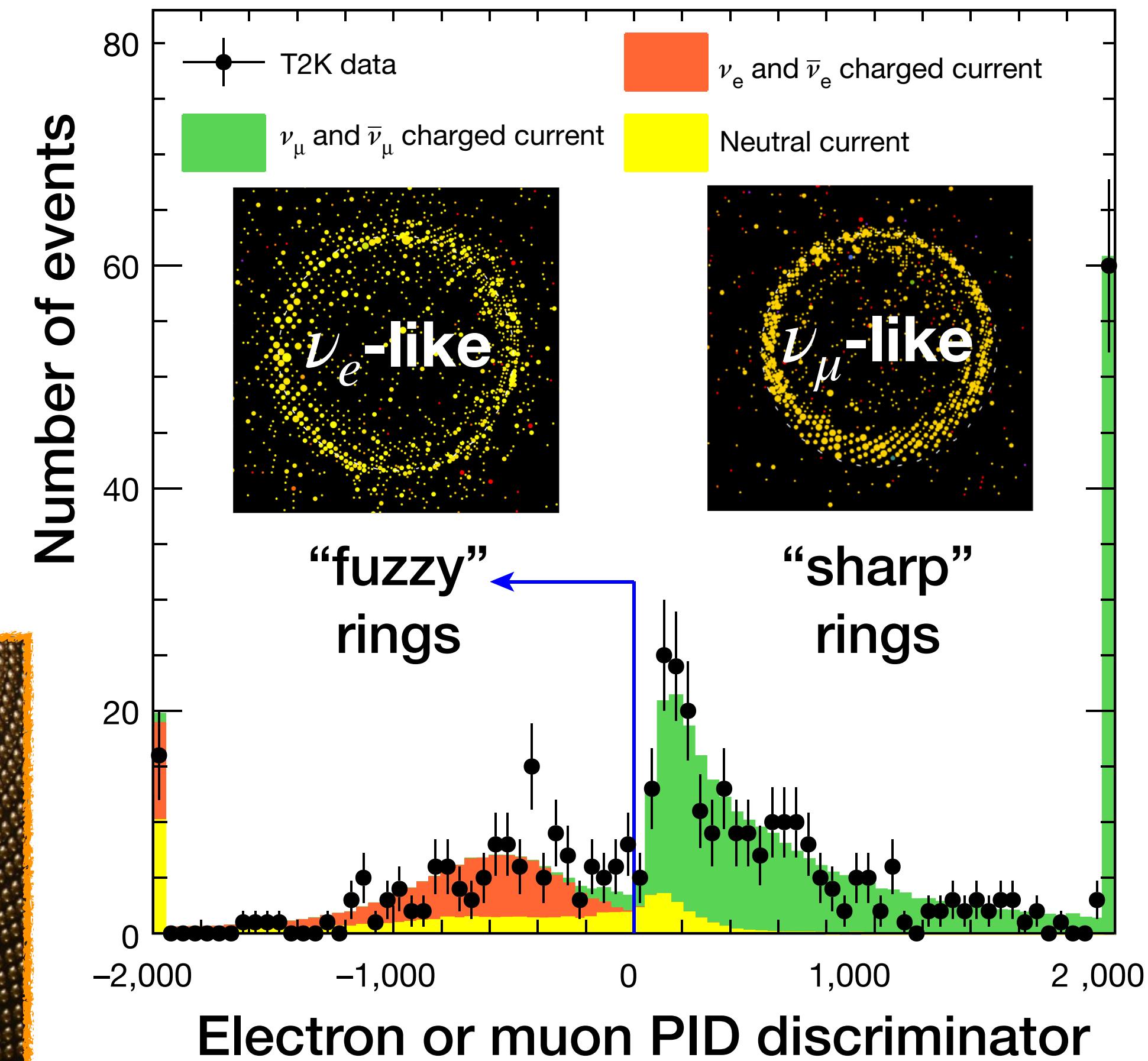
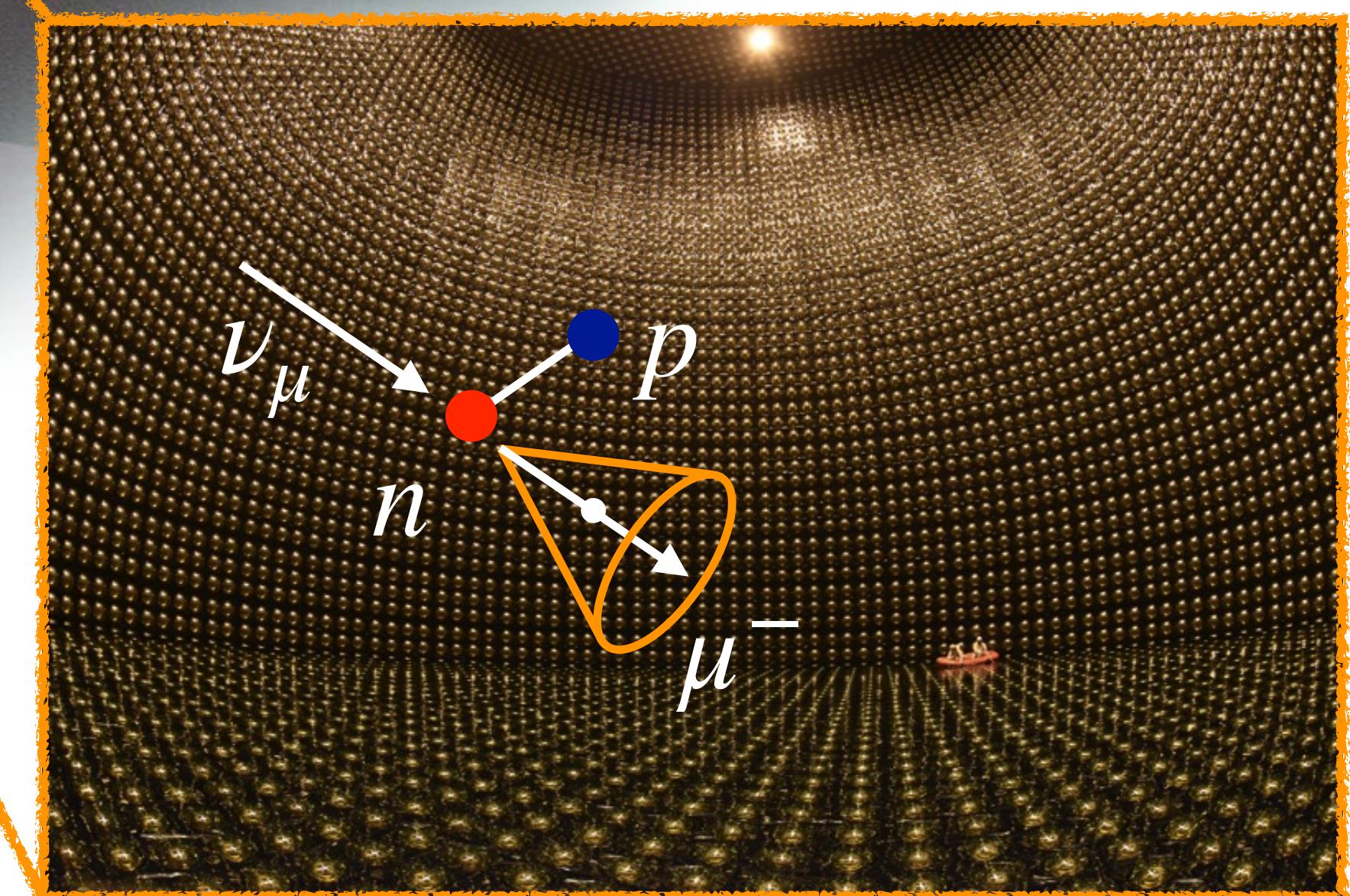
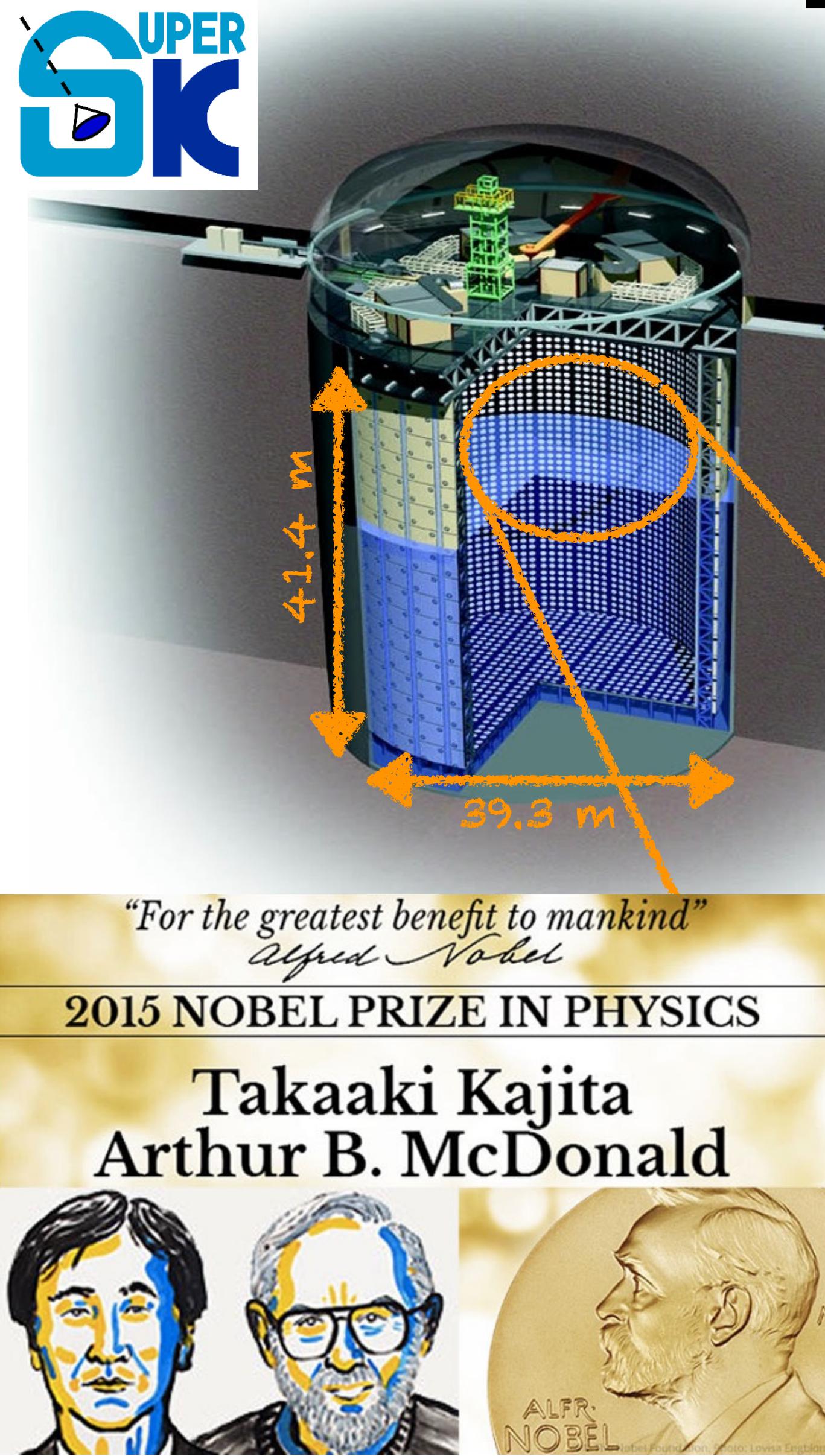
"For the greatest benefit to mankind"  
Alfred Nobel

2015 NOBEL PRIZE IN PHYSICS

Takaaki Kajita  
Arthur B. McDonald



# The Super-Kamiokande detector



PID based on "fuzziness" of Cherenkov rings

# Analysis strategy: flux model

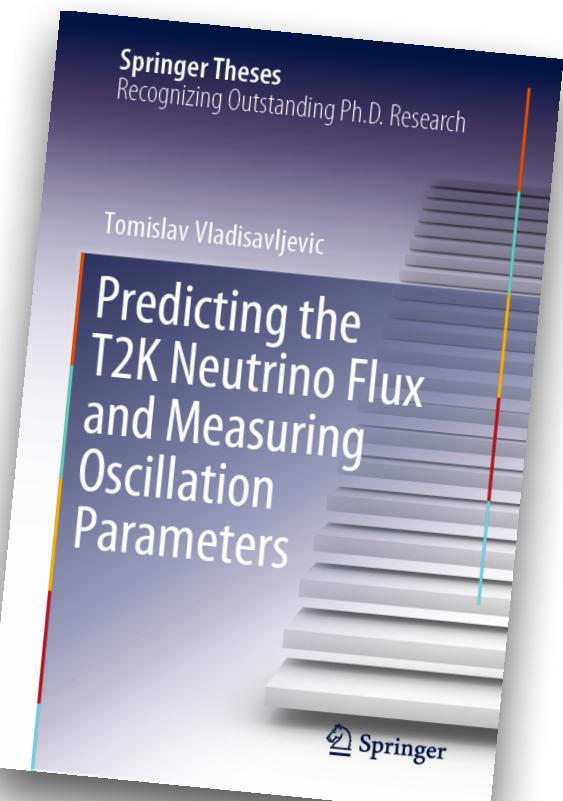
NA61/SHINE data



Neutrino flux model

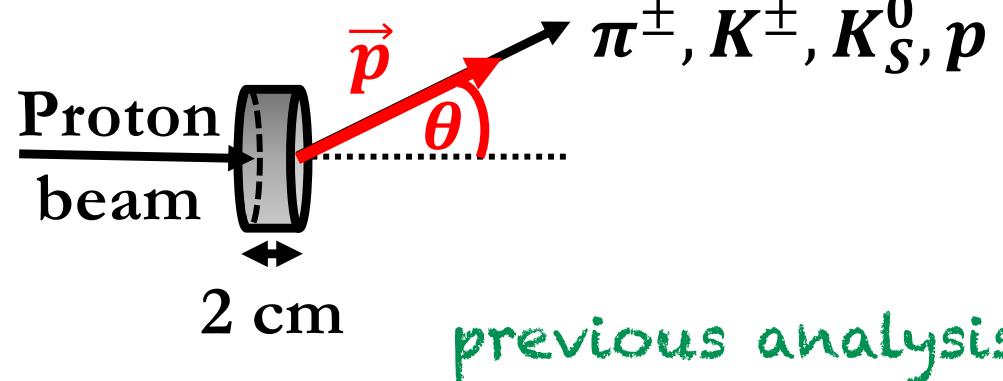
- The addition of NA61 replica-target data reduces the T2K flux prediction uncertainty from  $\sim 10\%$  to  $\sim 6\%$  in the latest analysis

(click [here](#) for more details)

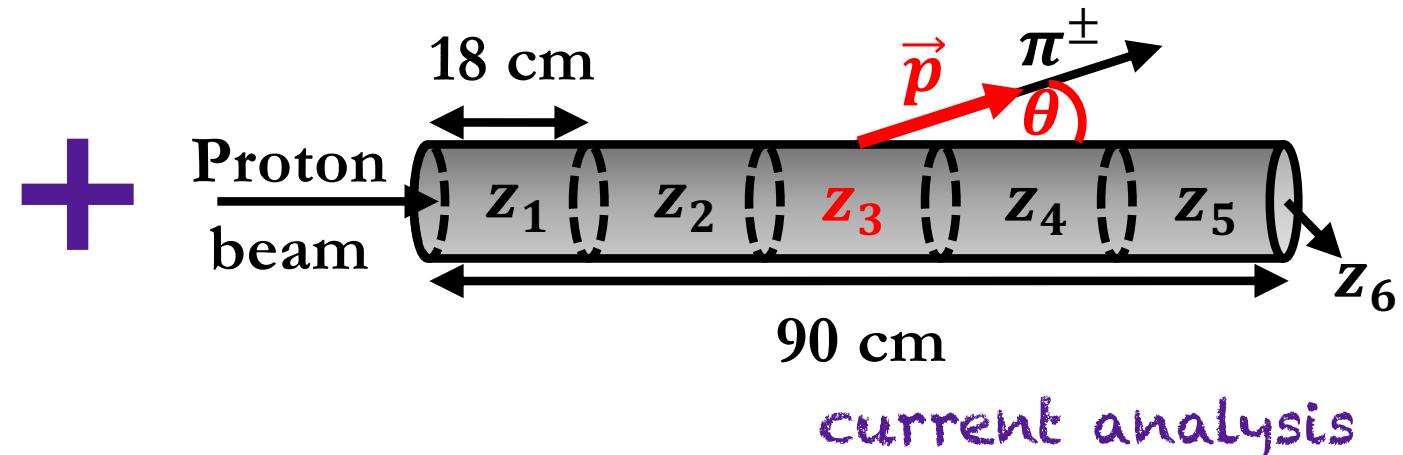


External NA61 hadron production data usage:

Thin-target data:



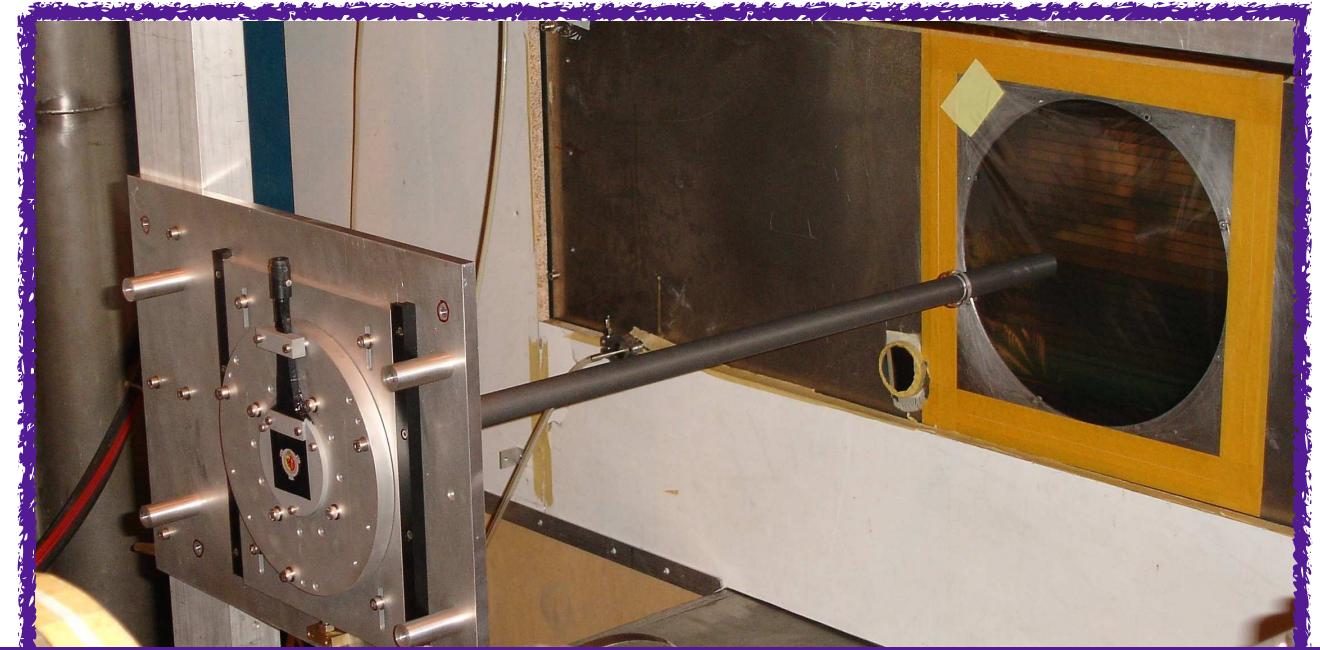
Replica-target data:



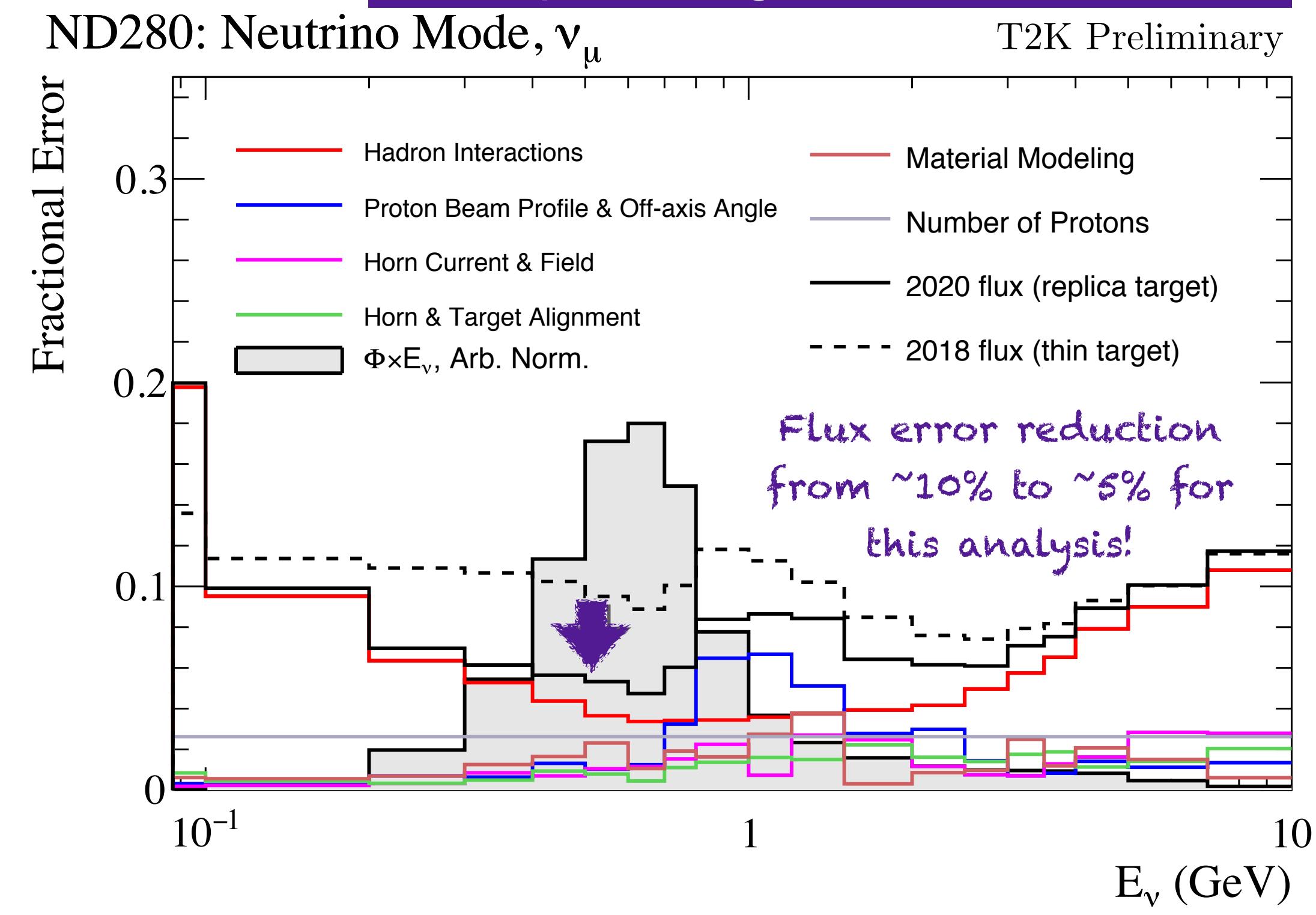
NA61 2009 thin-target data:  
Eur. Phys. J. C 76, 84 (2016)

NA61 2009 replica-target data:  
Eur. Phys. J. C 76, 617 (2016)

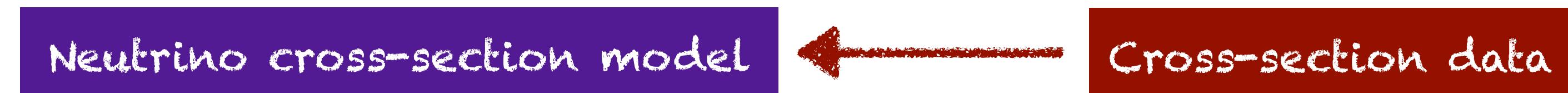
T2K beam monitor  
and INGRID data



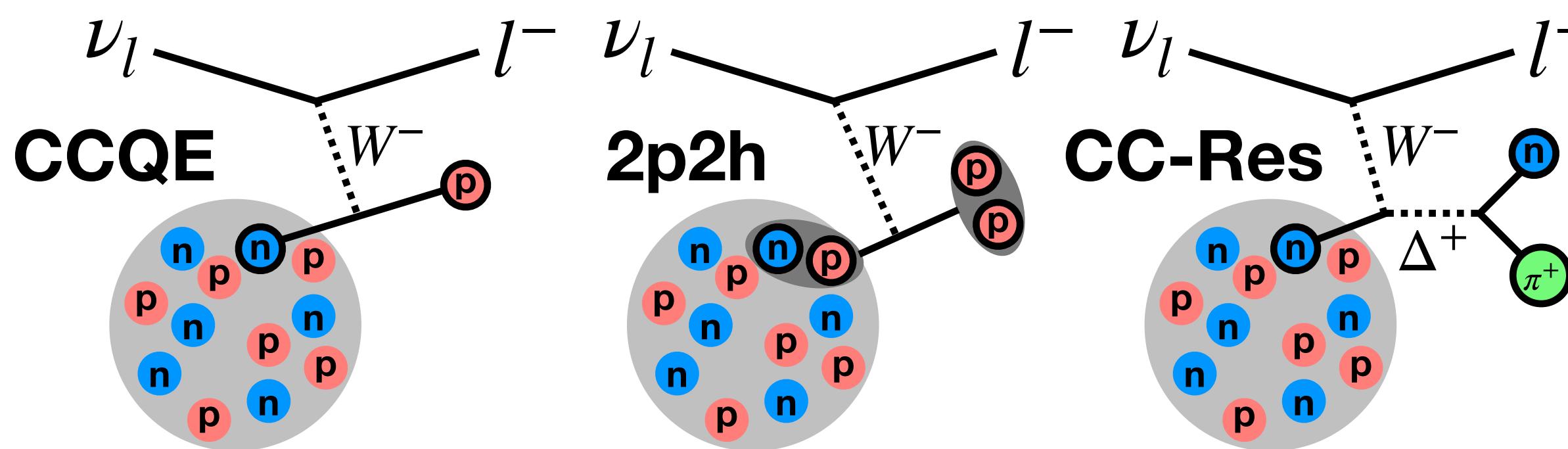
T2K replica-target mounted at NA61



# Analysis strategy: cross section model

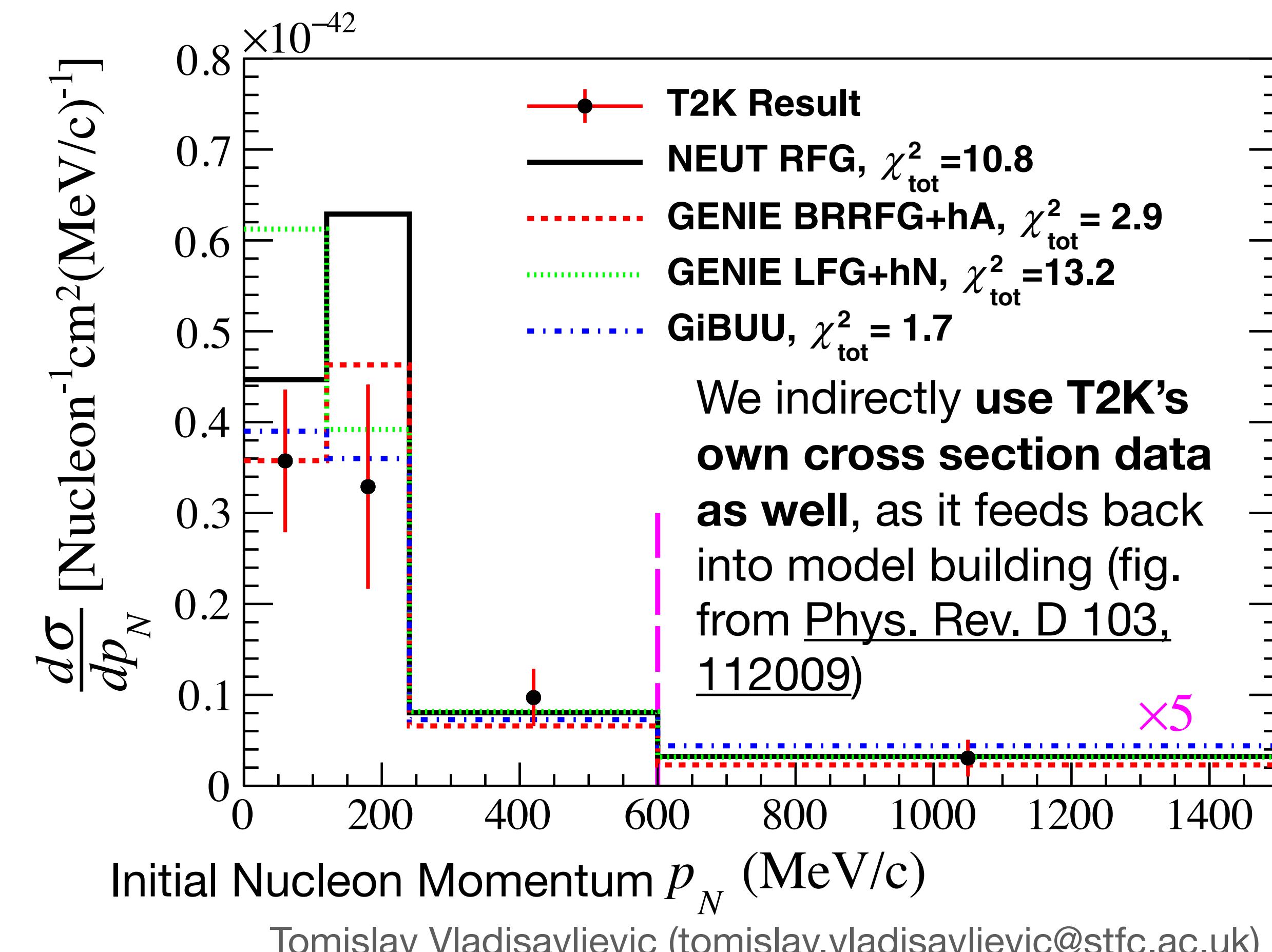


- At  $\sim 0.6$  GeV peak beam energy, neutrinos interact via **charged current quasi-elastic (CCQE)** scattering, with subdominant contributions from **multi-nucleon “2p2h”** and **resonant CC1 $\pi$  scattering**



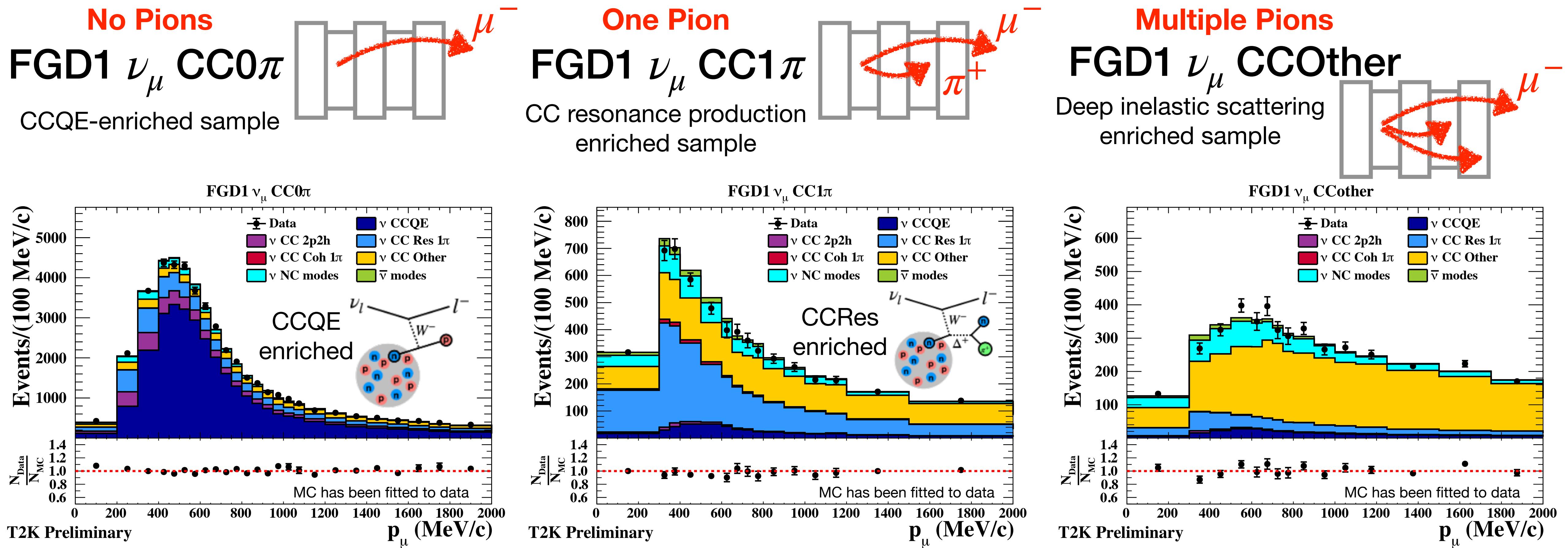
- Main update for our latest oscillation result: using the Spectral Function CCQE initial state nuclear model (strongly favoured by electron scattering data)

- Some of the most useful data for our model is still bubble chamber data from the 1980s!



# ND280 Analysis Samples

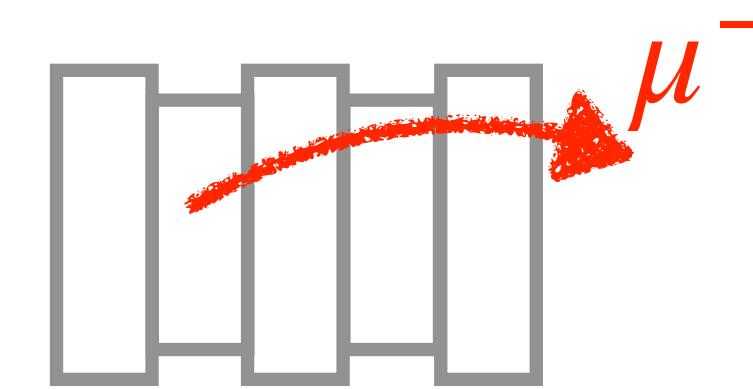
- $18 = 3 \times (1 + 2) \times 2$  classes of neutrino interactions defined at ND280, split by:
  - **3 Interaction Modes, based on pion content:**  $CC0\pi$ ,  $CC1\pi$ ,  $CCOther$  ( $CCN\pi$ ,  $N > 1$ )



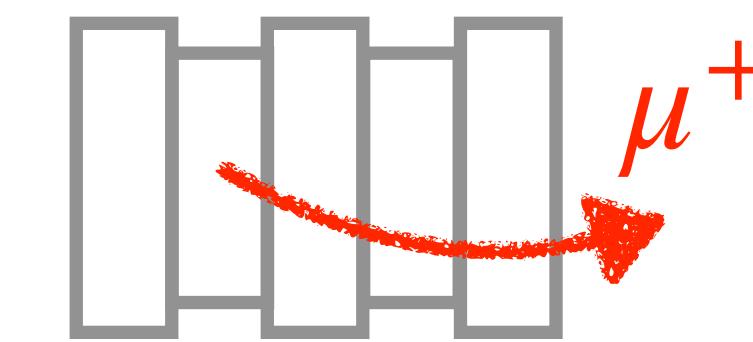
# ND280 Analysis Samples

- $18 = 3 \times (1 + 2) \times 2$  classes of neutrino interactions defined at ND280, split by:
  - **1+2 Modes of Operation:** Neutrino ( $\nu_\mu$  signal) and Anti-neutrino ( $\bar{\nu}_\mu$  signal and  $\nu_\mu$  background)
  - **Separation of anti-neutrino mode samples by pion content is a new feature of this analysis!**

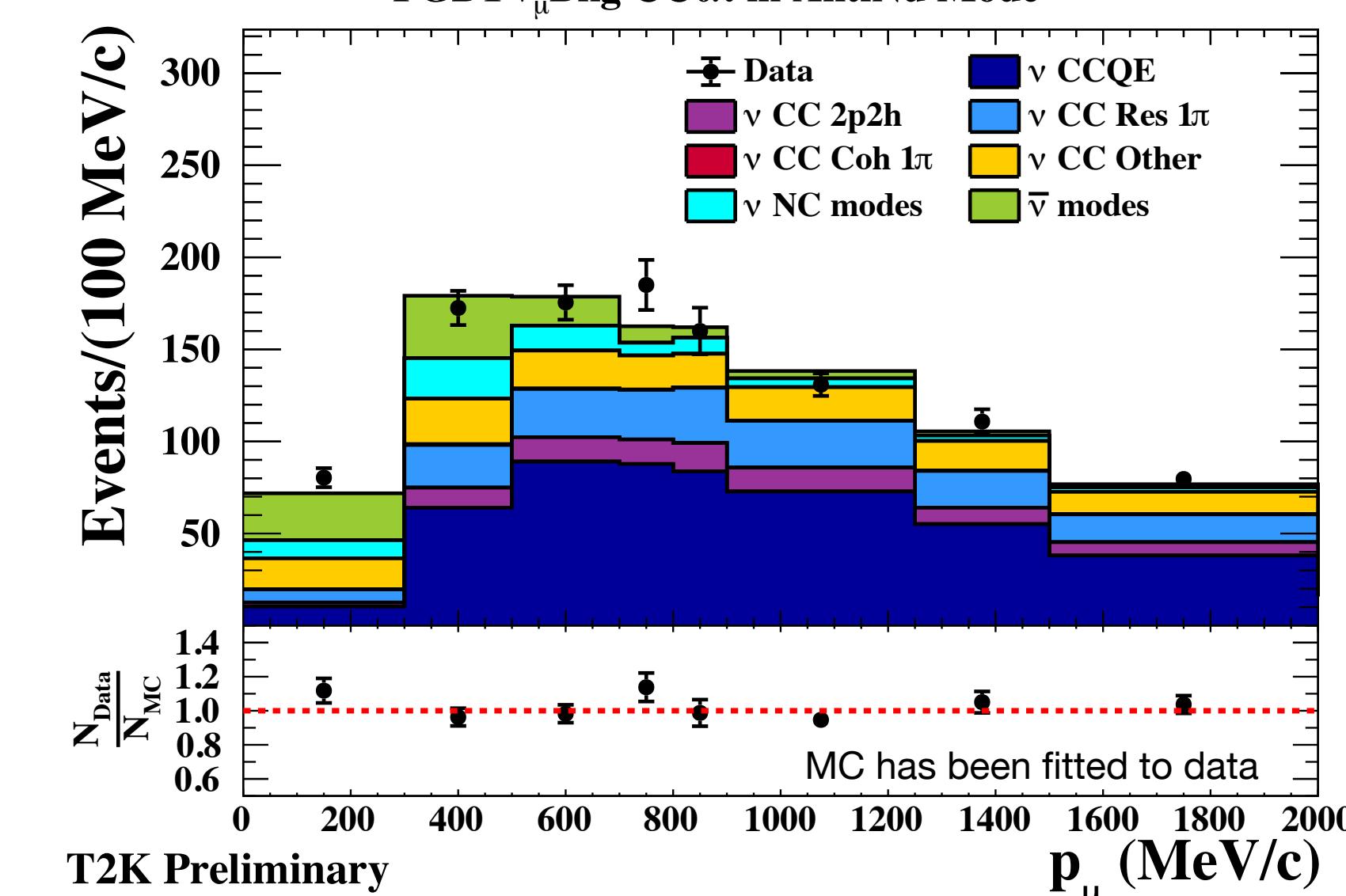
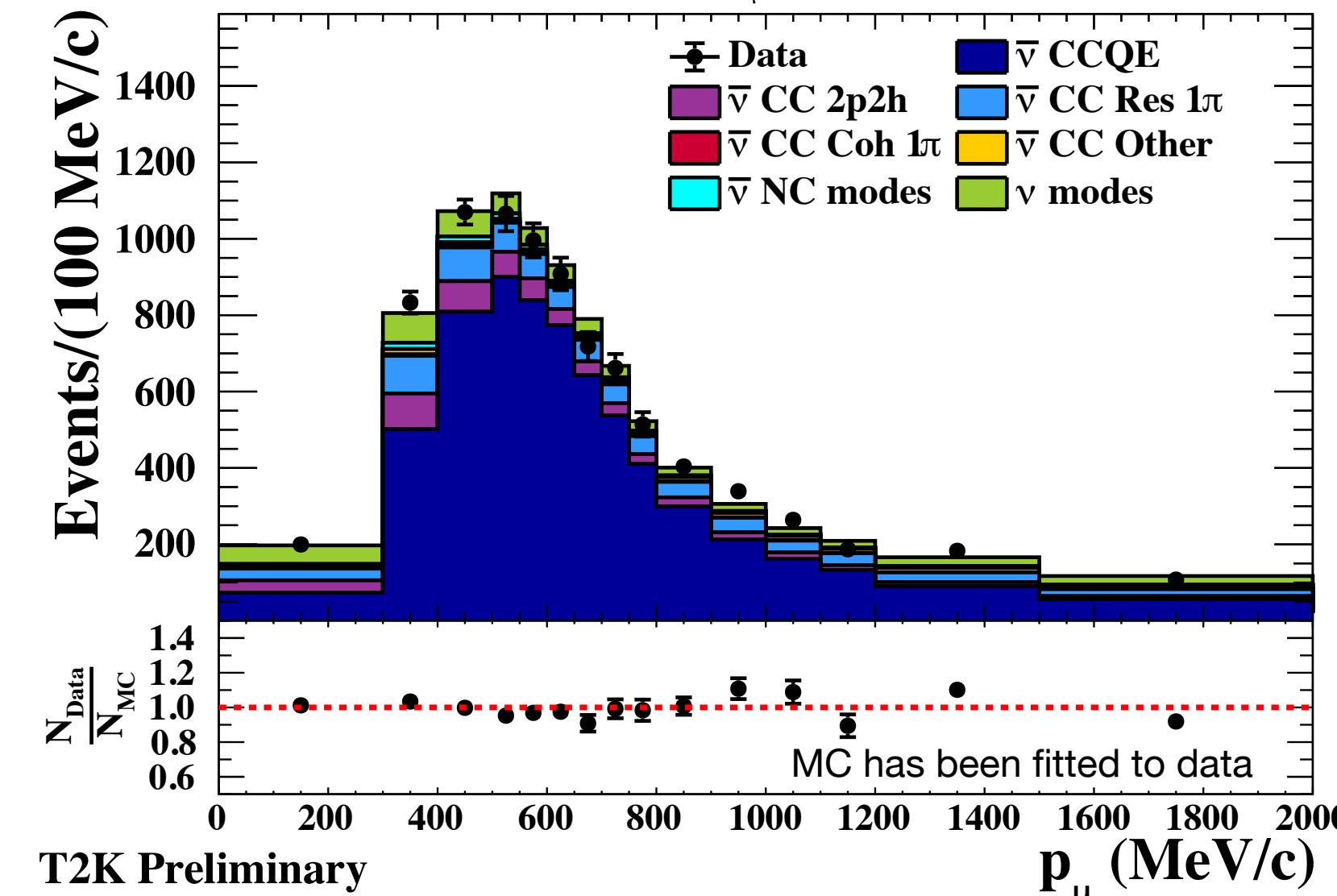
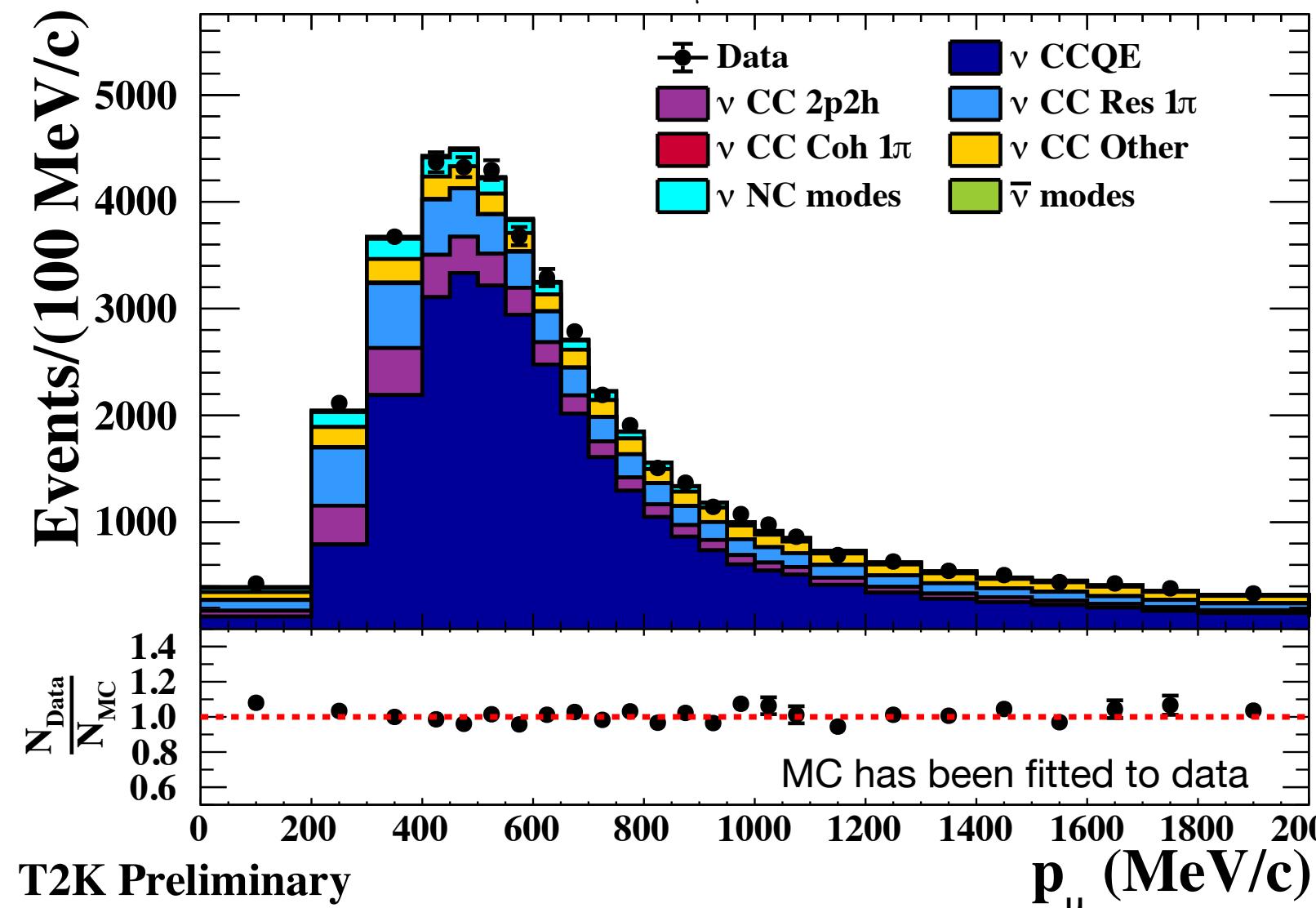
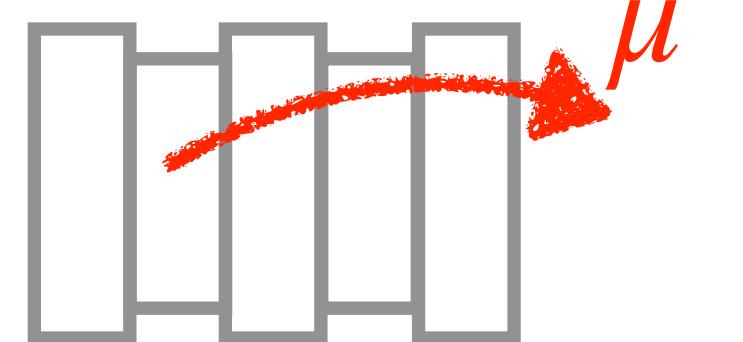
**$\nu$ -Mode Signal**  
 $\nu_\mu$  in  $\nu$ -Mode



**$\bar{\nu}$ -Mode Signal**  
 $\bar{\nu}_\mu$  in  $\bar{\nu}$ -Mode



**$\bar{\nu}$ -Mode Wrong-Sign  
Background**  
 $\nu_\mu$  in  $\bar{\nu}$ -Mode

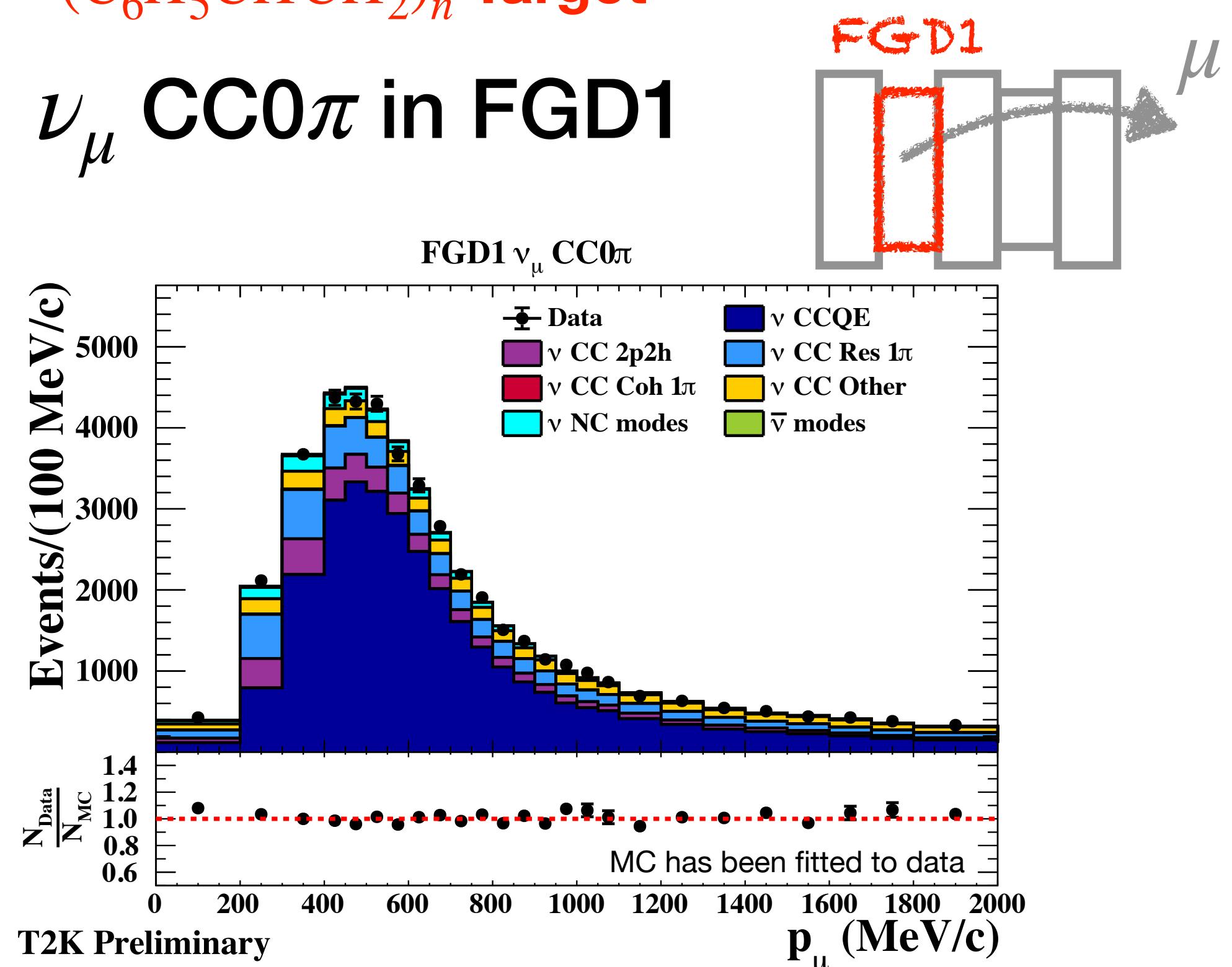


# ND280 Analysis Samples

- $18 = 3 \times (1 + 2) \times 2$  classes of neutrino interactions defined at ND280, split by:
  - **2 Detectors (Target Materials):** FGD1 (C) and FGD2 (C+O)

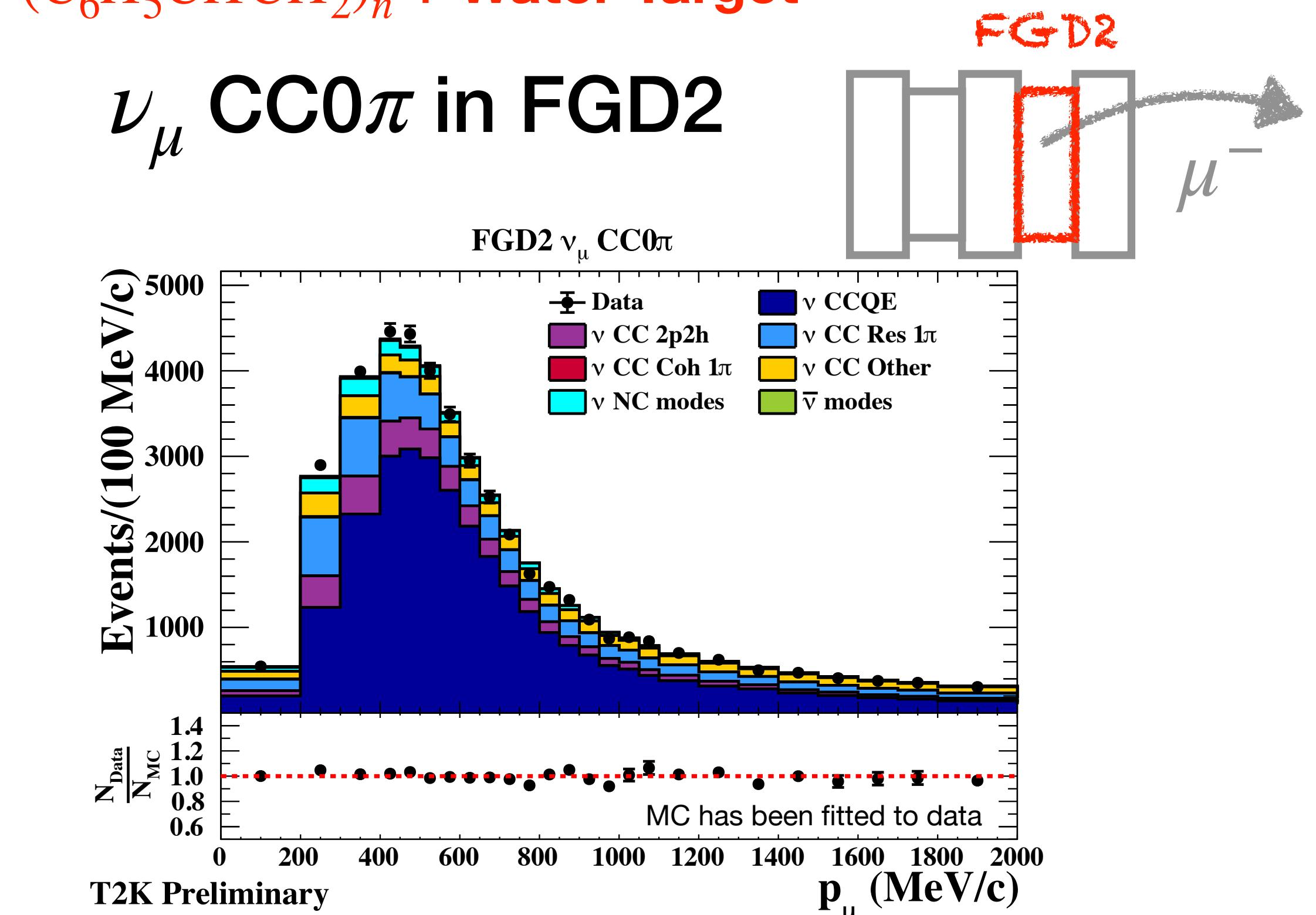
**Polystyrene Scintillator**  
 $(C_6H_5CHCH_2)_n$  Target

$\nu_\mu$  CC0 $\pi$  in FGD1



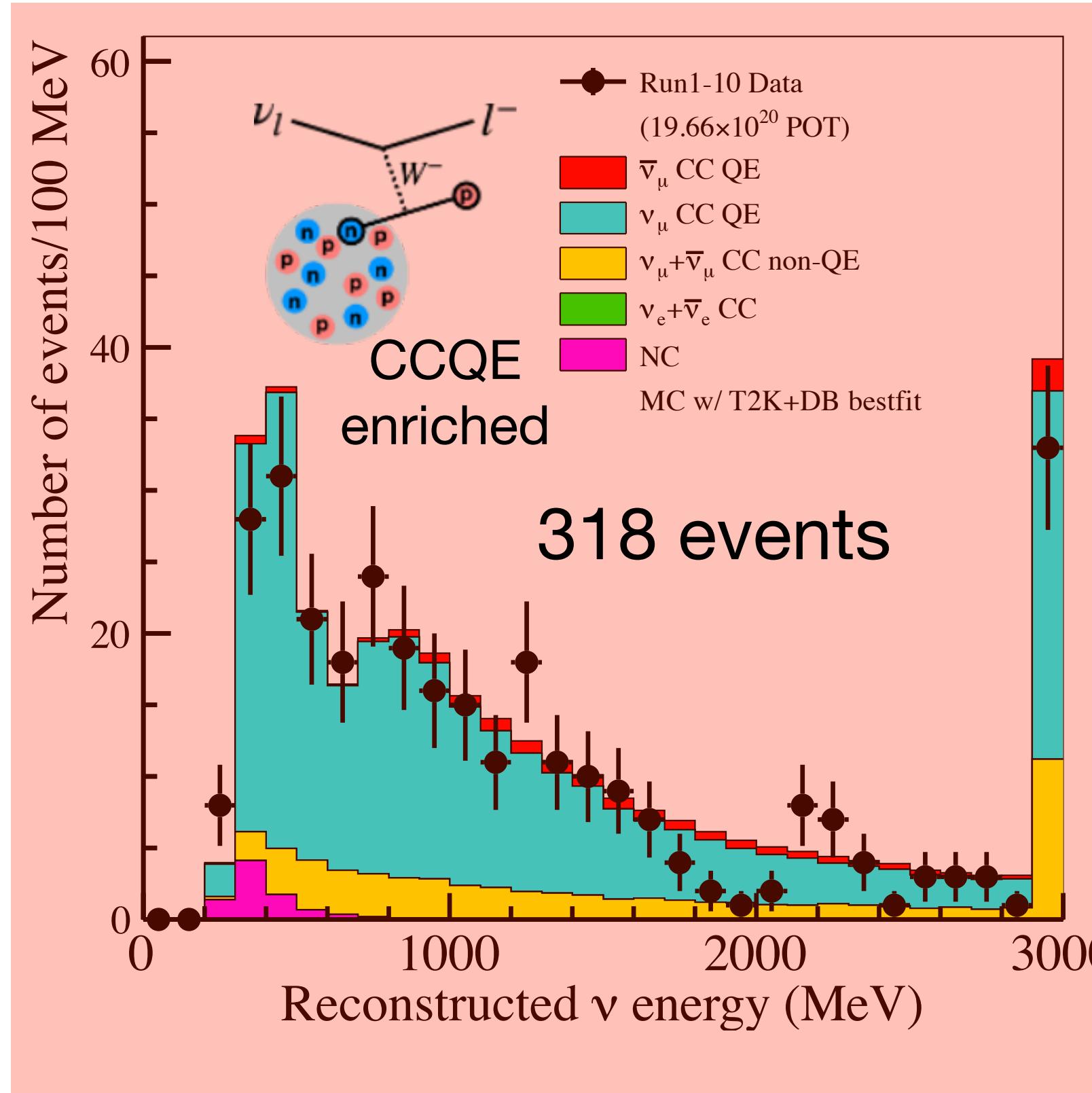
**Polystyrene Scintillator**  
 $(C_6H_5CHCH_2)_n +$  Water Target

$\nu_\mu$  CC0 $\pi$  in FGD2

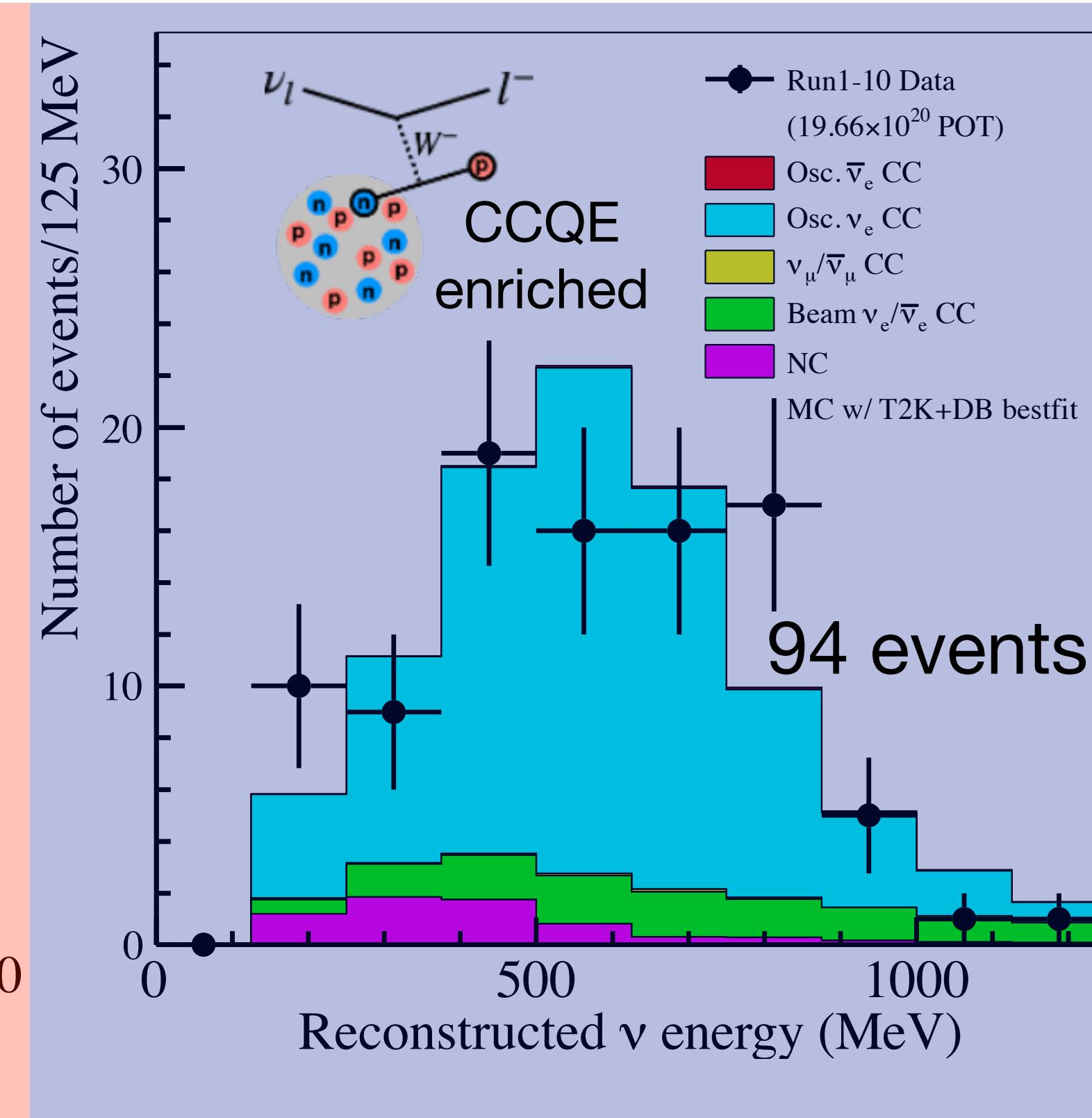


# Super-K data samples: neutrino mode

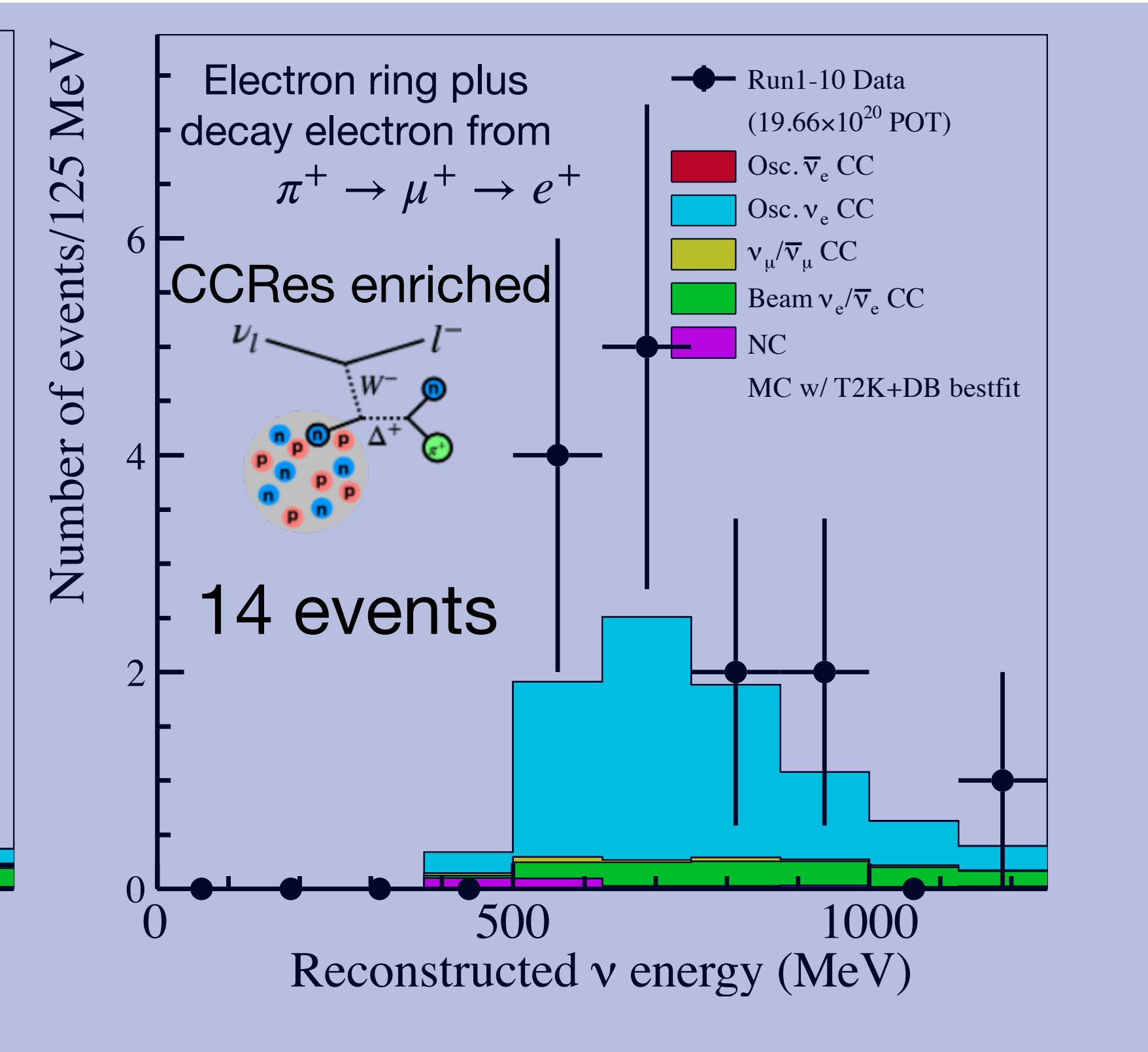
CC0 $\pi$  muon-like:  
1 muon-like ring



CC0 $\pi$  electron-like:  
1 electron-like ring



CC1 $\pi$  electron-like:  
1 electron-like ring + 1 Michel electron

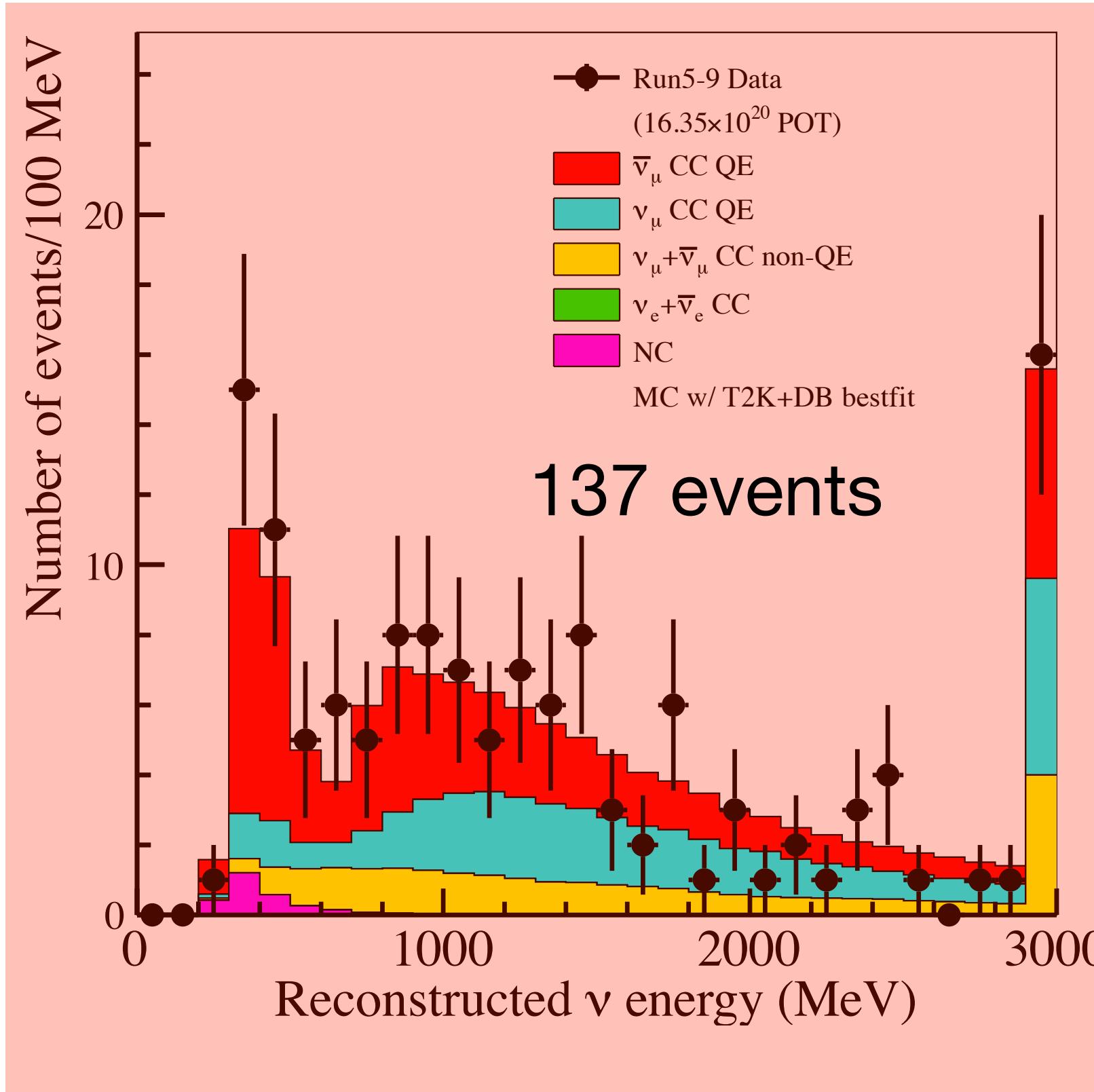


Direct probe of muon neutrino survival

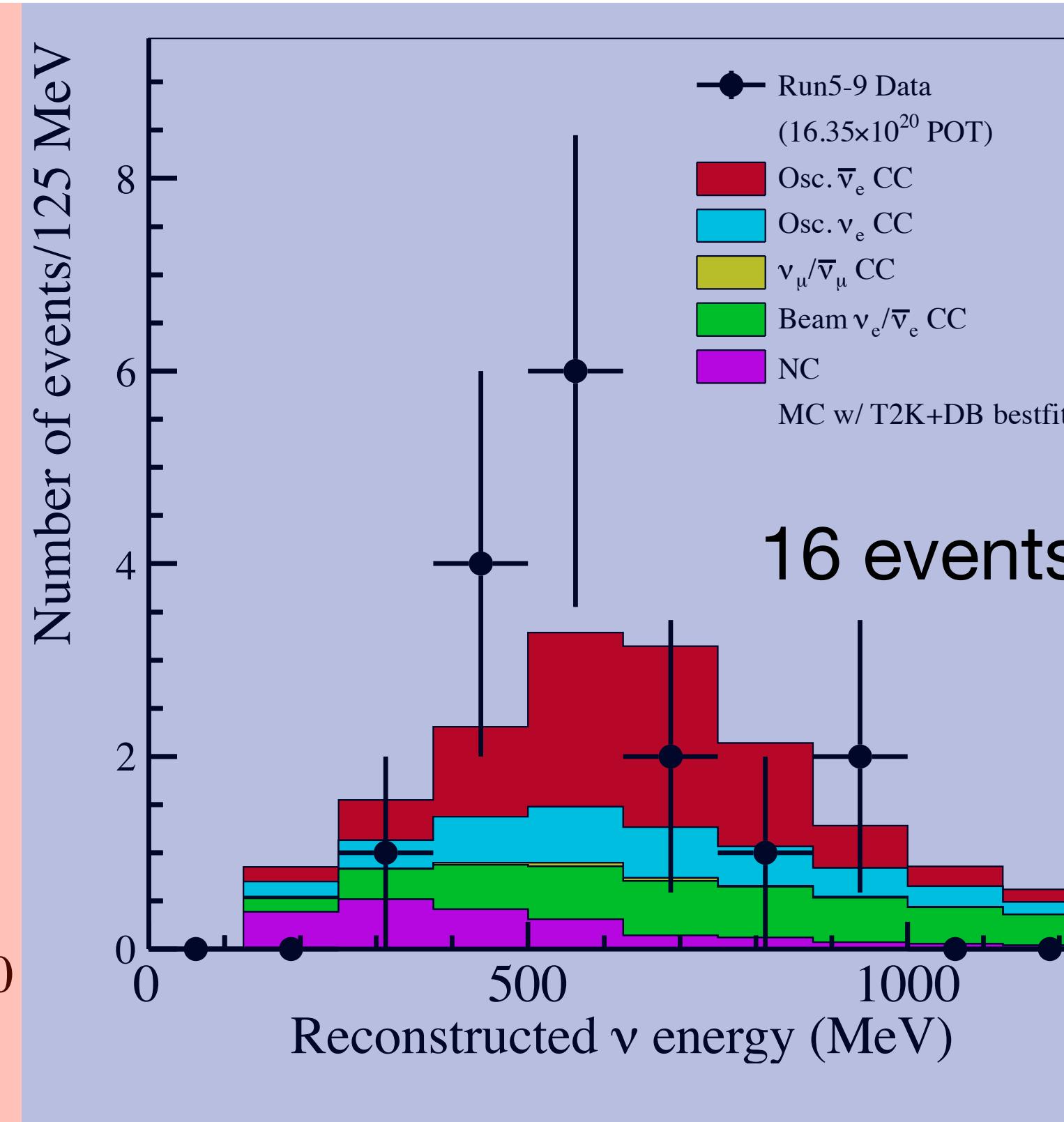
Direct probes of electron neutrino  
appearance

# Super-K data samples: anti-neutrino mode

CC0 $\pi$  muon-like:  
1 muon-like ring



CC0 $\pi$  electron-like:  
1 electron-like ring

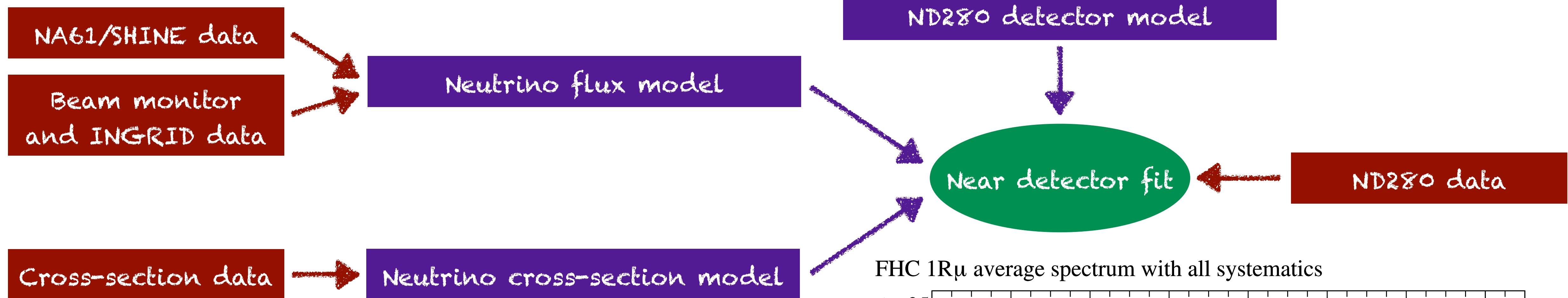


No CC1 $\pi$  electron-like sample in  
anti-neutrino mode, because of  
the high absorption of  $\pi^-$  in water

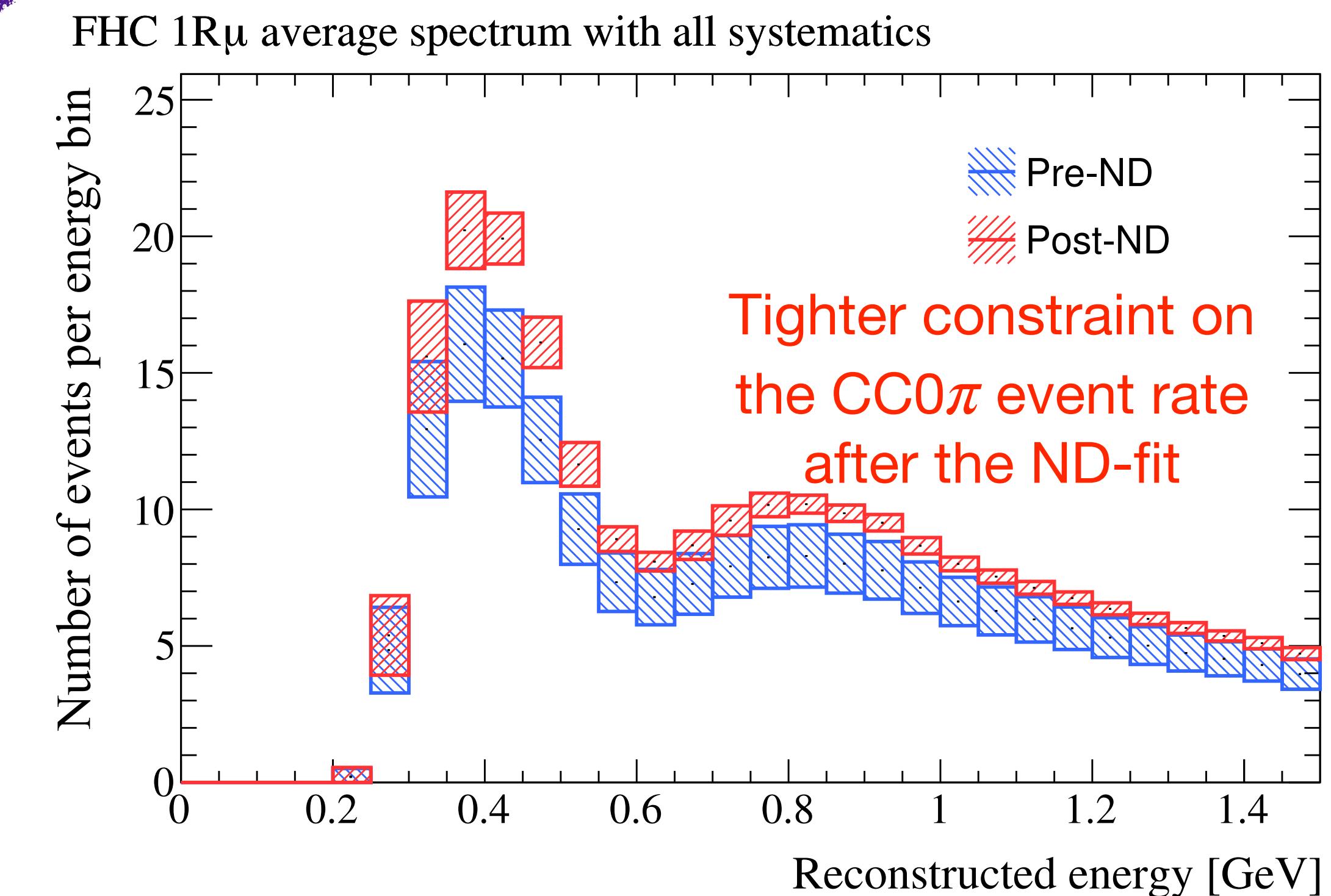
Direct probe of muon neutrino survival

Direct probe of  
electron neutrino appearance

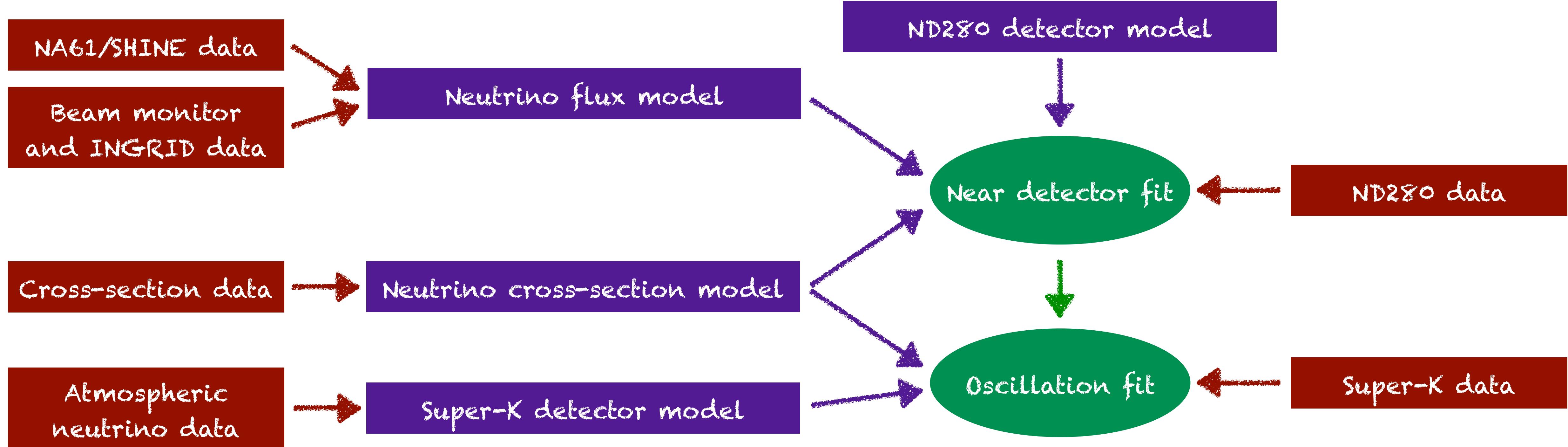
# Analysis strategy: near detector fit



- Fits to near detector data dramatically improve our predictive capabilities at Super-K (e.g. from **13.0%** to **4.7%** systematic uncertainty on the pre-fit CC $0\pi$  electron neutrino event rate)



# Analysis strategy: full overview



- Frequentist oscillation analyses: first fit to near detector data, then fit to Super-K data

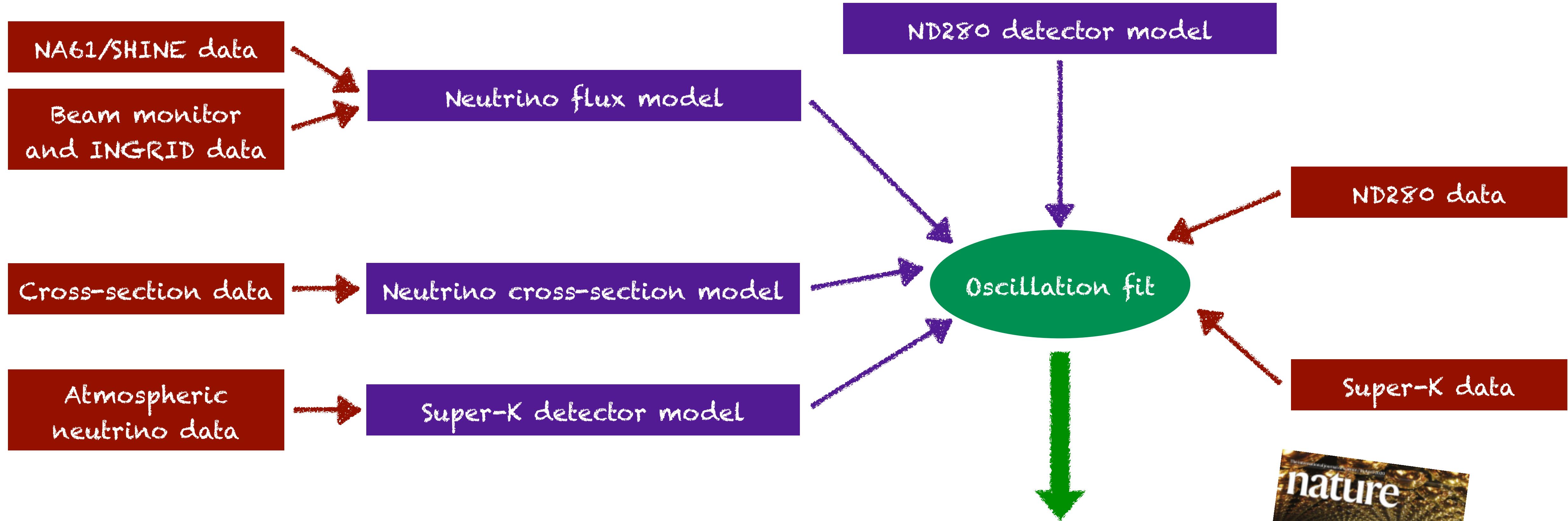
T2K oscillation result!

$$\Delta m_{32}^2, \sin^2 \theta_{23}, \delta_{CP}$$

For a recent published result  
check out: [Nature 580, 339–344 \(2020\)](#)



# Analysis strategy: full overview



- **MaCh3 (Bayesian) oscillation analysis:  
simultaneous fit to near and far detector data**
- T2K produces consistently converging results from  
the simultaneous and sequential fitting approaches

T2K oscillation result!

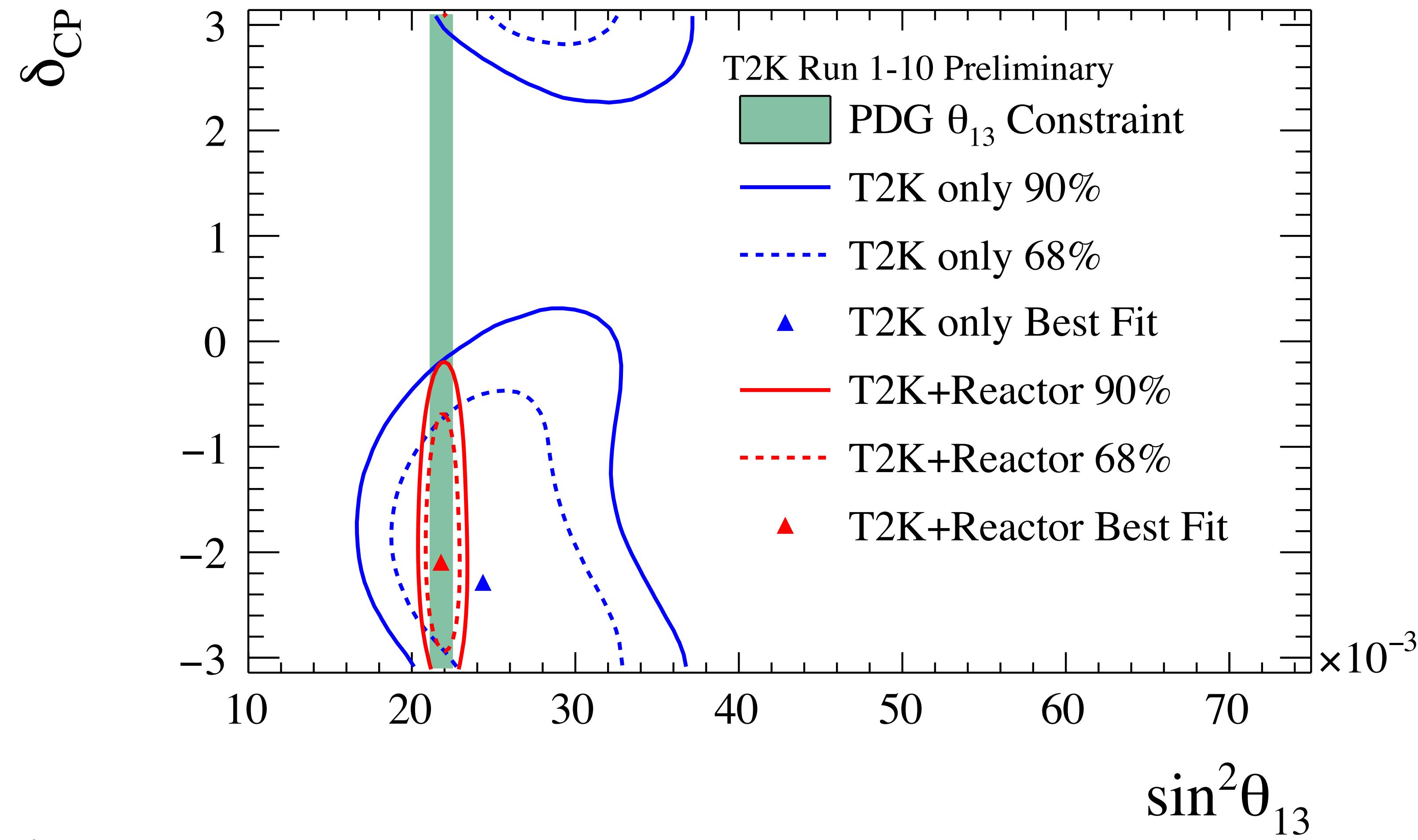
$$\Delta m_{32}^2, \sin^2 \theta_{23}, \delta_{CP}$$

For a recent published result  
check out: [Nature 580, 339–344 \(2020\)](#)



# $\delta_{CP}$ constraint, and external inputs from reactor neutrino experiments

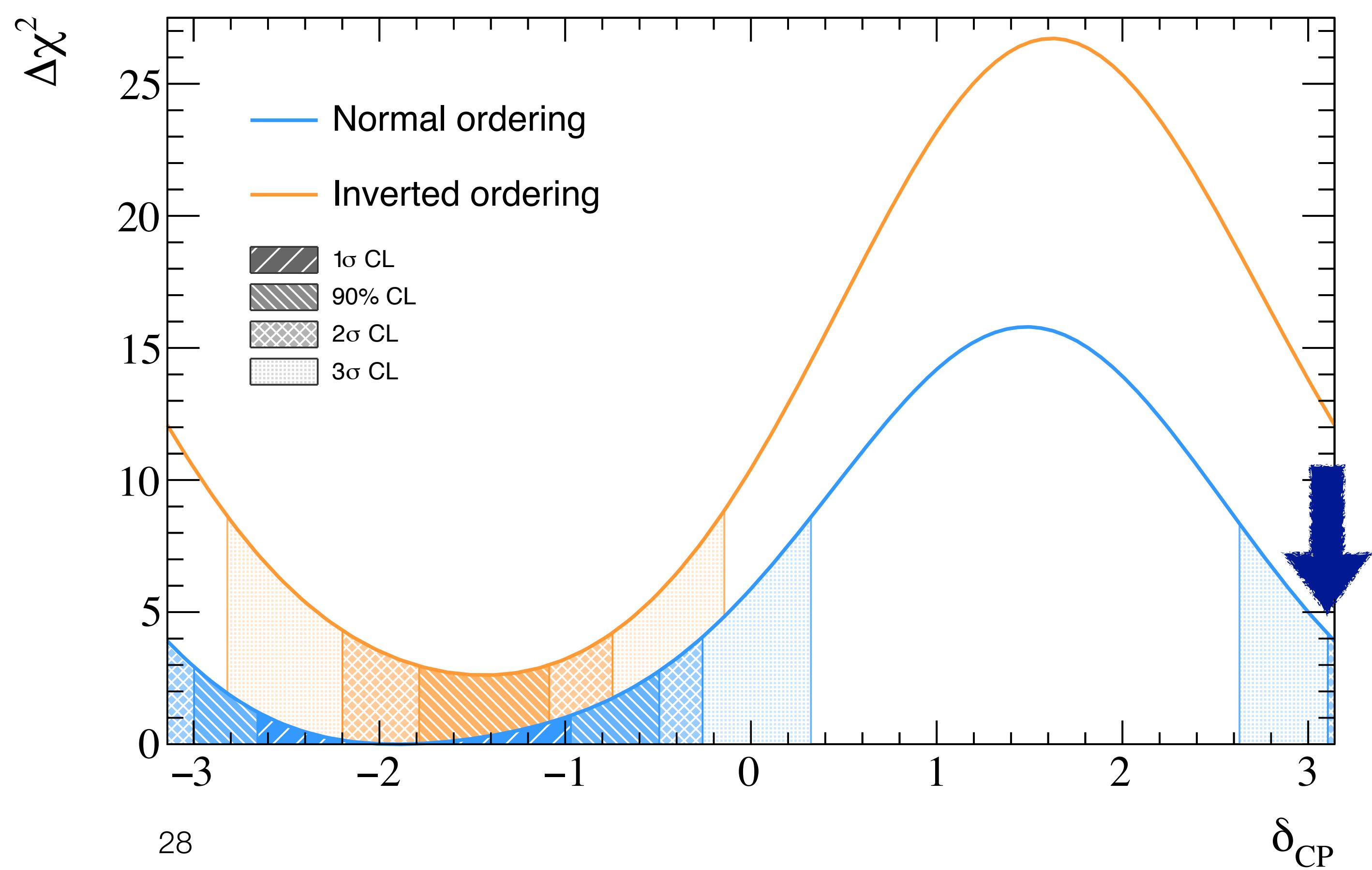
- We measure  $\delta_{CP}$  both with and without the constraint on  $\theta_{13}$  from reactor neutrino experiments (2019 PDG Constraint)
- The **tightest  $\delta_{CP}$  constraint from T2K incorporates the  $\theta_{13}$  measurement from combined reactor data**, which will be assumed in all subsequent plots



# $\delta_{CP}$ constraint

- CP conservation excluded at 90% confidence level, with a preference for maximal CP violation ( $\delta_{CP} = -\pi/2$ )
- $\delta_{CP} = \pi$  not quite excluded at  $2\sigma$
- Marginalised over both mass hierarchies, 35% of  $\delta_{CP}$  values excluded at  $3\sigma$
- Largest  $\Delta\chi^2$  change from a series of robustness studies (cross section modelling bias) would cause the left (right) edge of the 90% CL region to move by 0.073 (0.080)

Confidence level	Interval (NH)	Interval (IH)
$1\sigma$	$[-2.66, -0.97]$	
90%	$[-3.00, -0.49]$	$[-1.79, -1.09]$
$2\sigma$	$[-\pi, -0.26] \cup [3.11, \pi]$	$[-2.20, -0.75]$
$3\sigma$	$[-\pi, 0.32] \cup [2.63, \pi]$	$[-2.82, -0.14]$

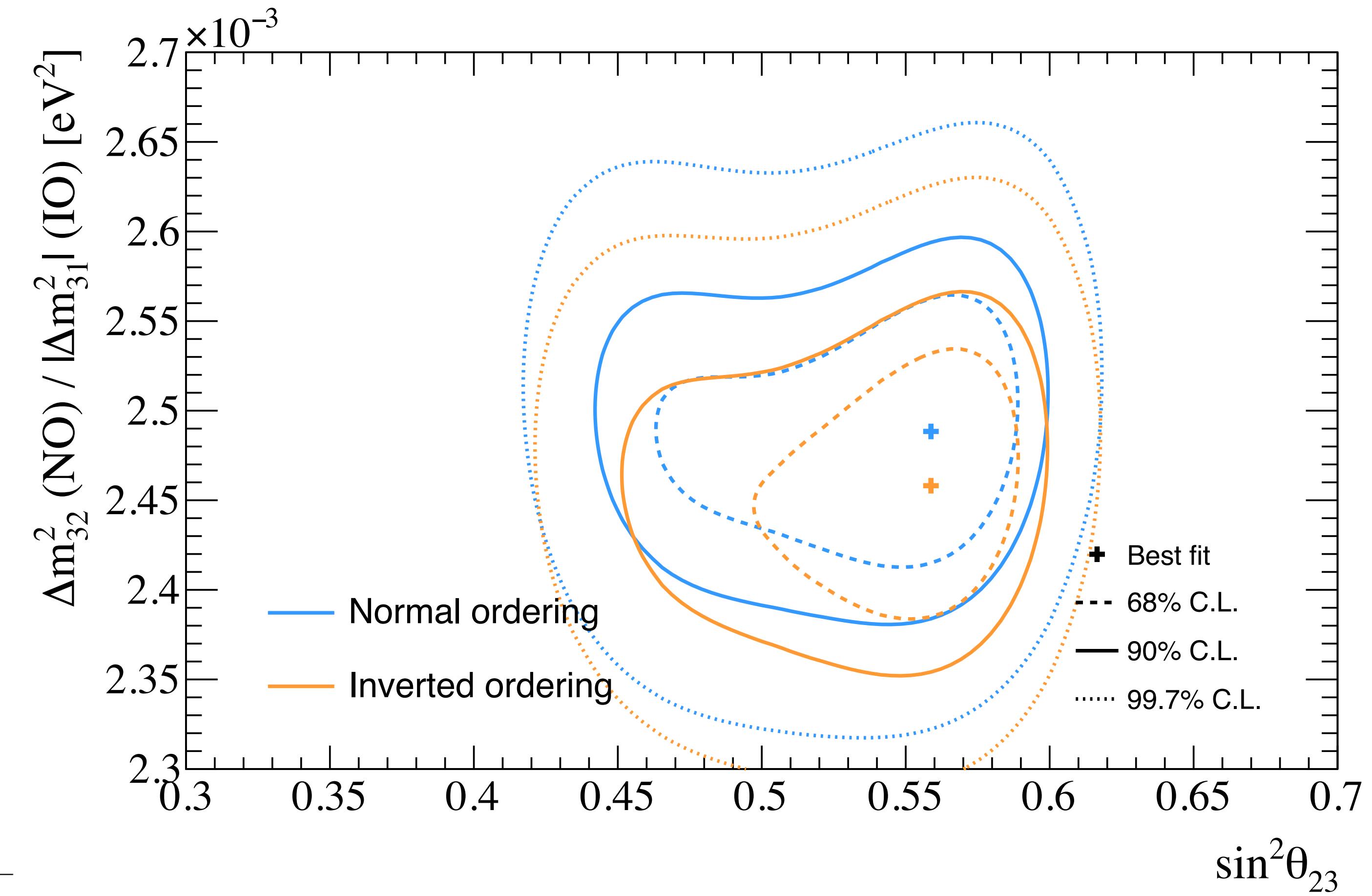


# Result in atmospheric ( $\Delta m_{32}^2$ , $\sin^2 \theta_{23}$ ) sector

- Fitting to all 5 samples at Super-K
- **Upper octant (normal ordering)** preferred with Bayes factor **3.4 (4.2)**
- No definitive conclusions can be drawn yet: Bayes factors within  $(10^{\frac{1}{2}}, 10)$  are classified as either “substantial” (Jeffreys scale) or “positive” (Kass and Raferty scale)

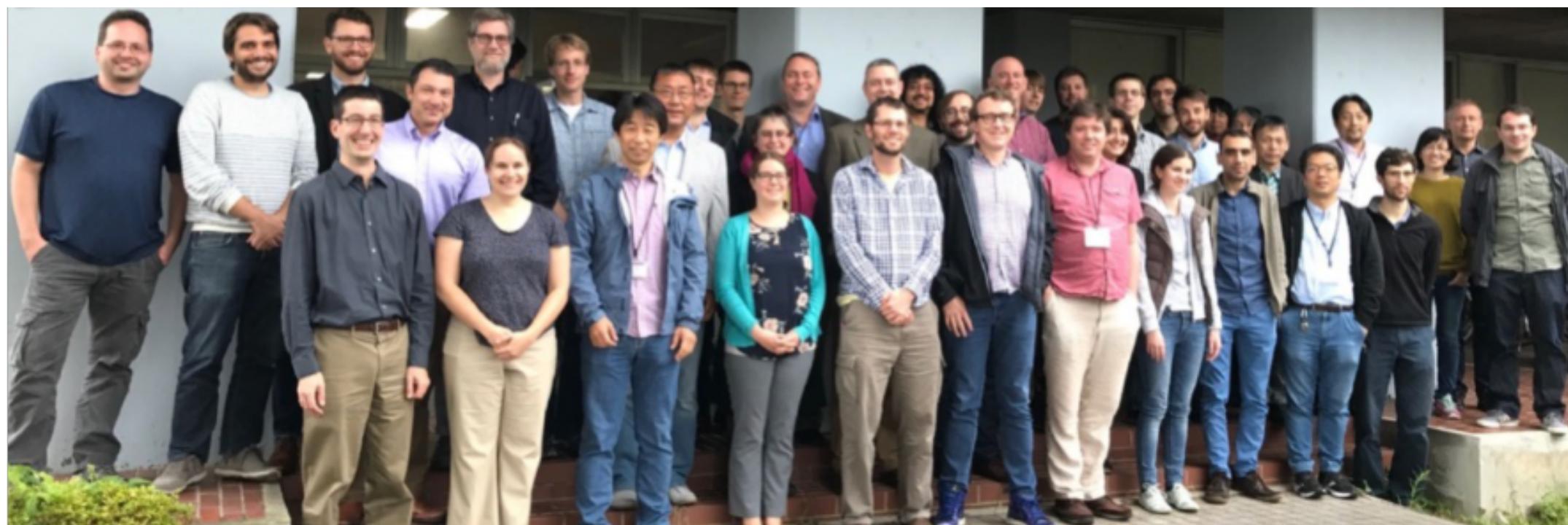
	$\sin^2 \theta_{23} < 0.5$	$\sin^2 \theta_{23} > 0.5$	Sum
NH ( $\Delta m_{32}^2 > 0$ )	0.195	0.613	0.808
IH ( $\Delta m_{32}^2 < 0$ )	0.035	0.157	0.192
Sum	0.230	0.770	1.000

	$\sin^2 \theta_{23}$	$\Delta m_{32}^2 (\times 10^{-3}) \text{ eV}^2$	$\sin^2 \theta_{23}$
2D best fit	0.488	2.46	
68% C.I. ( $1\sigma$ ) range	0.470 – 0.550	2.416 – 2.544 & –2.568 – –2.496	
90% C.I. range	0.447 – 0.580	2.376 – 2.584 & –2.616 – –2.436	



# Ongoing joint fits

- **T2K-NOvA and T2K-SuperK joint fits** are ongoing (agreements have been signed, and working groups formed)
- **Combining different detector/baseline/energy configurations resolves degeneracies**
- Expect an improved measurement sensitivity beyond a simple increase in statistics



T2K-NOvA meeting in J-PARC, Japan



T2K-NOvA meeting in Fermilab, USA

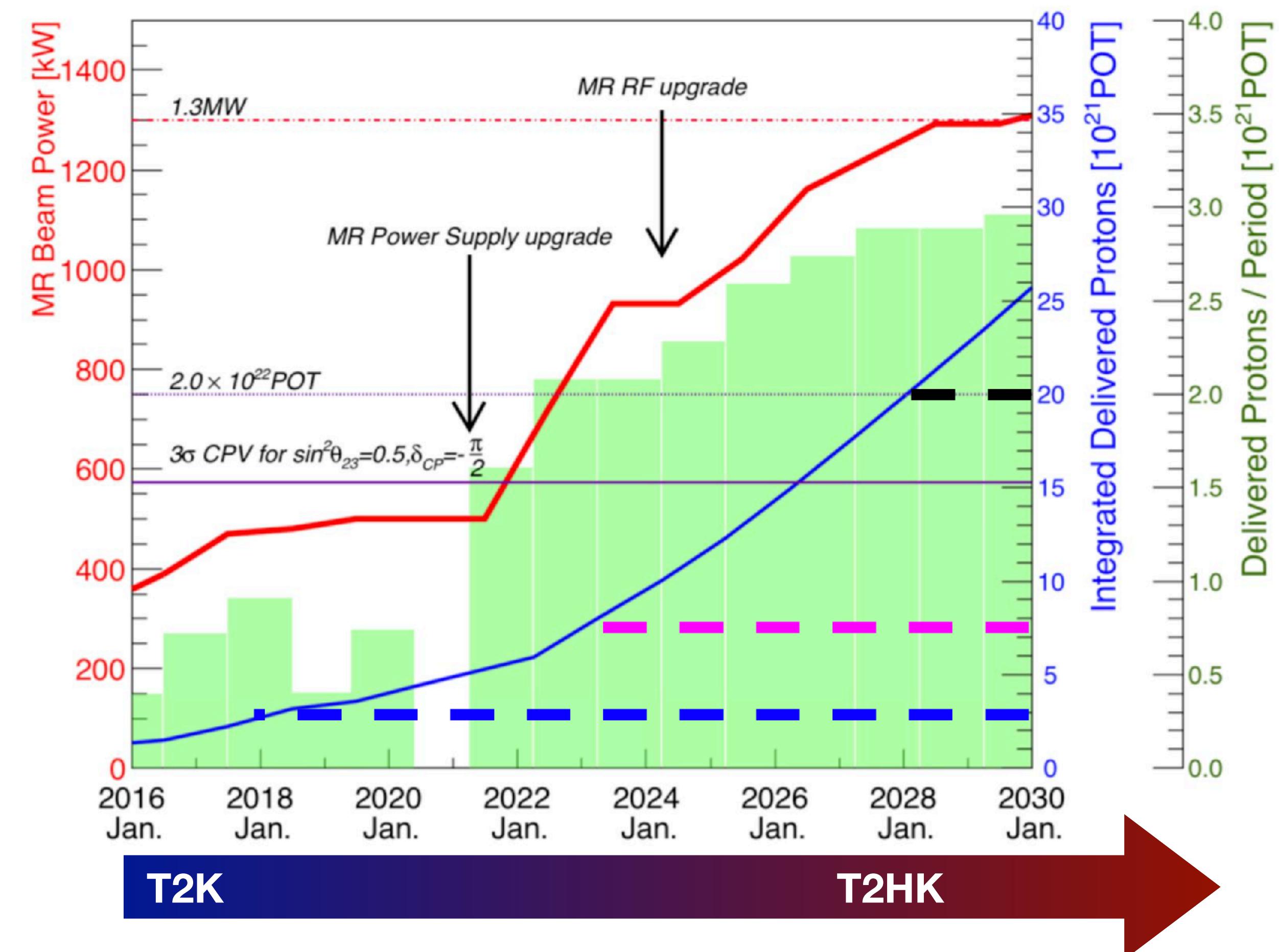
# Future of T2K

- T2K data taking has been extended until the start of Hyper-K, as per KEK's statement from 6 July 2020:



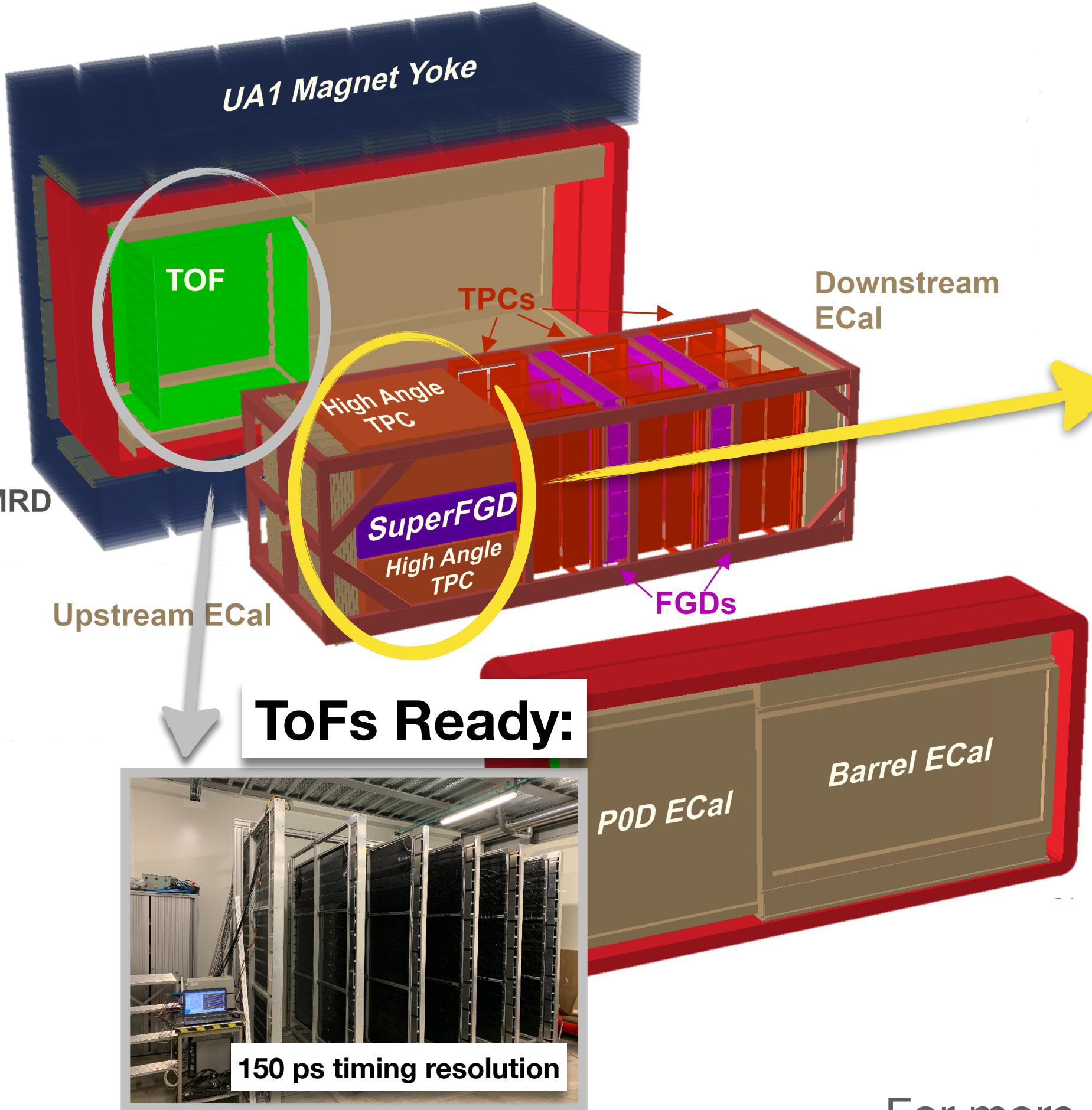
- KEK recognizes the importance of T2K operation with the upgraded near detector (ND280) for further advancement of CPV measurements. This will also strengthen international collaboration on T2K and HK, contribute to the success of the neutrino program in Japan, and enhance the physics program of HK with an upgraded beam intensity from the start of its operation.
- In light of the above, KEK will make its best effort to provide approximately four months of T2K beam operations per year until the start of HK.

- More statistics expected with beam power increase from ~500 kW towards 1.3 MW
- Reduced systematic uncertainties expected with the upgraded ND280 detector



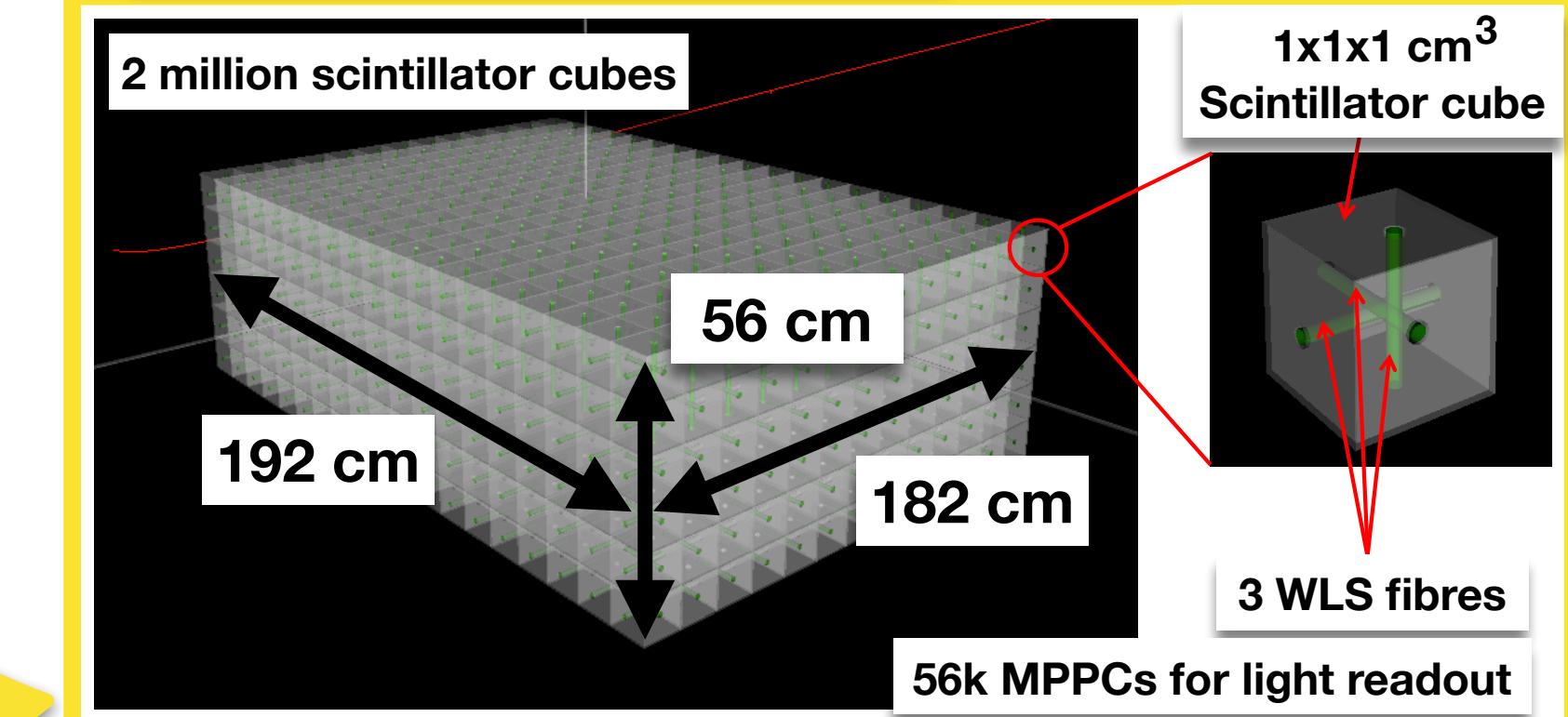
# T2K Upgrade

- **SuperFGD:** novel active target detector, surrounded by high angle time projection chambers (HTPCs) and enclosed with time of flight (ToF) detectors
- **Improved detector acceptance, efficiency, PID, timing resolution etc., feeding into reduction of cross section systematics**
- **Installation at J-PARC planned for the second half of 2022** (this includes a delay estimate due to Covid-19)

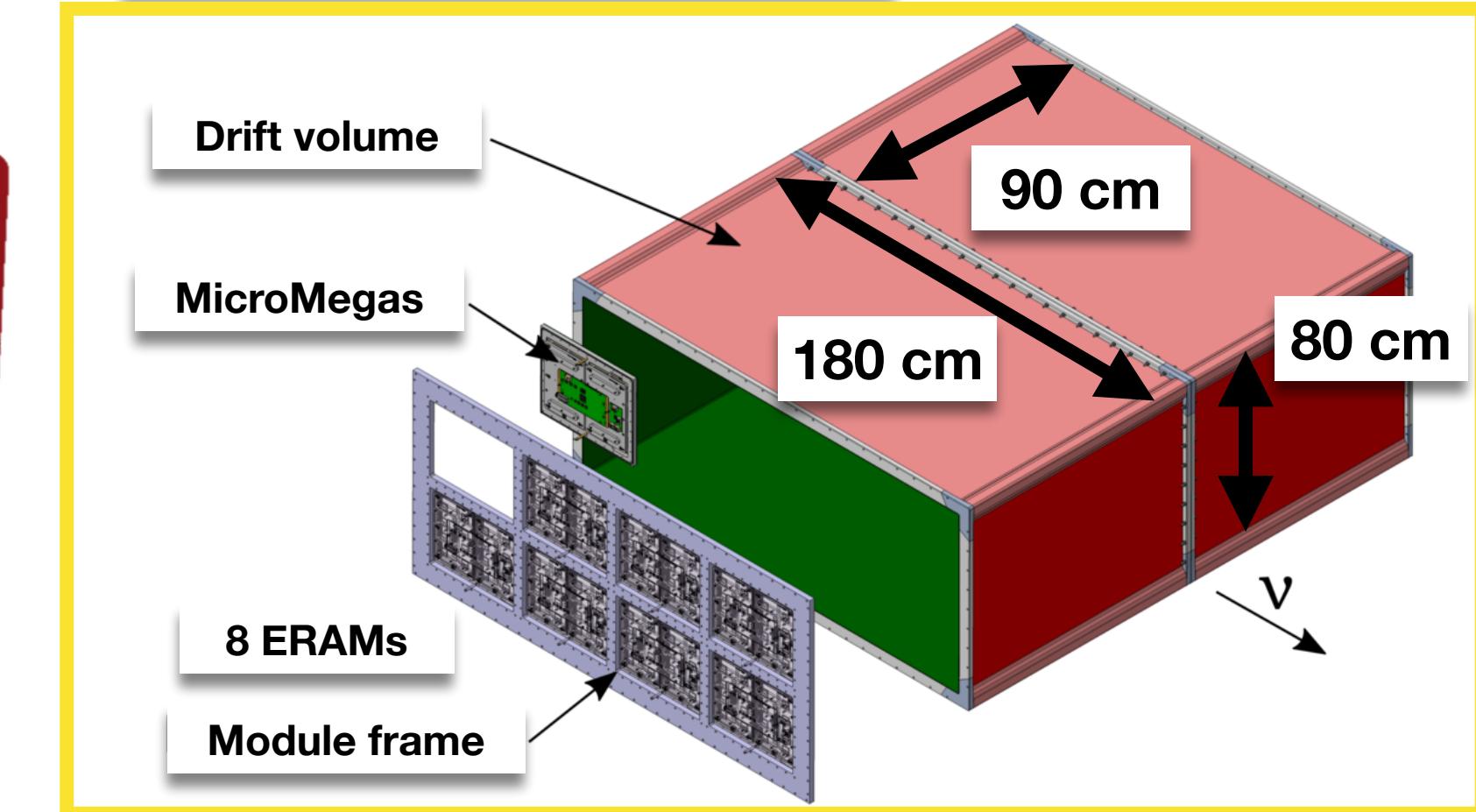


P0D sub-detector to be replaced with **SuperFGD**, **HTPCs** and **ToFs**

## SuperFGD Design:

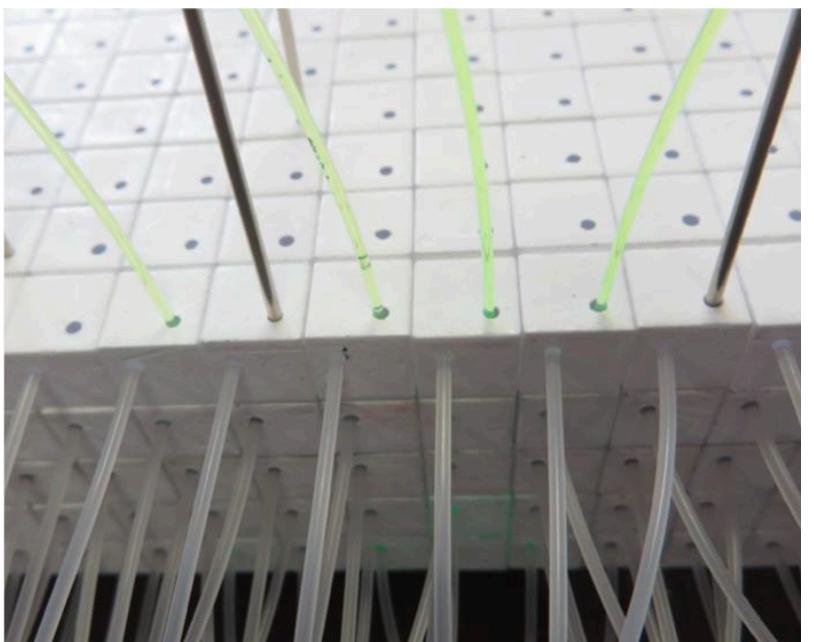
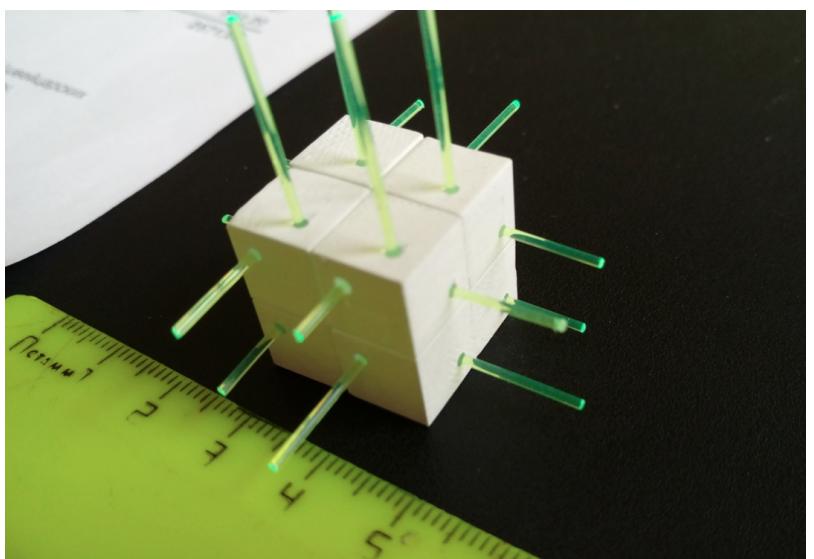


## 2 HTPCs Design:

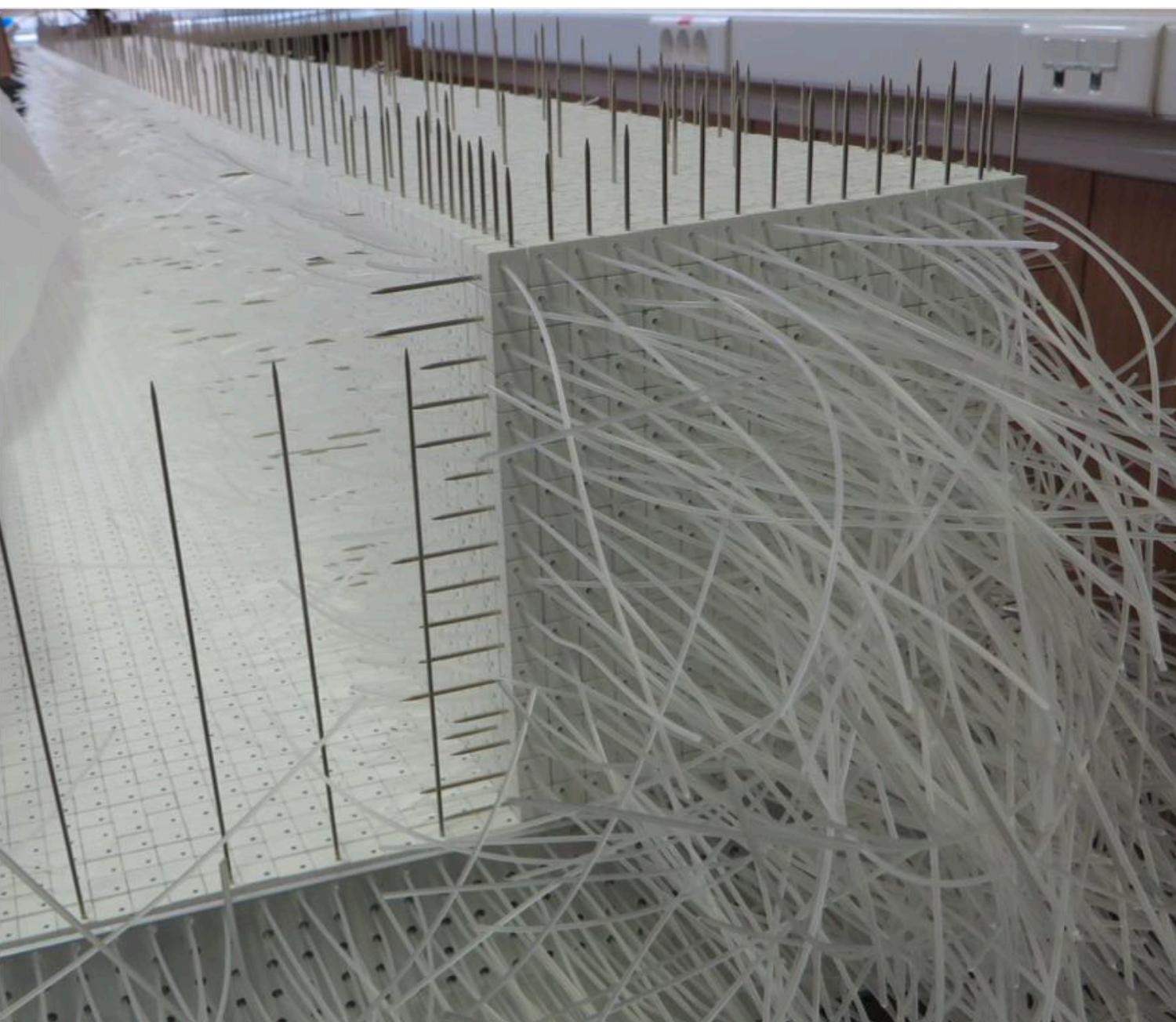


For more details refer to the ND280 Upgrade Technical Design Report, at [arXiv:1901.03750](https://arxiv.org/abs/1901.03750)

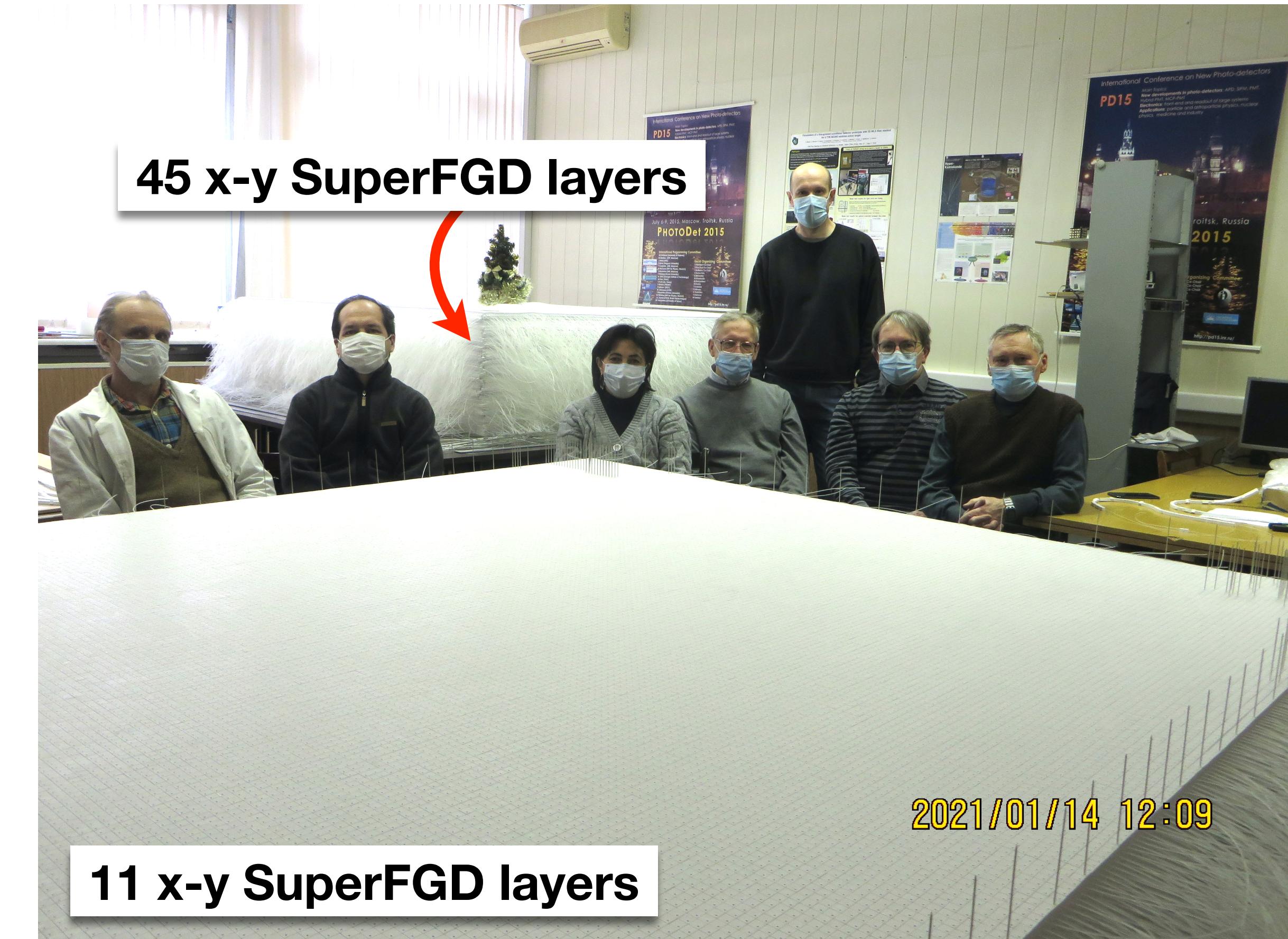
# T2K Upgrade: SuperFGD Preparation



- All 2.1 million cubes have been assembled into 56 x-y layers at **the Institute for Nuclear Research, Russia**

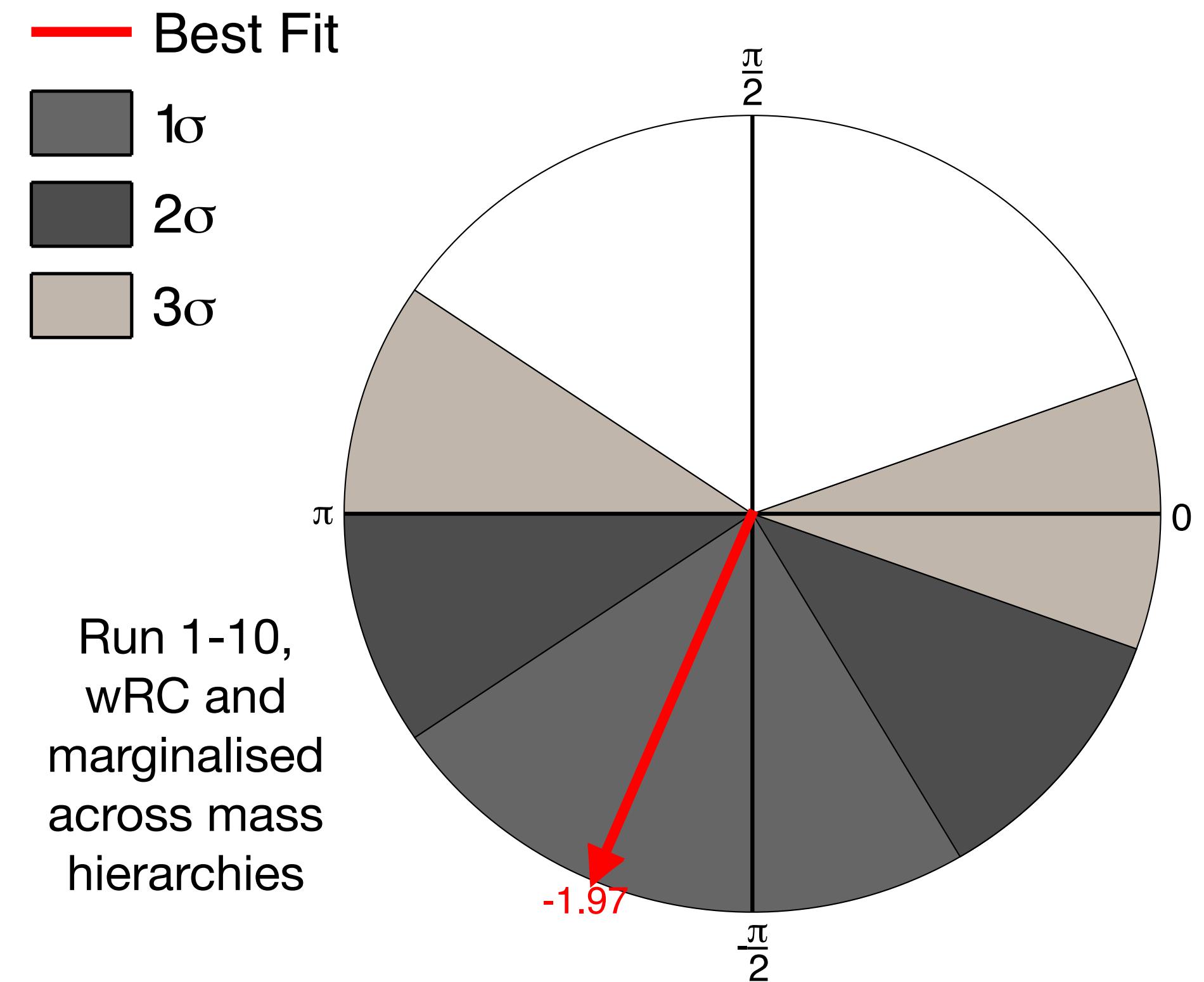


INR SuperFGD Assembly Team:



# Conclusions

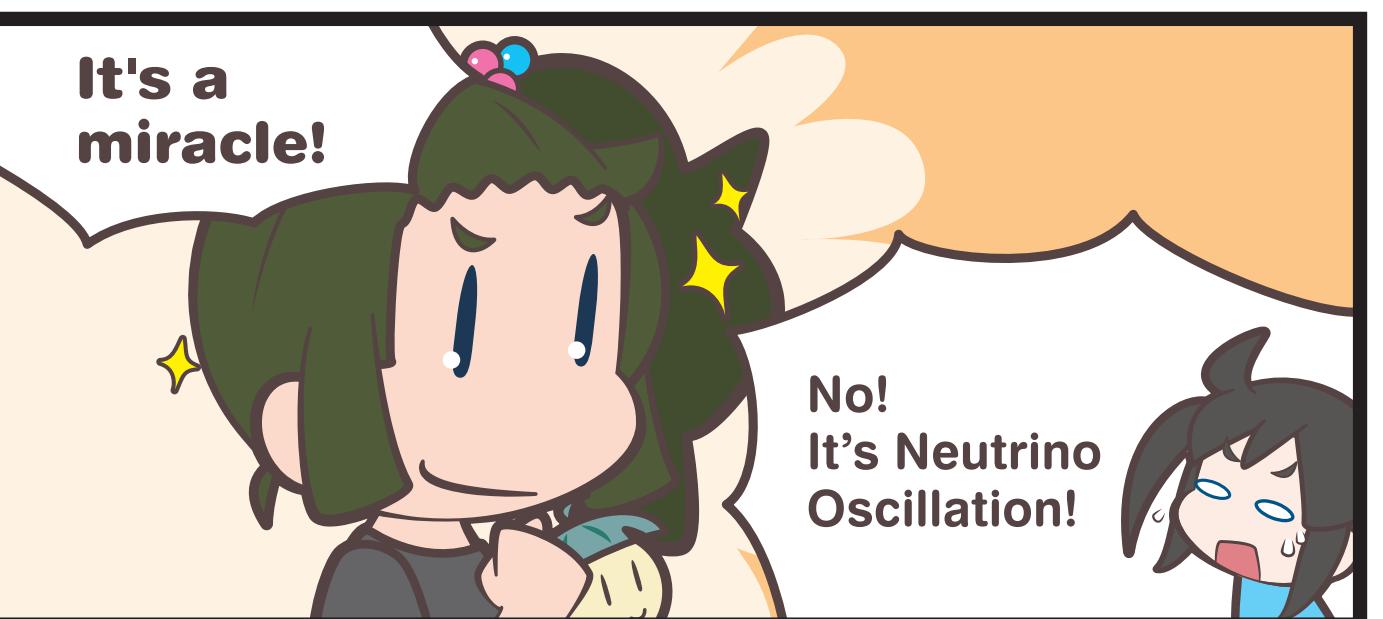
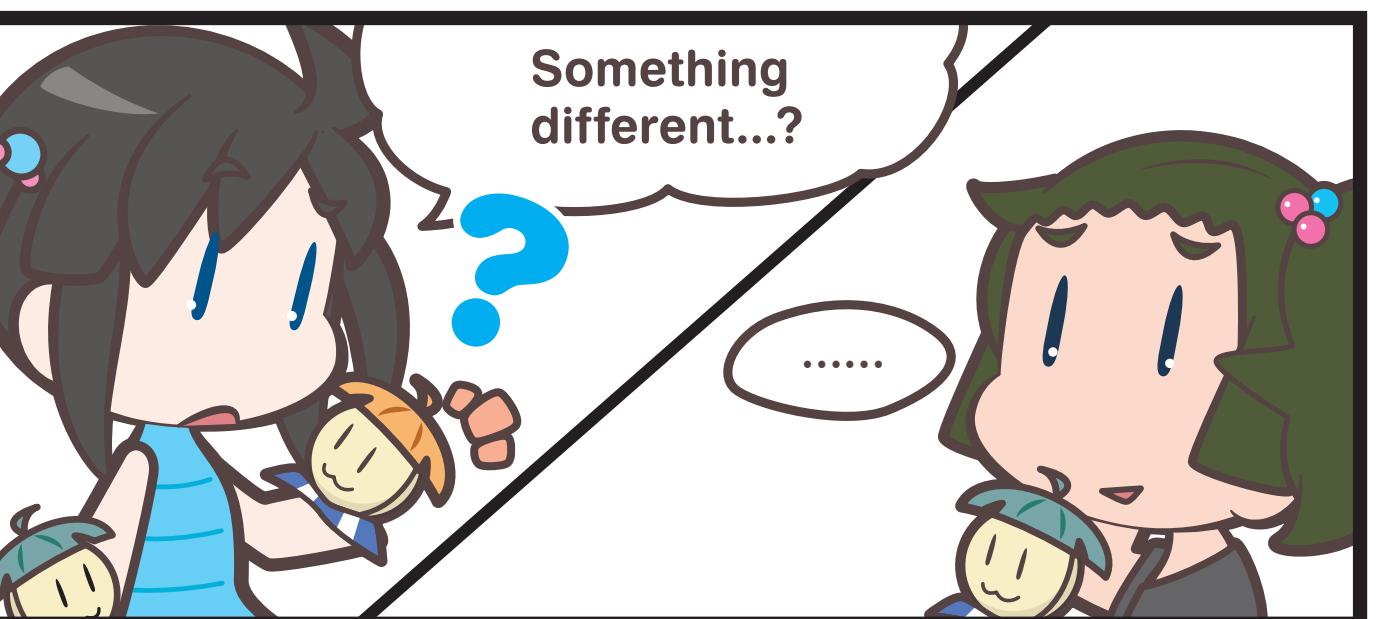
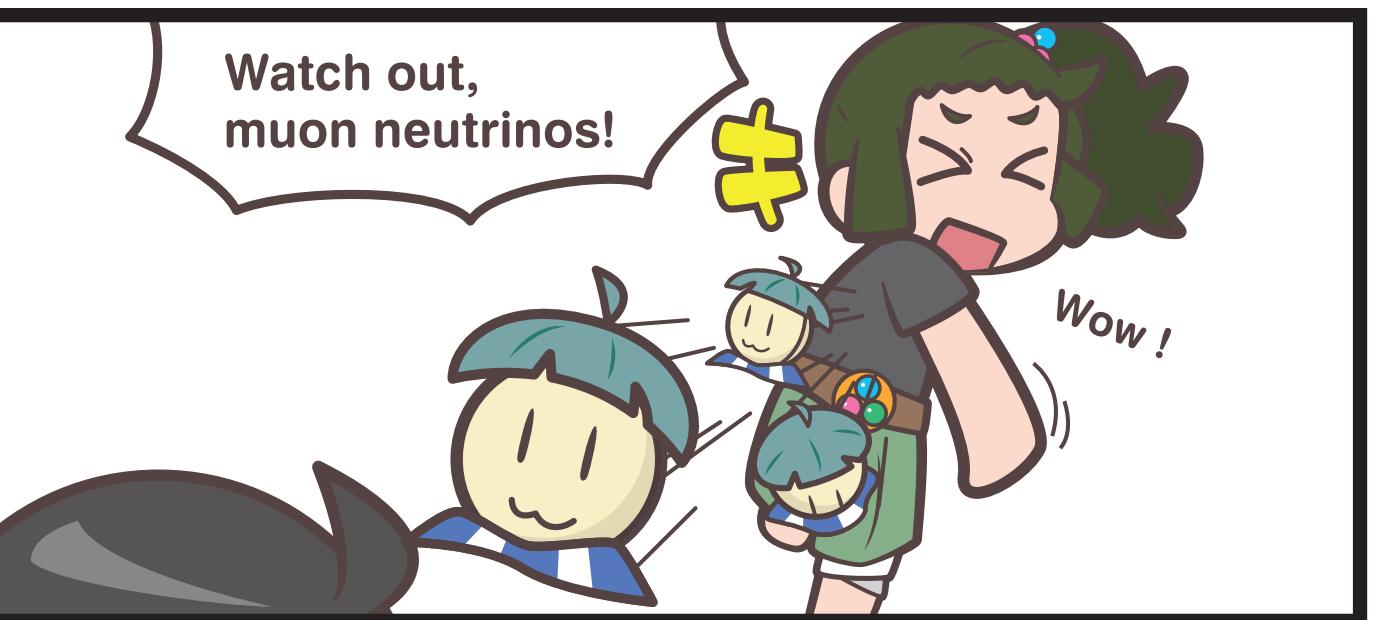
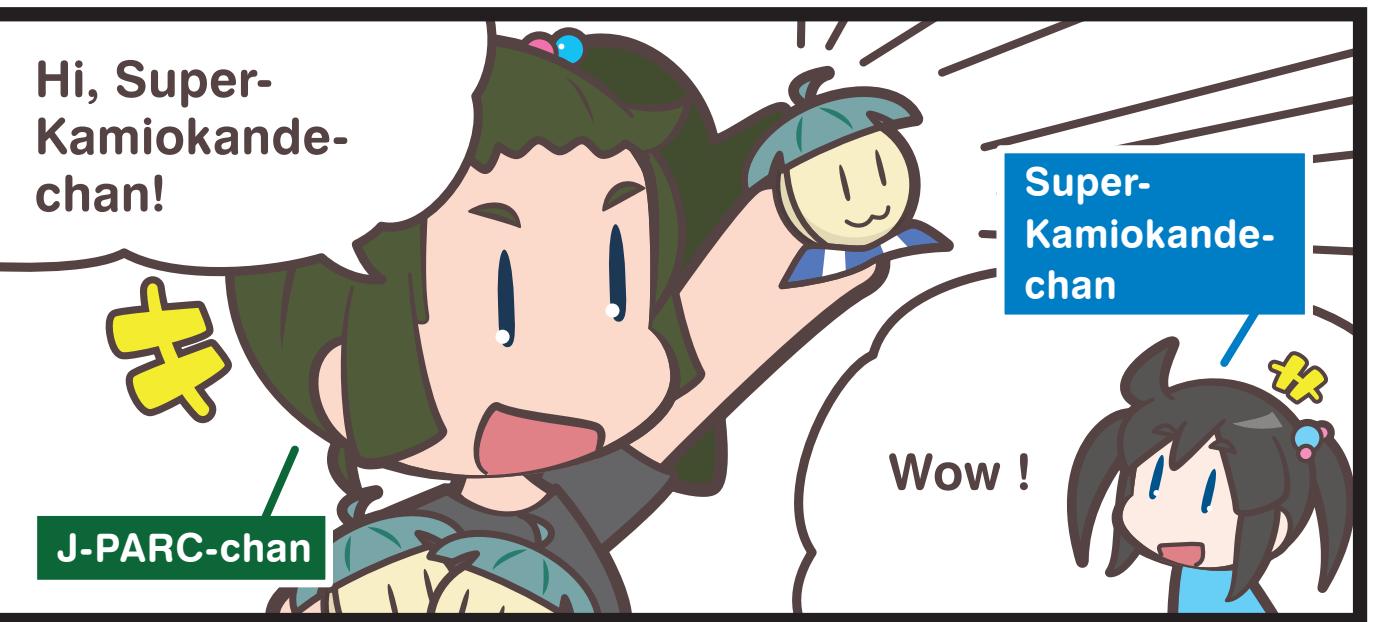
- Presented T2K's oscillation result for Runs 1-10 (data up to Feb 2020,  $3.6 \times 10^{21}$  POT), with 33% more POT in  $\nu$ -mode compared to earlier
- $\delta_{CP}$  conserving values excluded at 90% confidence level, with  $\delta_{CP} = \pi$  not quite excluded at  $2\sigma$
- Preference for maximum CP violation, normal mass ordering and upper  $\theta_{23}$  octant
- Ambitious upgrade of near detector planned for second half of 2022
- Future T2K data-taking has been approved until the start of Hyper-K



# Thank you!

## Any questions?

### T2K!



**J-PARC-chan:**  
lives in Tokai-mura, Naka-gun, Ibaraki, Japan.



**Super-Kamiokande-chan:**  
lives in Kamioka-cho, Hida-city, Gifu, Japan.



# Backup

**Tomislav Vladisavljevic**  
**RAL, UK**

On behalf of the T2K Collaboration

XX Lomonosov Conference on Elementary Particle Physics  
20 August 2021, Online

# Neutrino oscillations, general expression

- Neutrino oscillations occur because of different flavour eigenstates (interacting states) and mass eigenstates (propagating states)

$$\begin{aligned}
 \theta_{23} &\approx 48^\circ & \theta_{13} &\approx 8^\circ & \theta_{12} &\approx 34^\circ \\
 |\Delta m_{32}^2| &\approx |\Delta m_{31}^2| \approx 2.5 \times 10^{-3} \text{ eV}^2 & \delta_{CP} &= ? & \Delta m_{21}^2 &\approx 7.5 \times 10^{-5} \text{ eV}^2 \\
 s_{ij} &= \sin \theta_{ij} & c_{ij} &= \cos \theta_{ij}
 \end{aligned}$$

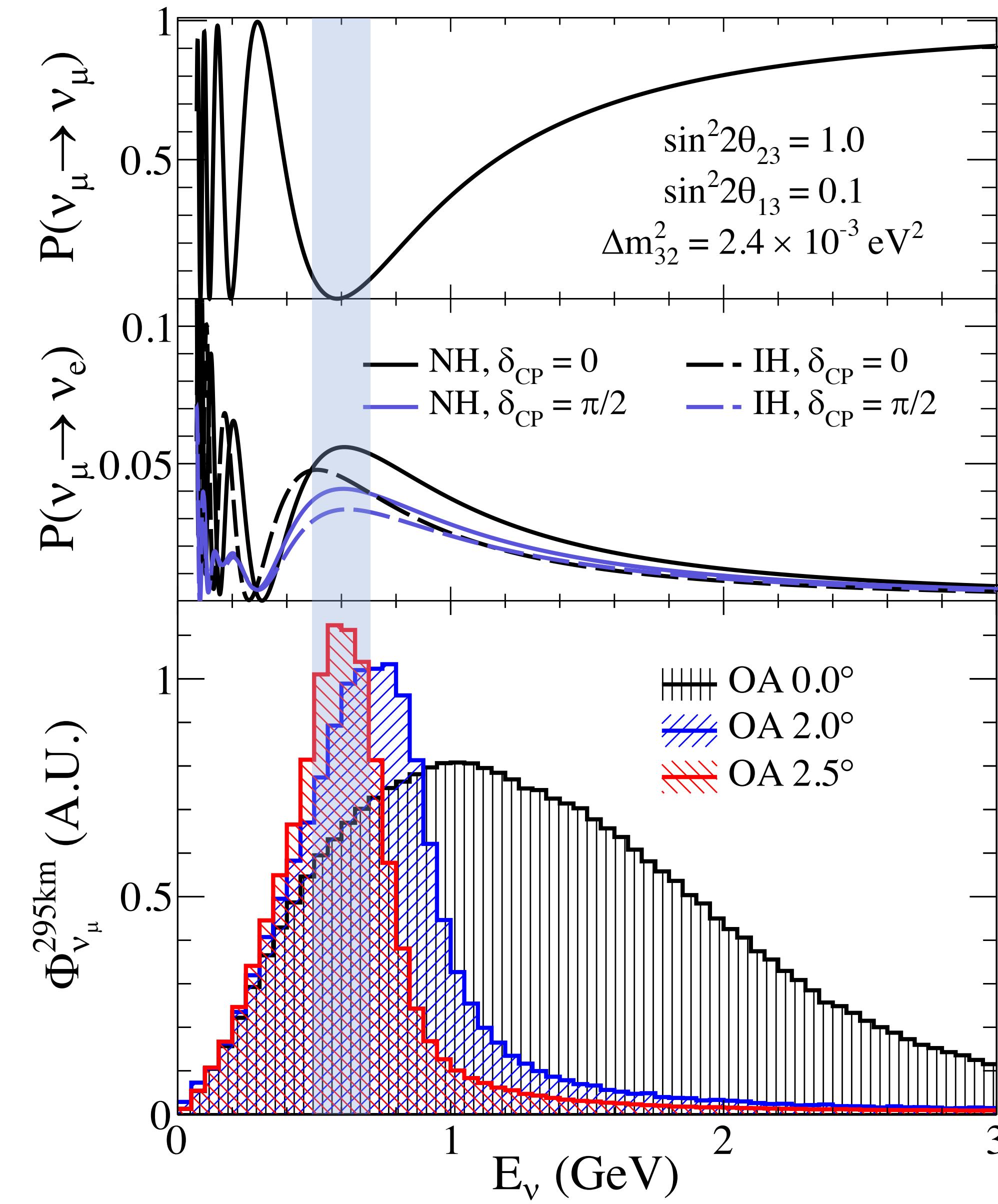
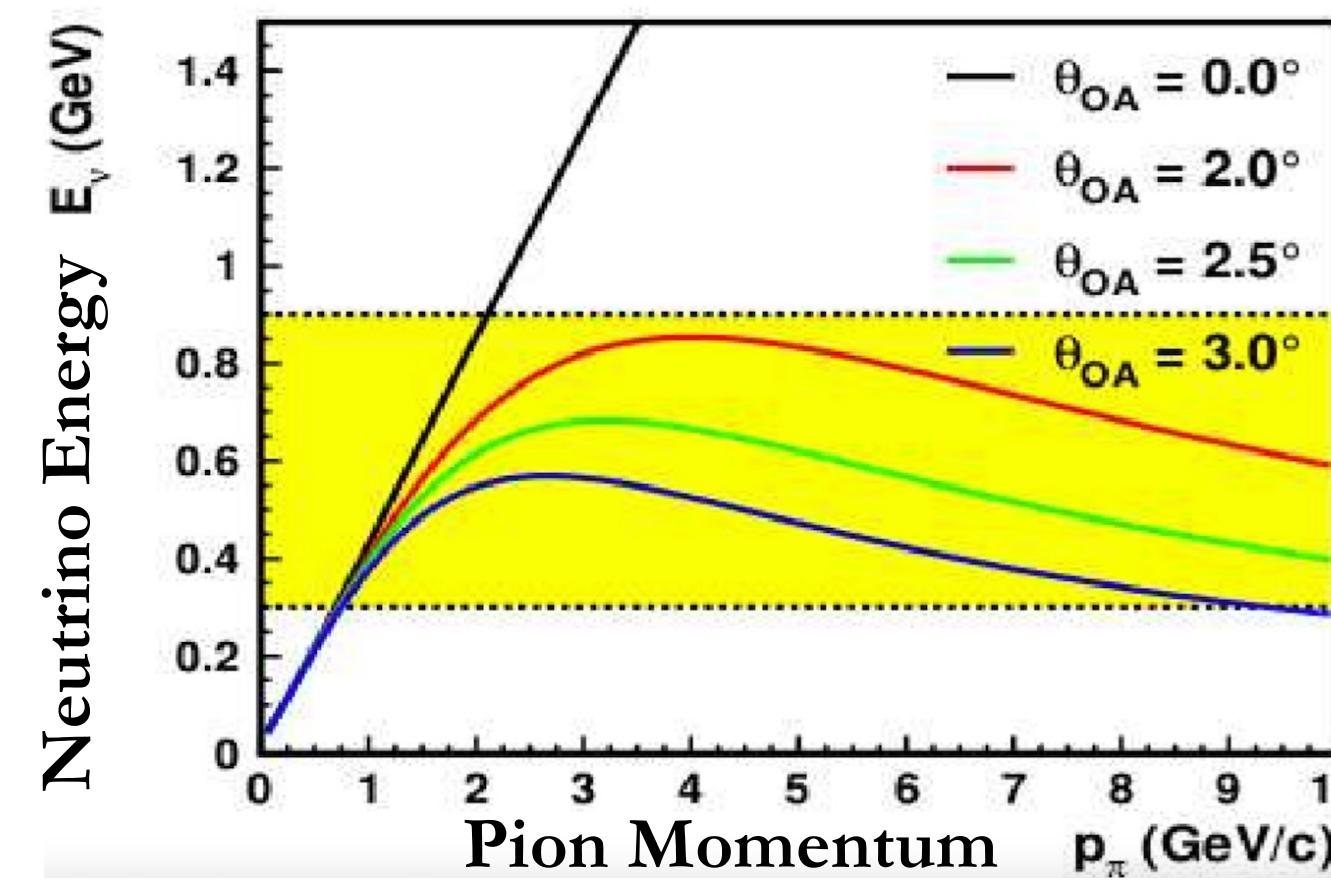
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

- Three flavour (anti)neutrino oscillation formula in vacuum:

$$P(\overset{\textcolor{purple}{(}}{\bar{\nu}_\alpha} \rightarrow \overset{\textcolor{purple}{(}}{\bar{\nu}_\beta}) = \delta_{\alpha\beta} - 4 \sum_{k>j} \operatorname{Re} \left\{ U_{\alpha k} U_{\beta k}^\star U_{\alpha j}^\star U_{\beta j} \right\} \sin^2 \frac{\Delta m_{kj}^2 L}{4E} \mp 2 \sum_{k>j} \operatorname{Im} \left\{ U_{\alpha k} U_{\beta k}^\star U_{\alpha j}^\star U_{\beta j} \right\} \sin 2 \frac{\Delta m_{kj}^2 L}{4E}$$

# The T2K Off-Axis Neutrino Flux

- SK and ND280 are placed at  $2.5^\circ$  off-axis angle with respect to the primary proton beam direction
- On-axis near detector (INGRID) used to monitor beam stability and direction
- Off-axis beam makes the  $\nu_\mu$  flux more narrow and peaked around the energies needed for observing the first oscillation maximum at SK (295 km baseline)
- High energy tail gets reduced



# $\nu_e$ ( $\bar{\nu}_e$ ) appearance probability at T2K

- Electron (anti)neutrino appearance probability:

$$P_{T2K}(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \frac{\Delta m_{32}^2 L}{4E}$$
$$+ \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \frac{\Delta m_{21}^2 L}{4E}$$
$$+ J \cos \left( \frac{\Delta m_{32}^2 L}{4E} \right) \pm \delta_{CP}$$

CP-conserving terms

Can be neglected at L/E of long baseline accelerator experiments

CP-violating term for  $\delta_{CP} \neq 0, \pi$

Plus sign for neutrinos

where  $J = \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \sin \frac{\Delta m_{32}^2 L}{4E} \sin \frac{\Delta m_{21}^2 L}{4E}$

# Appearance probability at T2K including matter effect

- Electron (anti)neutrino appearance probability around oscillation maximum:

$$P_{T2K}(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \frac{\Delta m_{32}^2 L}{4E} \left( 1 \pm \frac{2a}{\Delta m_{31}^2} (1 - 2 \sin^2 \theta_{13}) \right)$$

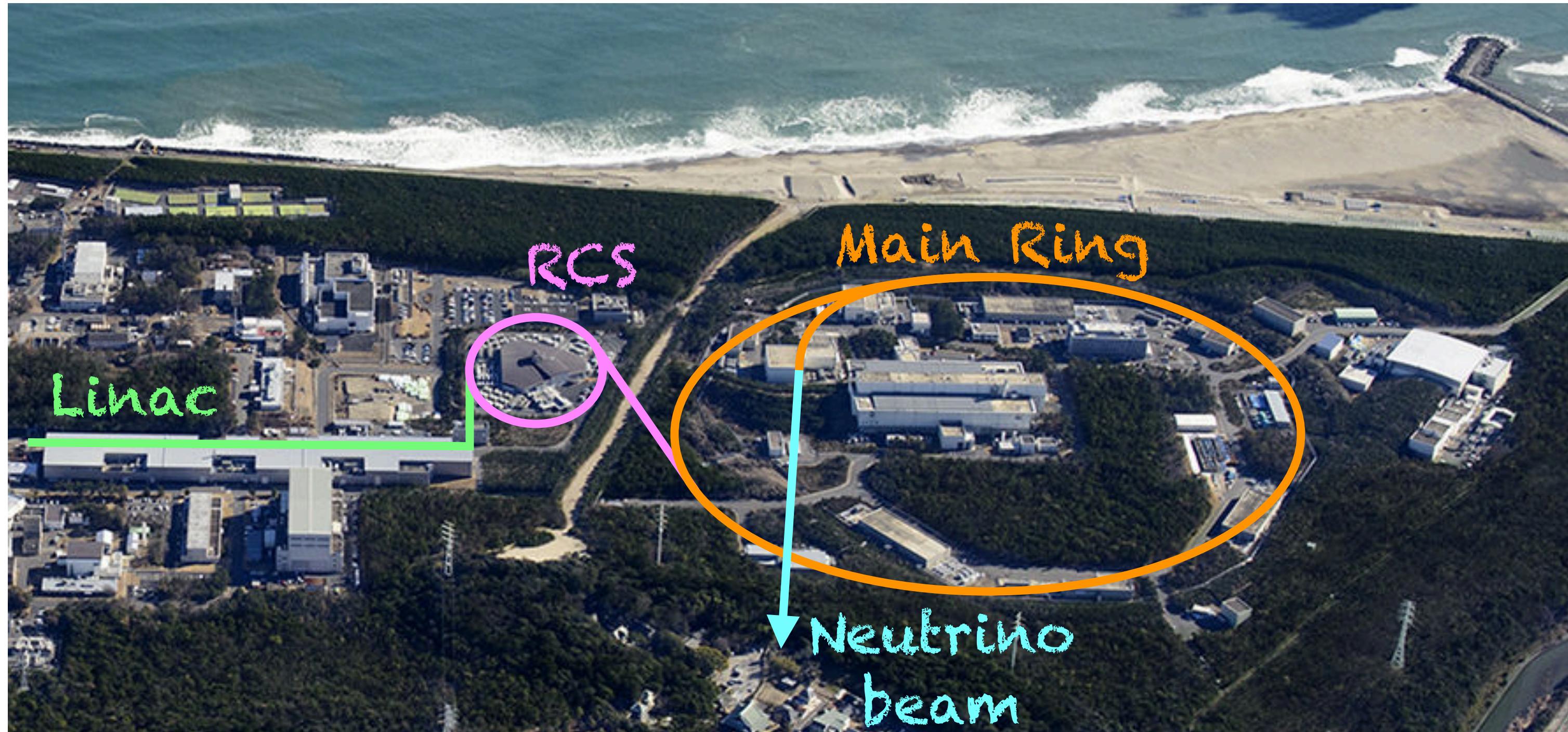
*Minus sign for anti-neutrinos*

$$\pm \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \sin \delta_{CP} \sin^2 \frac{\Delta m_{32}^2 L}{4E} \sin \frac{\Delta m_{21}^2 L}{4E}$$

*Plus sign for anti-neutrinos*

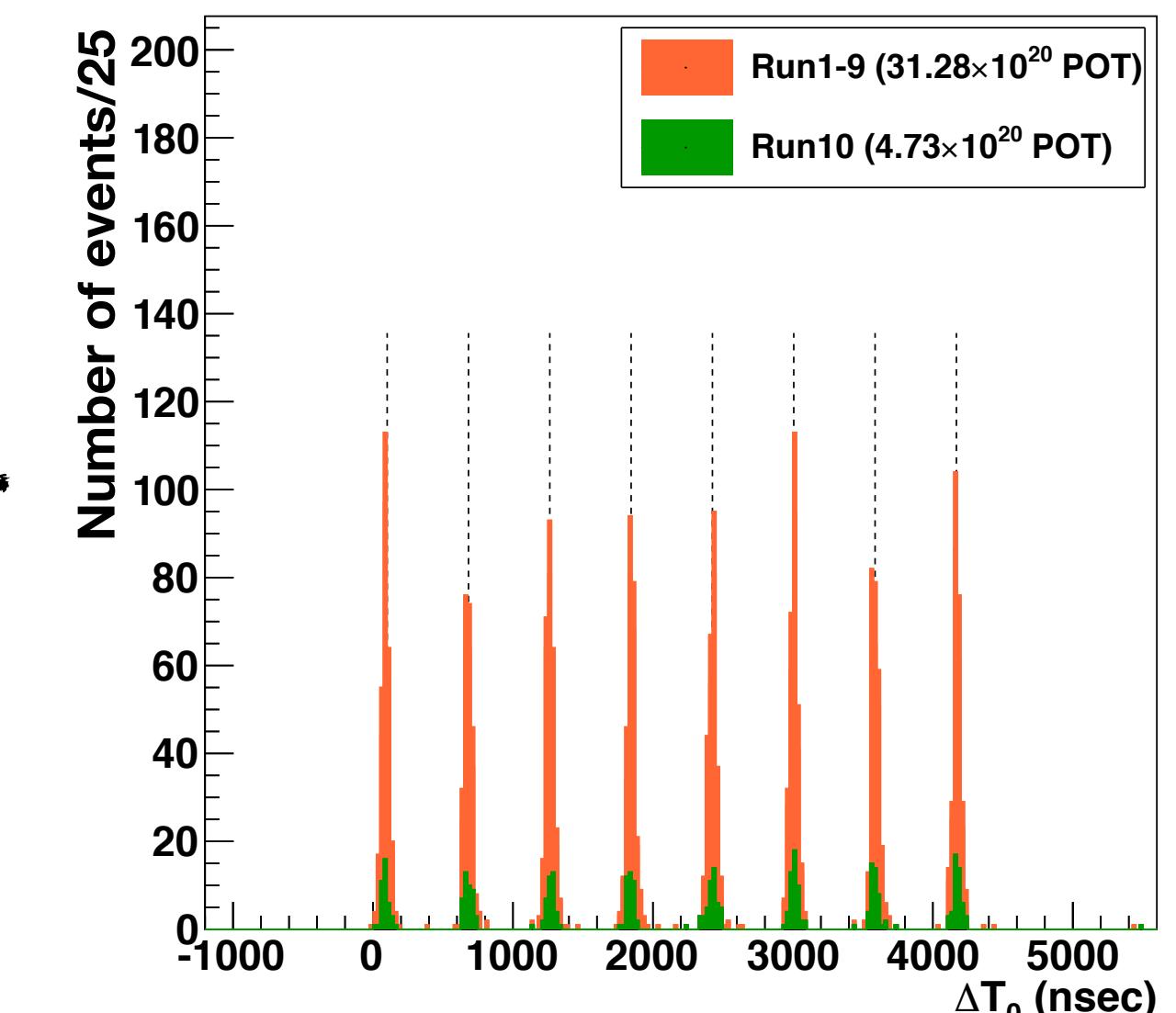
where  $a = 2\sqrt{2} G_F n_e E$  depends on the electron density in Earth's crust  $n_e$ , and Fermi's constant  $G_F$

# J-PARC facility

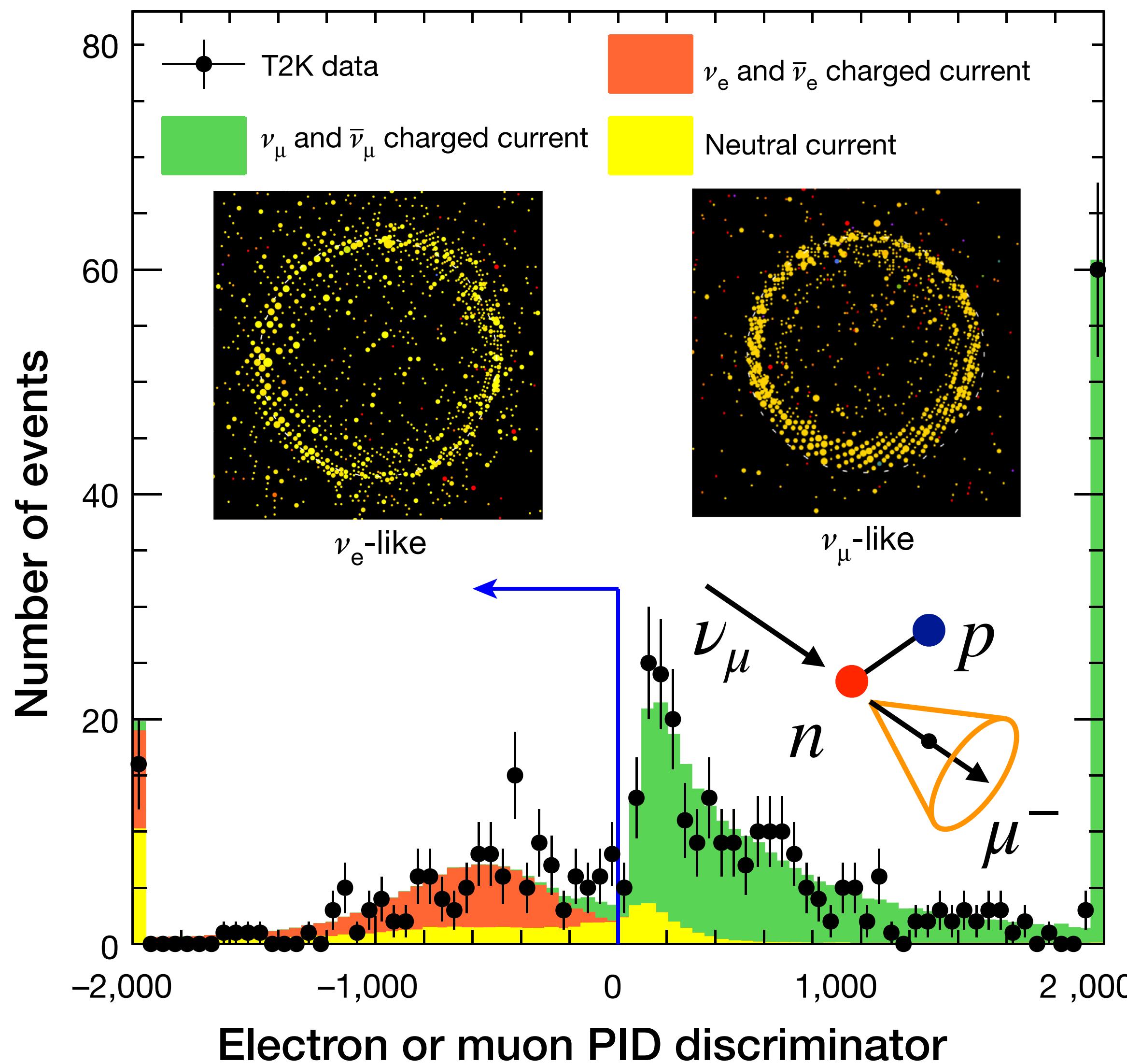


- 8 bunches per spill, with a spill repetition cycle of 2.48 s
- 80 ns bunch width, and 580 ns bunch separation

- 400 MeV Linac ( $H^-$ ) to 3 GeV Rapid Cycling Synchrotron ( $p$ ) to 30 GeV Main Ring ( $p$ ) to 600 MeV Neutrino beam ( $\nu$ )
- Low duty cycle for neutrino beam, so beam induced events are selected based on event timing



# Particle identification at Super-K



Electrons undergo more scattering, hence produce “fuzzy” Cherenkov rings

Good PID achieved using times, charge and position of all hit PMTs

$$L(\Gamma, \theta) = \prod_j^{unhit \ PMTs} P_j(unhit | \Gamma, \theta) \prod_i^{hit \ PMTs} \{1 - P_i(unhit | \Gamma, \theta)\}$$

Event topology hypothesis (1R muon-like, 1R electron-like etc.)

Event kinematics

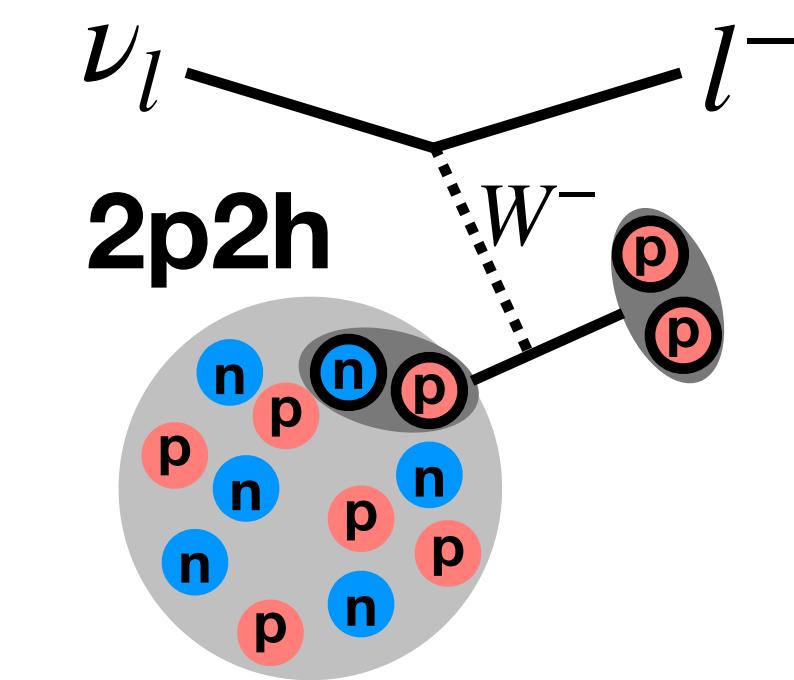
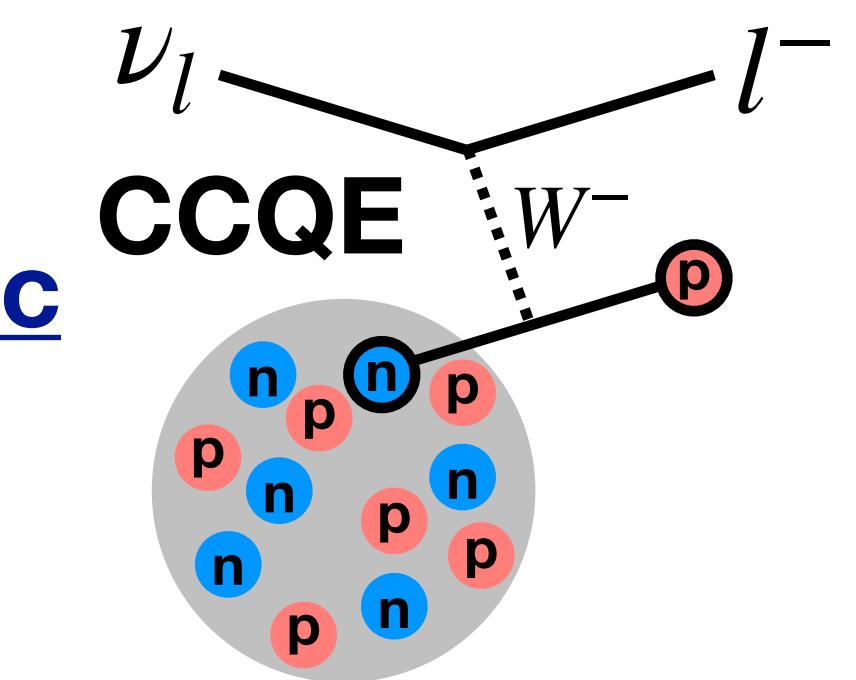
Charge signal likelihood

Timing likelihood

$$PID = \ln \left( \frac{L_{\mu}^{max}}{L_e^{max}} \right)$$

# Cross section model updates

- Updates for our latest oscillation measurement (with the new [NEUT 5.4.0](#) interaction model):
  1. CCQE initial state nuclear model migrated from [Relativistic Fermi Gas plus Random Phase Approximation](#) to [Spectral Function](#) (strongly favoured by electron scattering data)
  2. Better (shell-dependent) treatment of removal energy comes with the SF CCQE model
  3. More sophisticated treatments of the energy dependent 2p2h uncertainty, deep inelastic scattering uncertainty and final state interaction correlations between near and far detectors



# Systematic uncertainties before/after ND fit

Before Near detector fit:

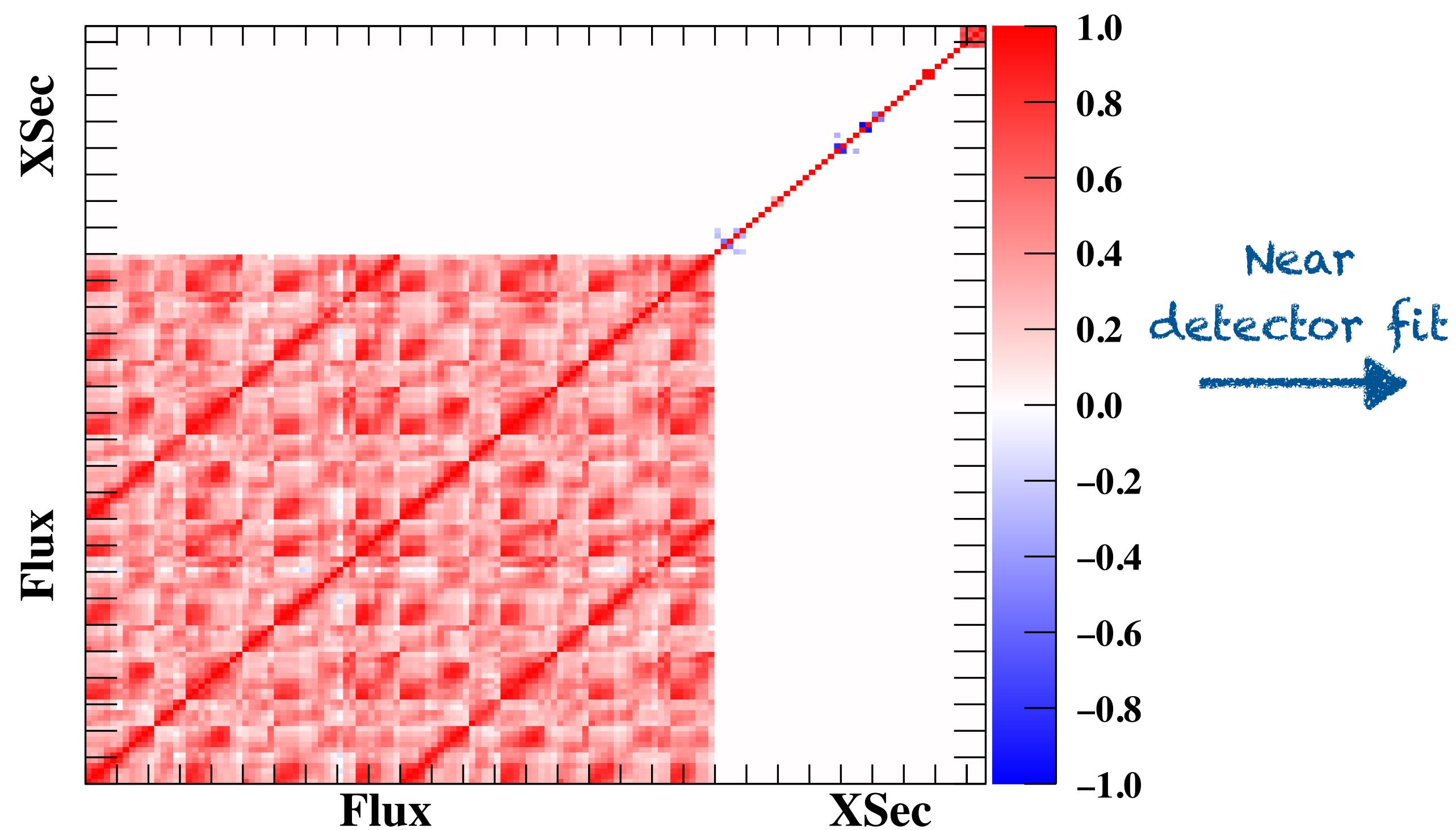
Error source (units: %)	1R $\mu$		1Re				FHC/RHC
	FHC	RHC	FHC	RHC	FHC CC1 $\pi^+$		
Flux	5.1	4.7	4.8	4.7	4.9		2.7
Cross-section (all)	10.1	10.1	11.9	10.3	12.0		10.4
SK+SI+PN	2.9	2.5	3.3	4.4	13.4		1.4
<b>Total</b>	11.1	11.3	13.0	12.1	18.7		10.7

After Near detector fit:

Error source (units: %)	1R $\mu$		1Re				FHC/RHC
	FHC	RHC	FHC	RHC	FHC CC1 $\pi^+$		
Flux	2.9	2.8	2.8	2.9	2.8		1.4
Xsec (ND constr)	3.1	3.0	3.2	3.1	4.2		1.5
Flux+Xsec (ND constr)	2.1	2.3	2.0	2.3	4.1		1.7
Xsec (ND unconstrained)	0.6	2.5	3.0	3.6	2.8		3.8
SK+SI+PN	2.1	1.9	3.1	3.9	13.4		1.2
<b>Total</b>	3.0	4.0	4.7	5.9	14.3		4.3

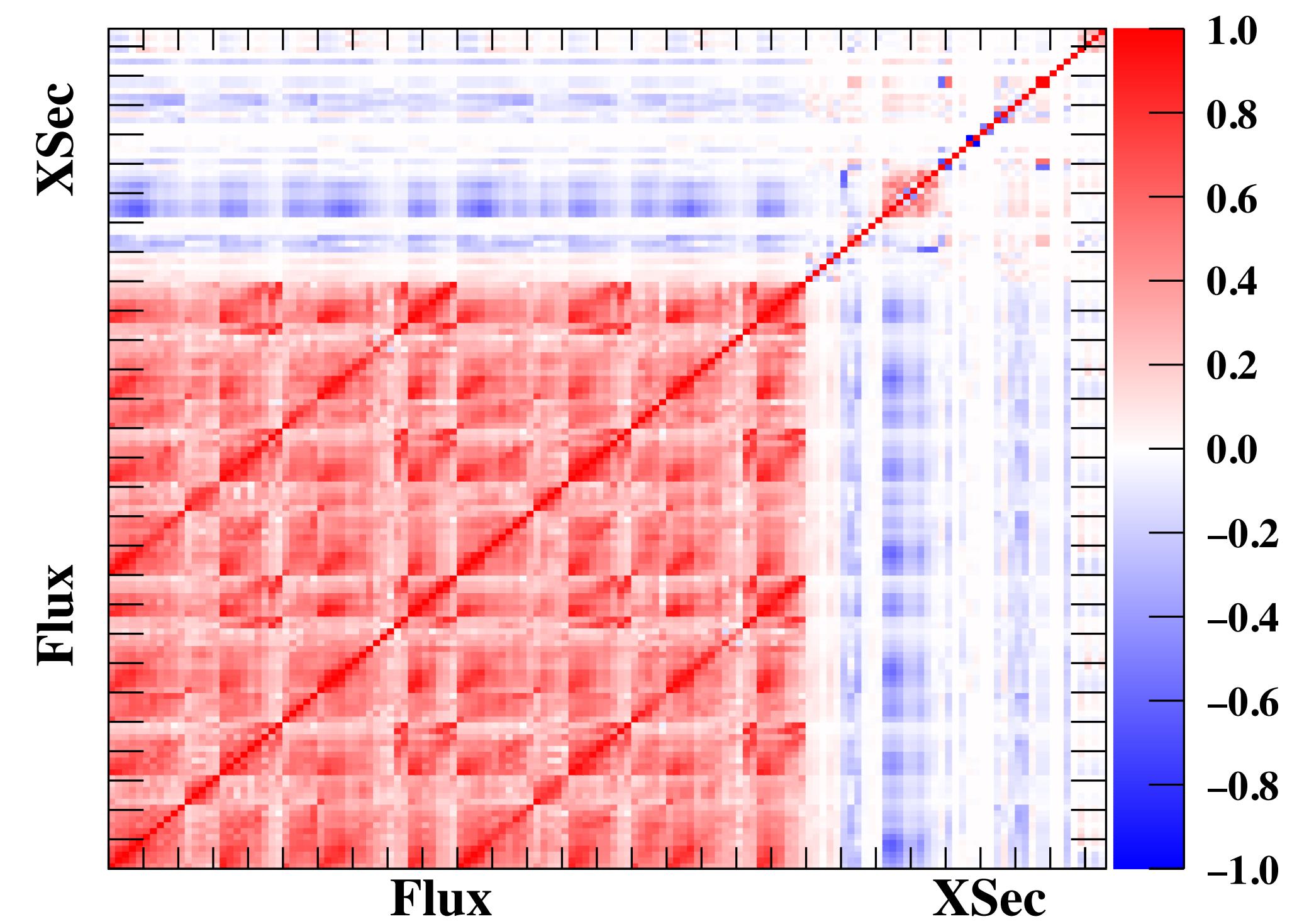
# More on the near detector fit

**Flux and Xsec Prefit Correlation Matrix**



Many more detector systematics enter the fit, which are not shown here for clarity

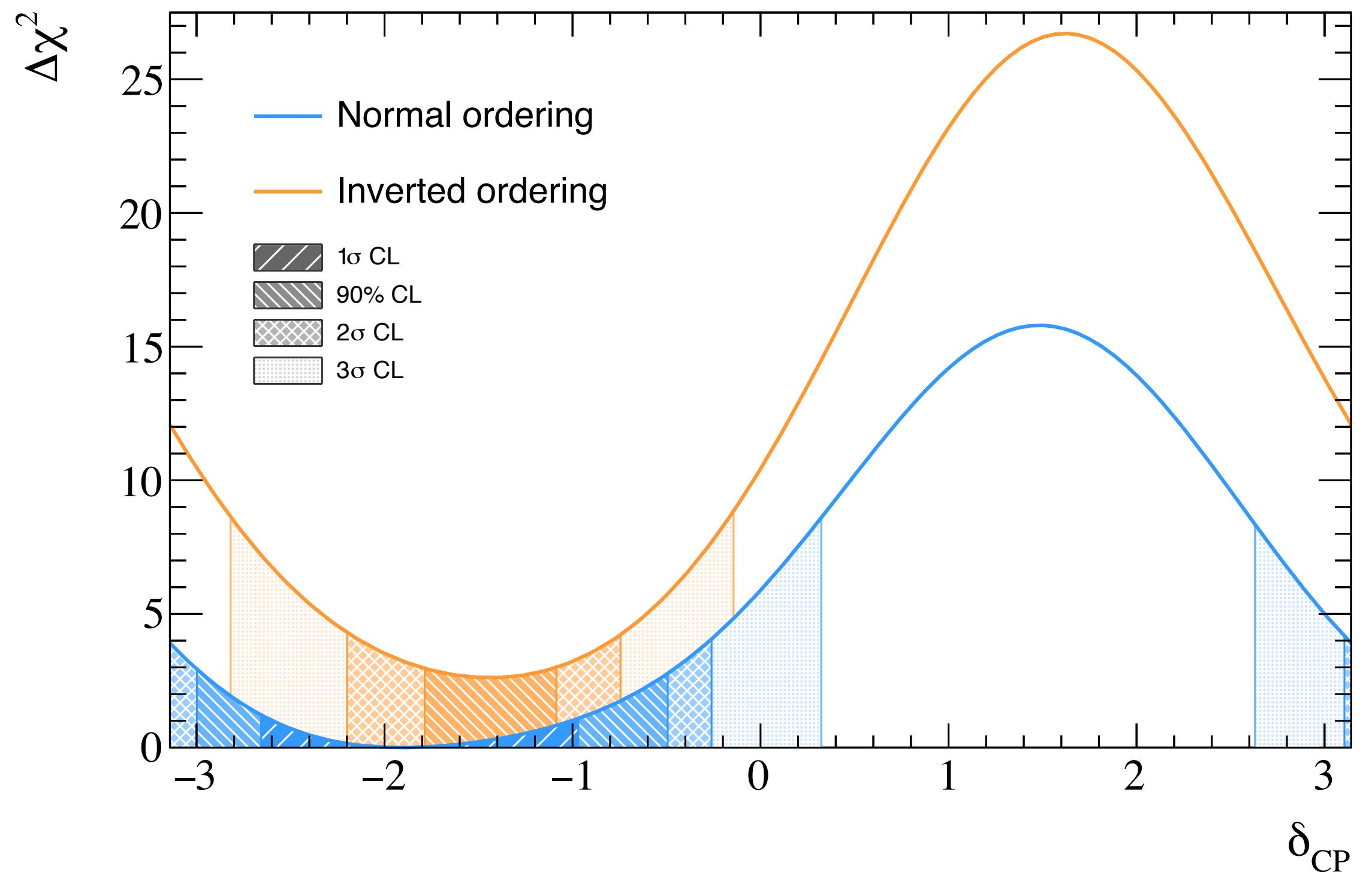
**Flux and Xsec Postfit Correlation Matrix**



Anti-correlations between flux and cross section parameters after ND-fit

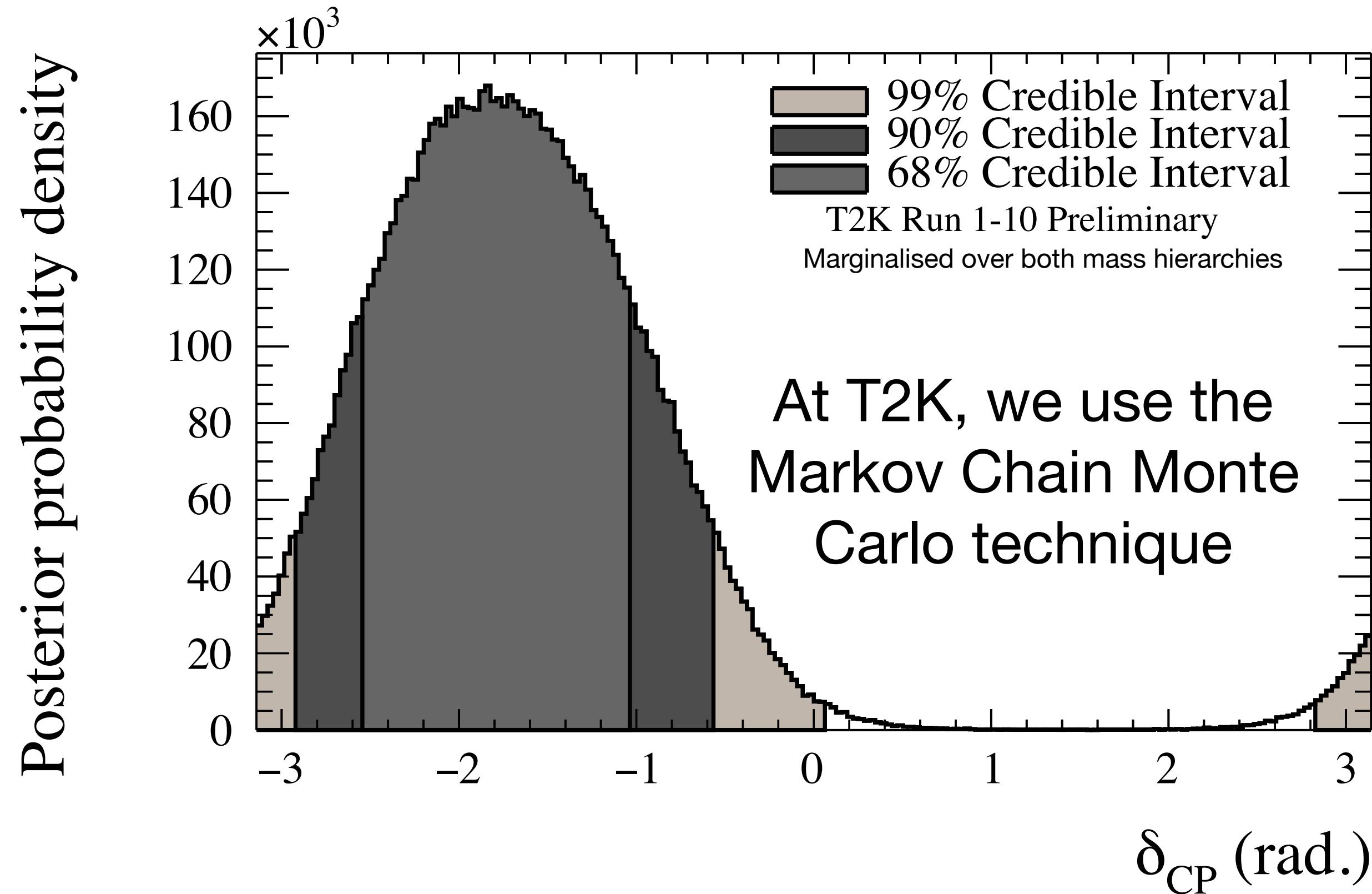
# Frequentist vs. Bayesian: Conceptual differences

Frequentist results produce  $\Delta\chi^2$  distributions



Areas with a lower  $\Delta\chi^2$  are more likely to contain the “true” value

Bayesian results produce posterior probability density distributions



Areas with a larger probability density are more likely to contain the “true” value

# Markov Chain Monte Carlo

- MCMC algorithm presents an efficient solution for sampling from a complex multivariate posterior likelihood distribution
- The algorithm provides a recipe for building the Markov chain
- At each step, an ensemble of model parameters is either accepted or rejected based on a comparison with a metric
- Parameters in the resulting Markov Chain have a distribution density proportional to the posterior likelihood

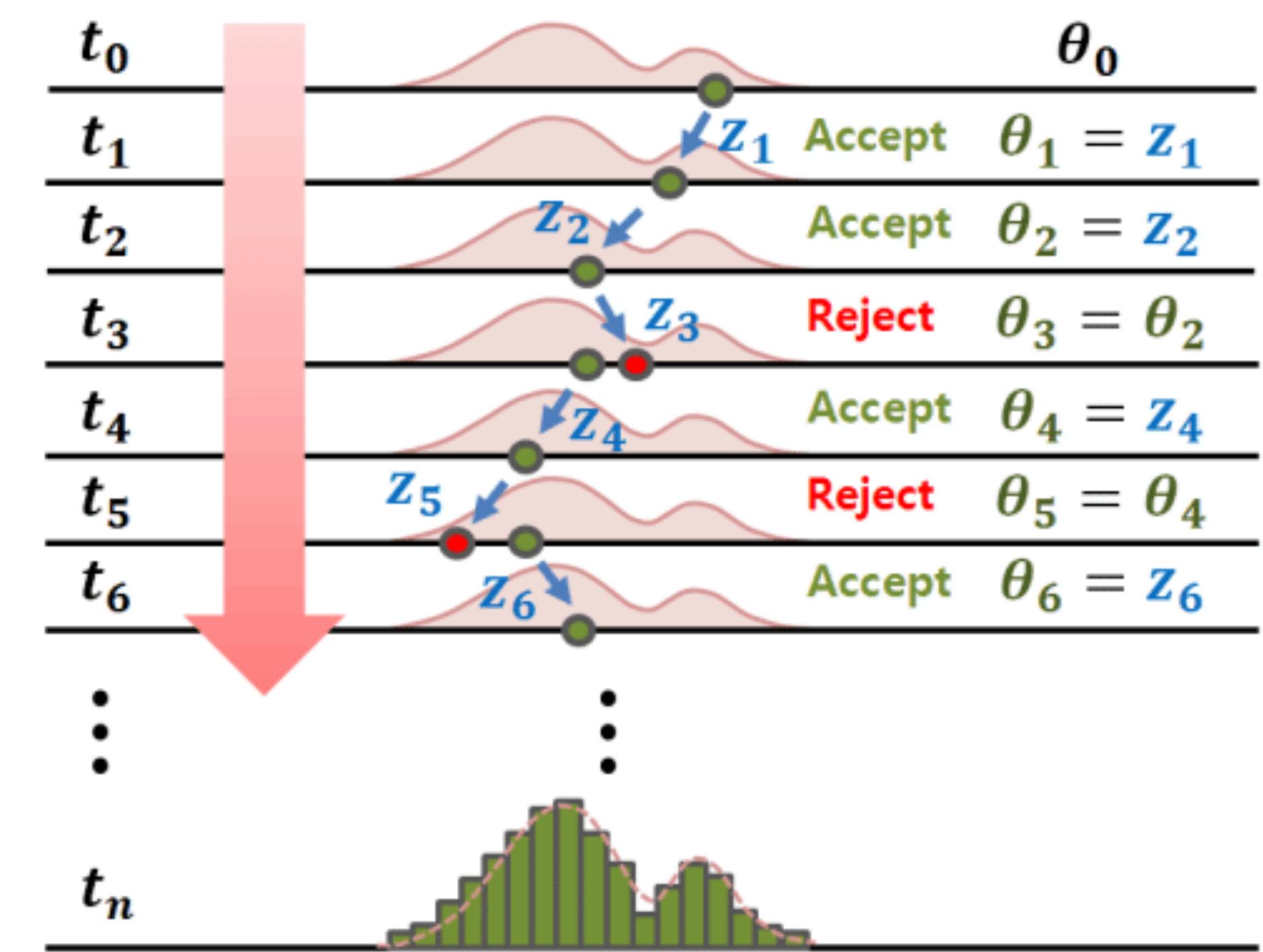
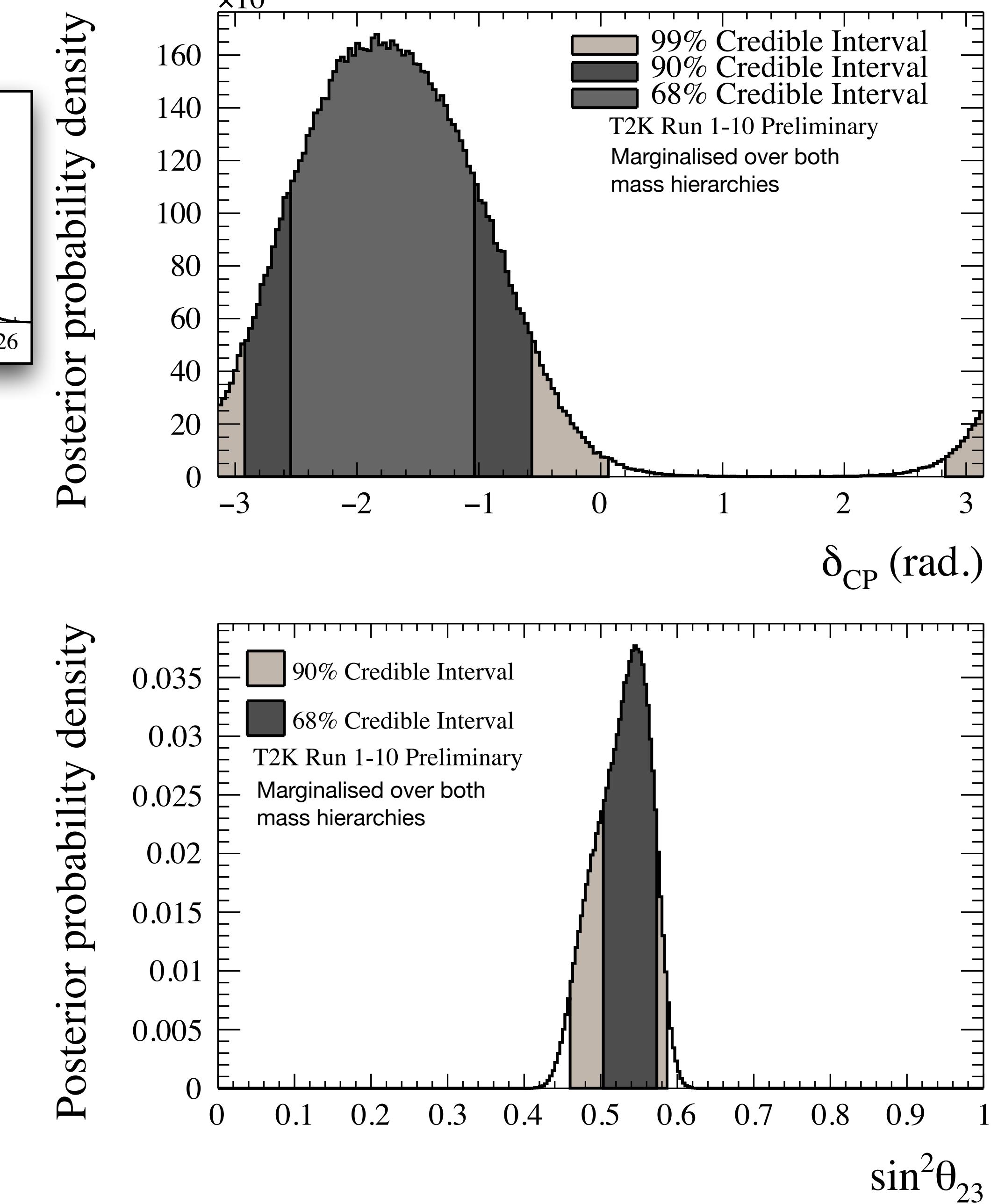
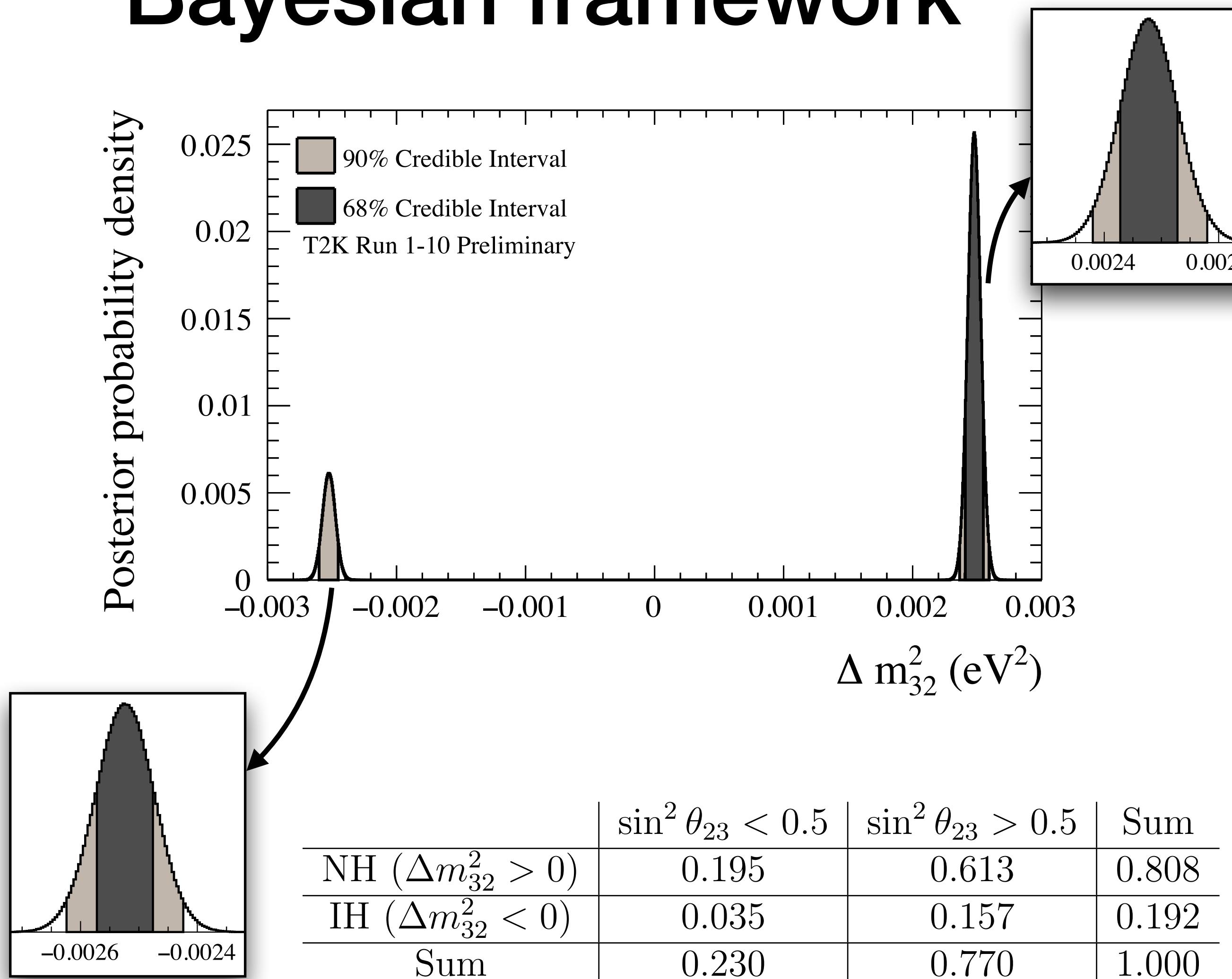


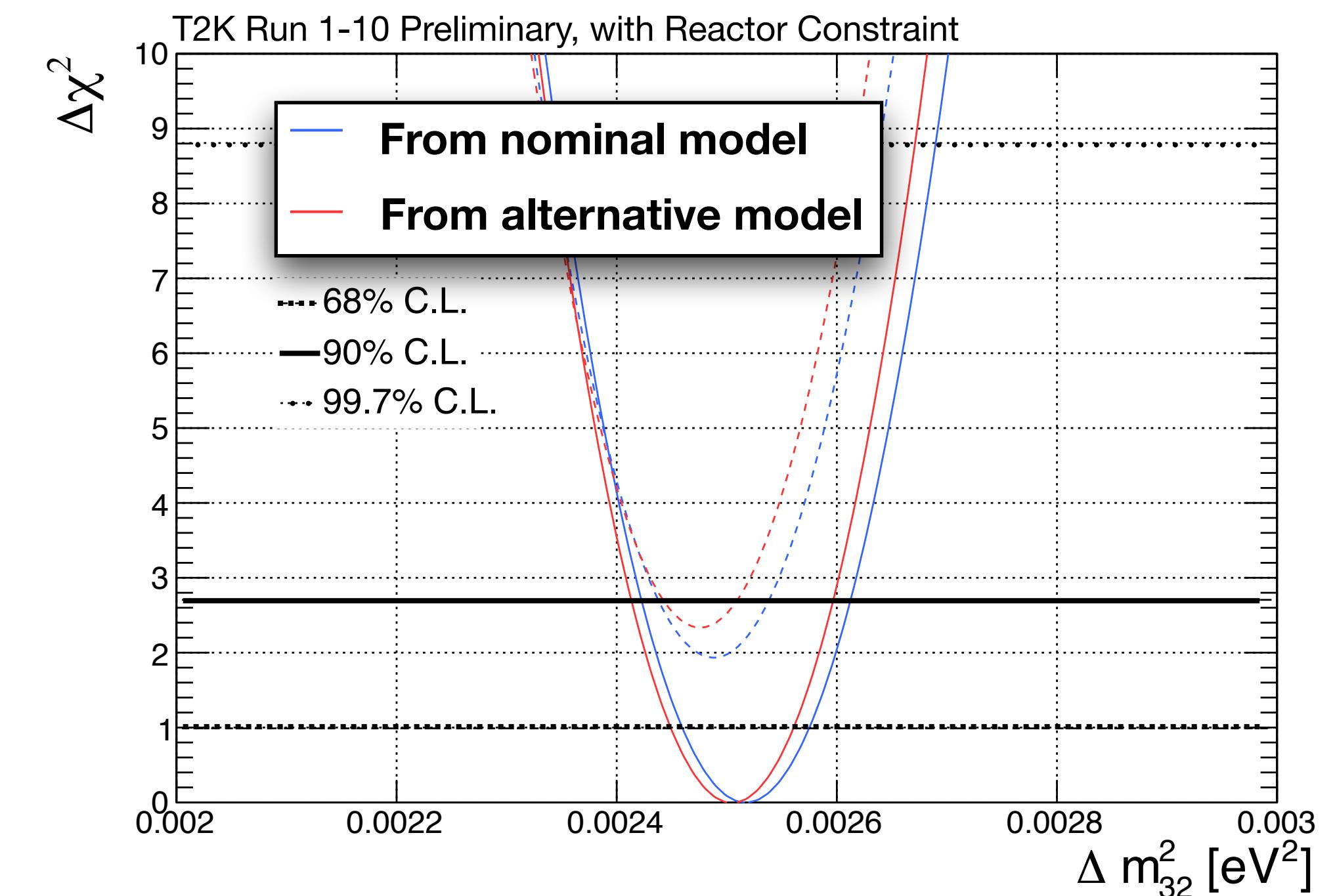
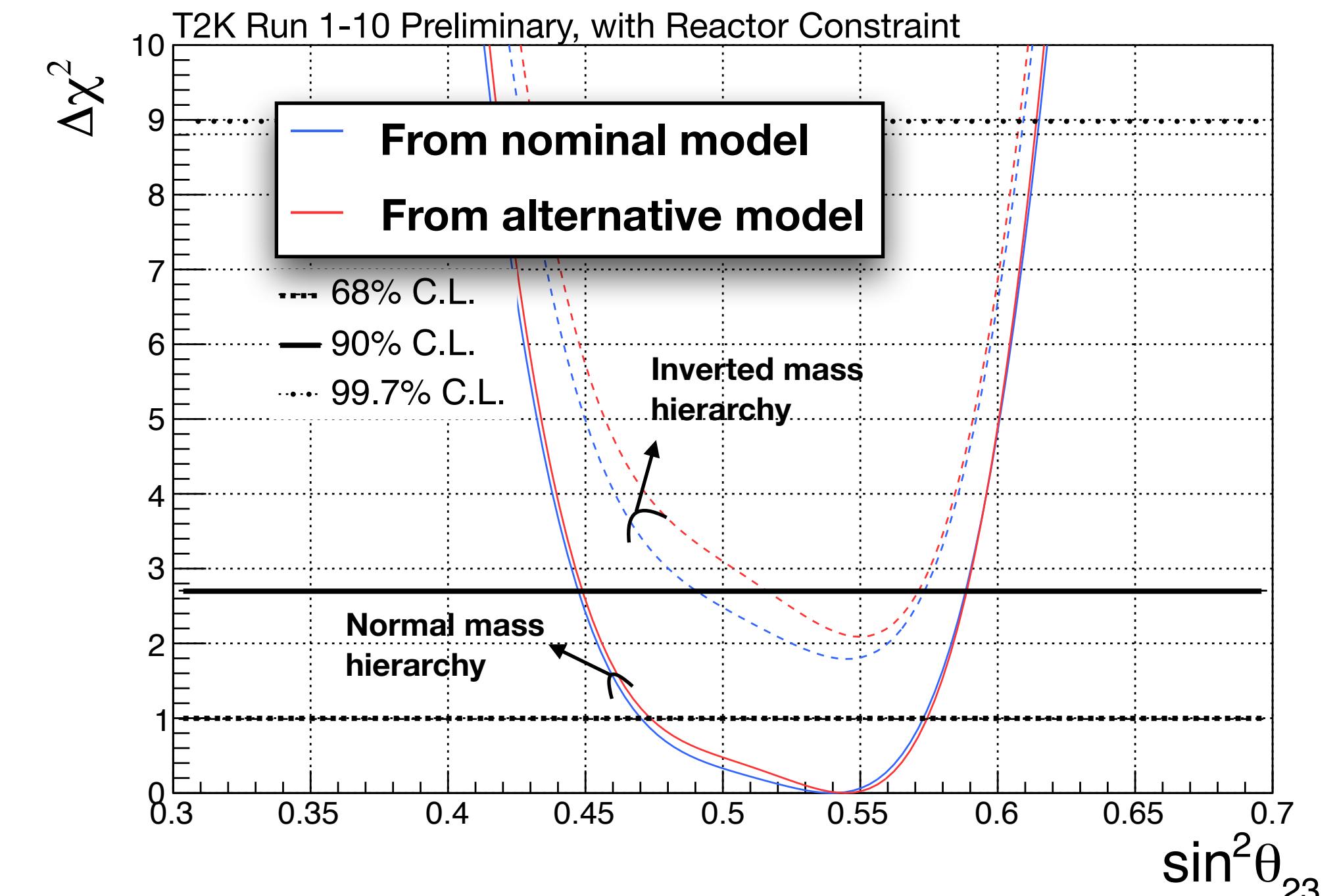
Figure from [Adaptive Markov chain Monte Carlo algorithms for Bayesian inference: recent advances and comparative study](#)

# Full inference in Bayesian framework



# Bias studies for oscillation analyses

- The sturdiness of our analyses is tested by performing **fits on simulated data from alternative theory and data driven models**
- No significant biases have been observed for  $\theta_{13}$  and  $\theta_{23}$
- **Slight bias observed for  $\Delta m_{32}^2$** , which has been accounted for by adding a **compensating uncertainty of  $1.4 \times 10^{-5} \text{ eV}^2$**
- The slight bias on  $\delta_{CP}$  is accounted for by reporting the maximum observed shift on the  $\delta_{CP}$  contour boundary (here the contour is completely statistics dominated)



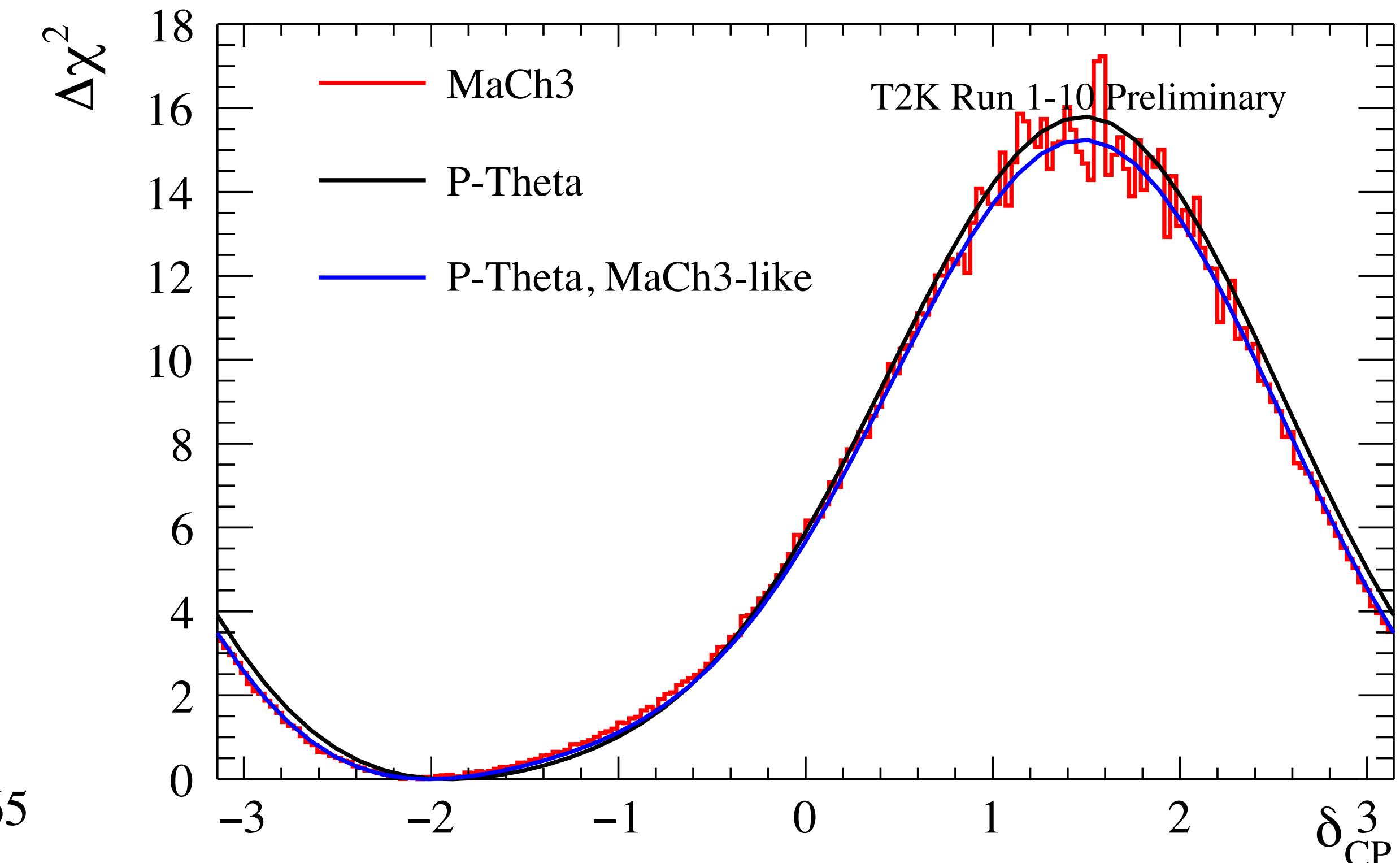
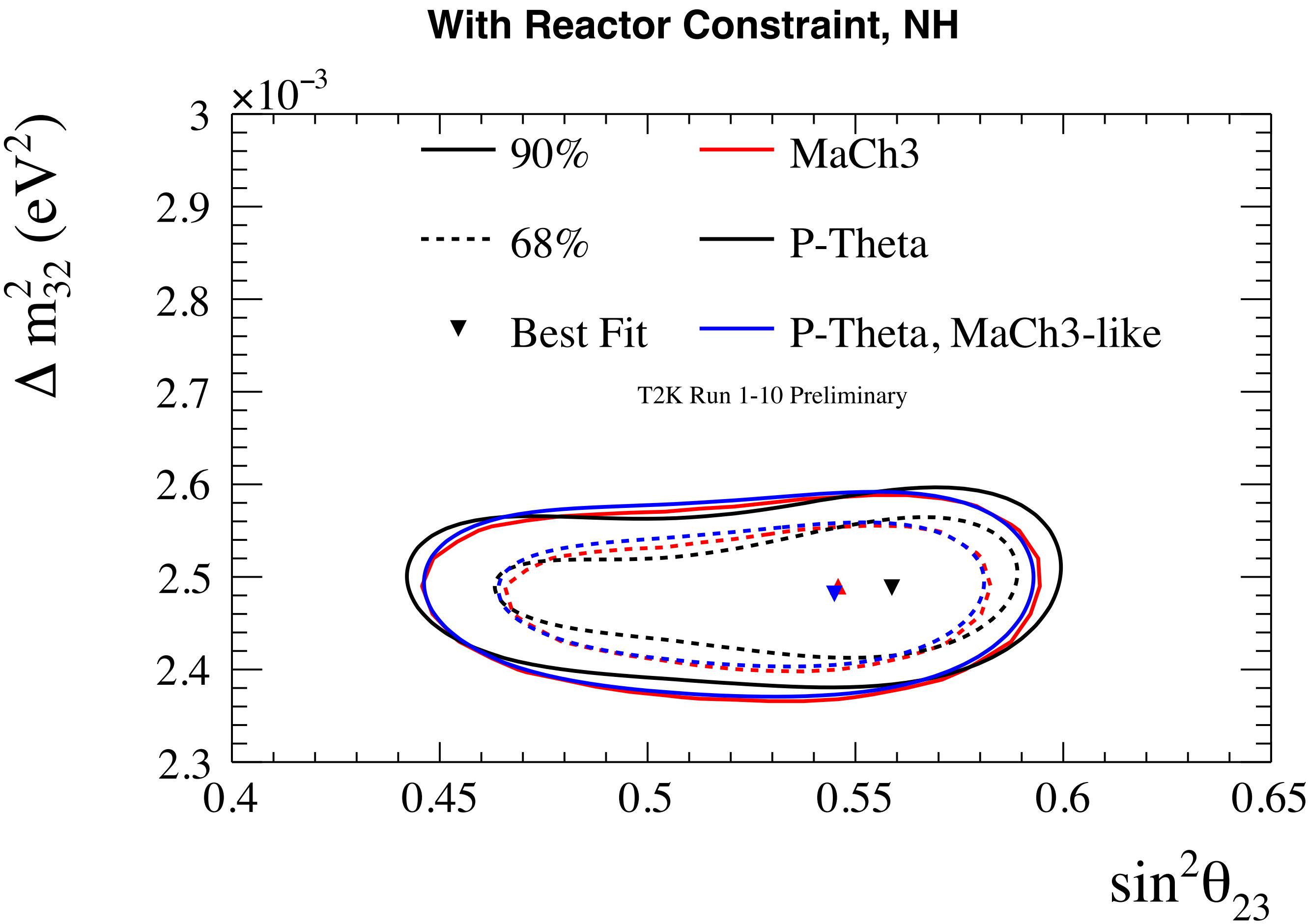
# Bayesian vs. Frequentist T2K Analyses Strategies

	<b>MaCh3</b>	<b>P-Theta</b>	<b>“MaCh3-like” P-Theta</b>
<b>Kinematic variables for CC0<math>\pi</math> 1R<sub>e</sub></b>	$(E_e^{rec}, \theta_e^{rec})$	$(p_e^{rec}, \theta_e^{rec})$	$(E_e^{rec}, \theta_e^{rec})$
<b>Likelihood</b>	Binned Poisson likelihood ratio	Binned Poisson likelihood ratio	Binned Poisson likelihood ratio
<b>Likelihood optimization</b>	Markov chain Monte Carlo	Grid search and gradient descent	Grid search and gradient descent
<b>Resulting contours/ limits</b>	Bayesian credible intervals	Frequentist confidence intervals with Feldman Cousins	Frequentist confidence intervals with Feldman Cousins
<b>Mass hierarchy inference</b>	Bayes factor from fraction of points in each hierarchy	Bayes factor from likelihood integration	Frequentists p-value from generated PDF
<b>Inputs from near detector</b>	Simultaneous near+far detector joint fit	Constrained covariance matrix	Constrained covariance matrix
<b>Systematics handling</b>	Marginalisation after fit	Marginalisation during fit	Marginalisation during fit

*T2K always employs marginalisation rather than profiling because of inherently better non-Gaussian systematics treatment*

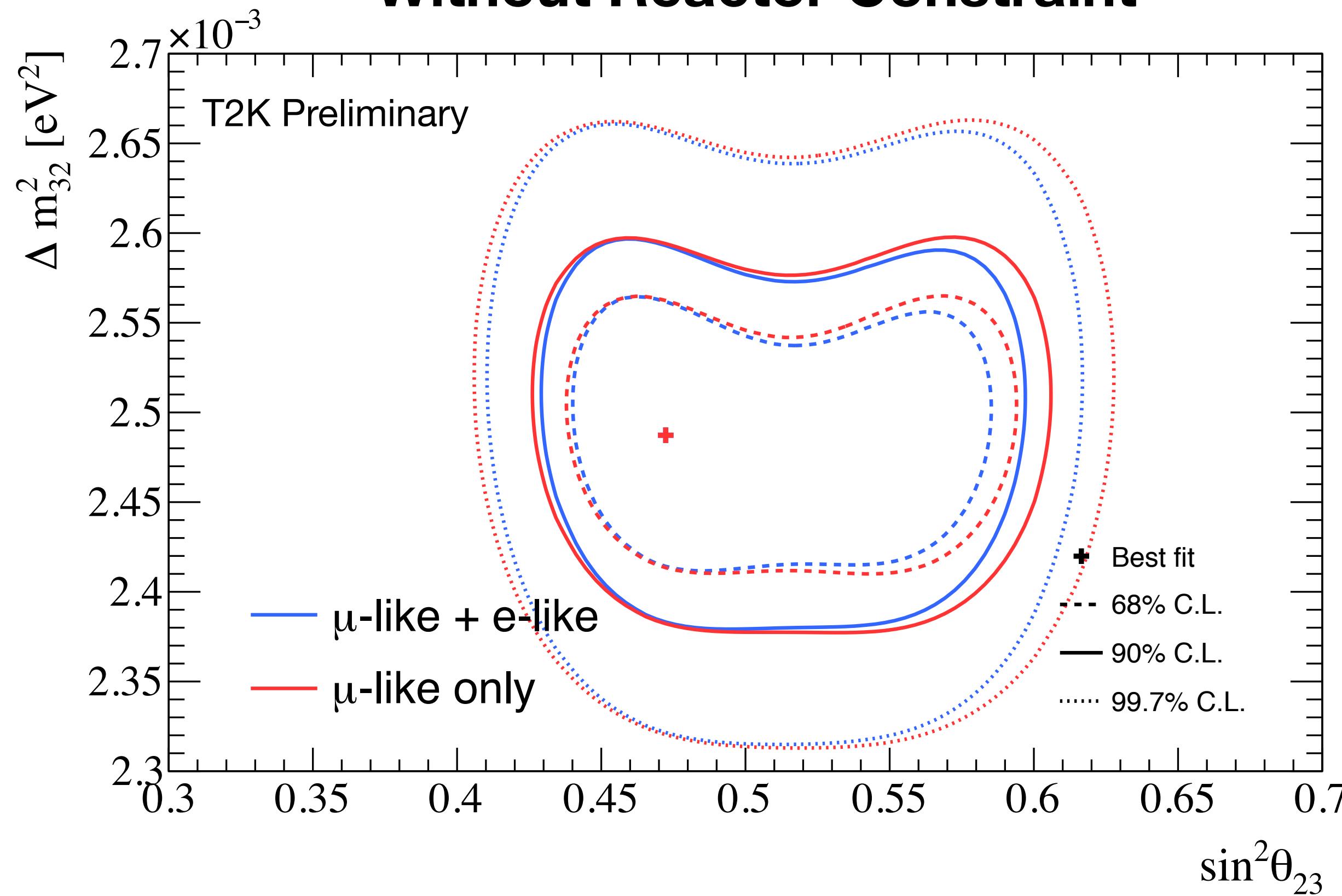
*Feldman  
Cousins method  
gives critical  
 $\Delta\chi^2$  values for  
the confidence  
intervals*

# Cross-checks between T2K oscillation analyses

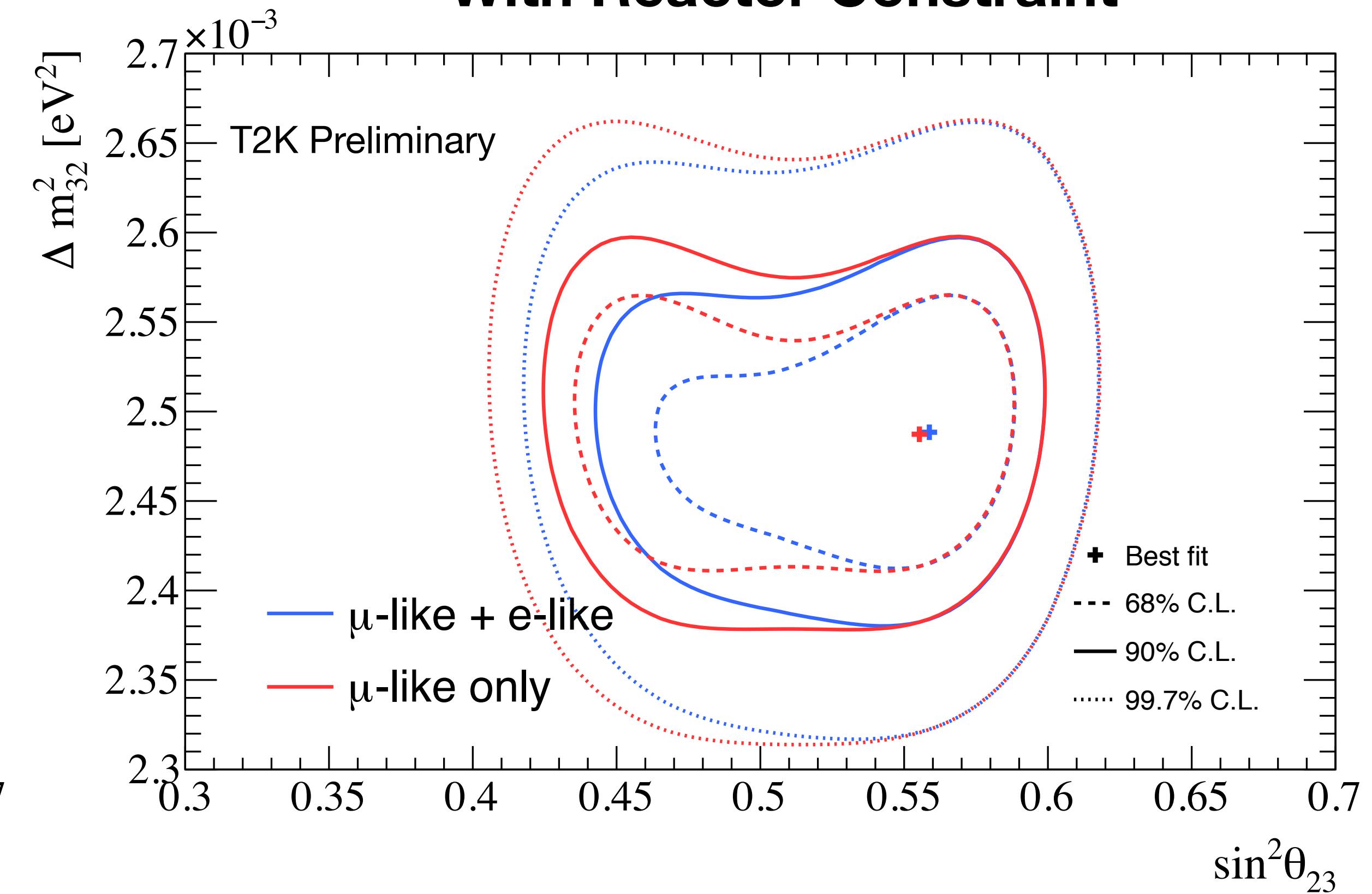


# $\theta_{23}$ octant preference

**Without Reactor Constraint**



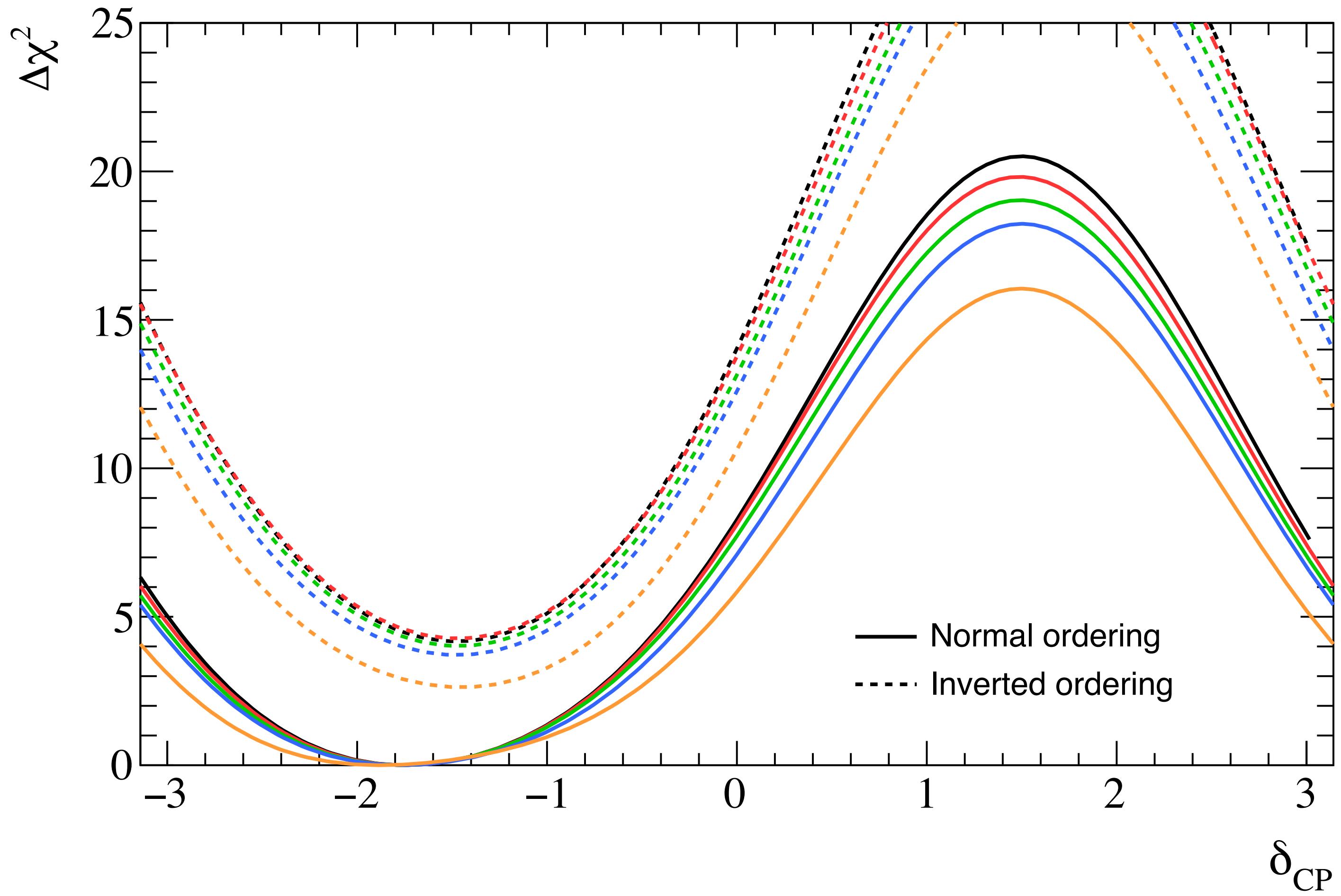
**With Reactor Constraint**



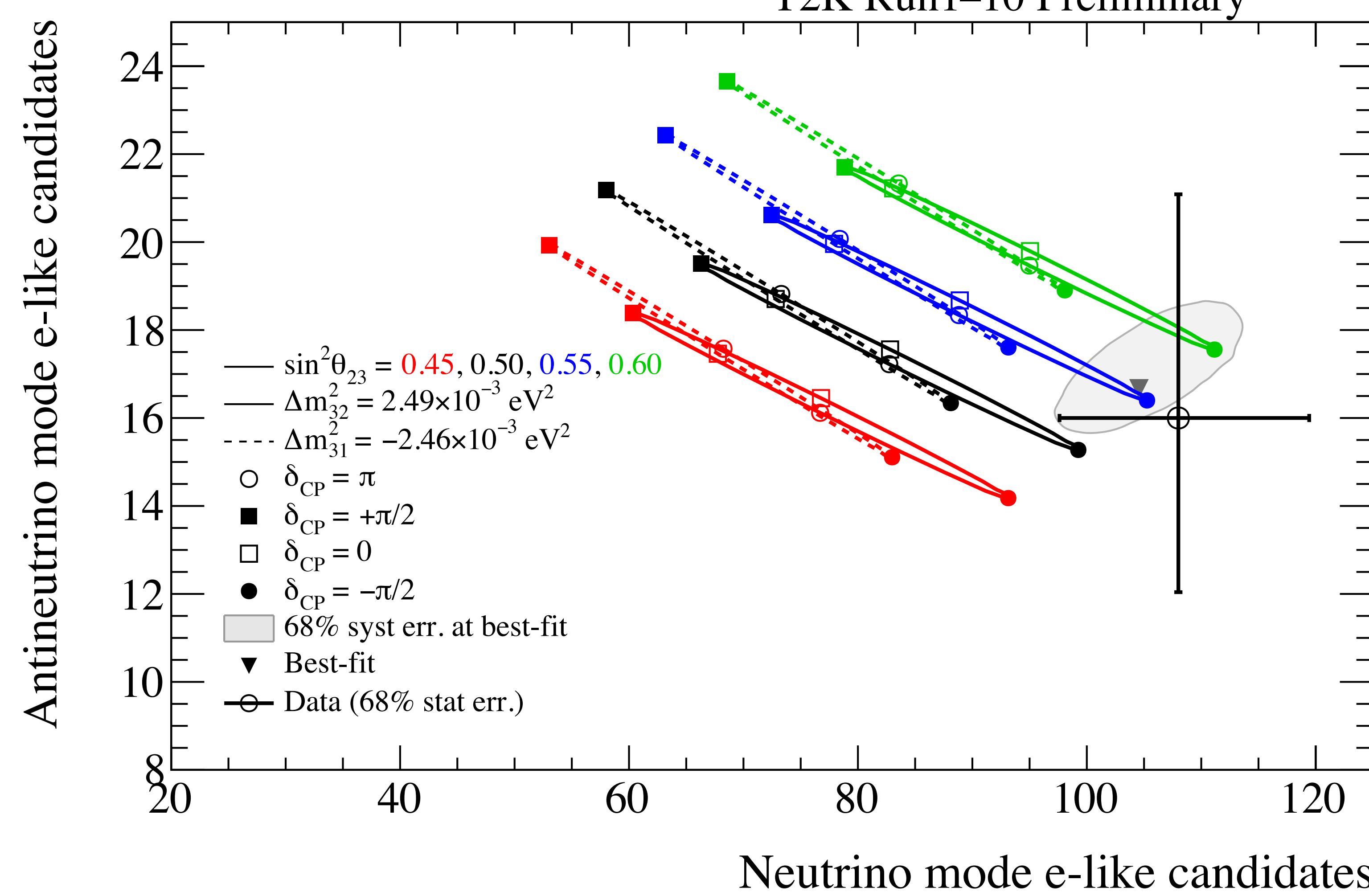
Upper  $\theta_{23}$  octant preference is driven by the electron-like Cherenkov ring sample, in combination with the tighter reactor constraint on  $\theta_{13}$

# Evolution of T2K's $\delta_{CP}$ result

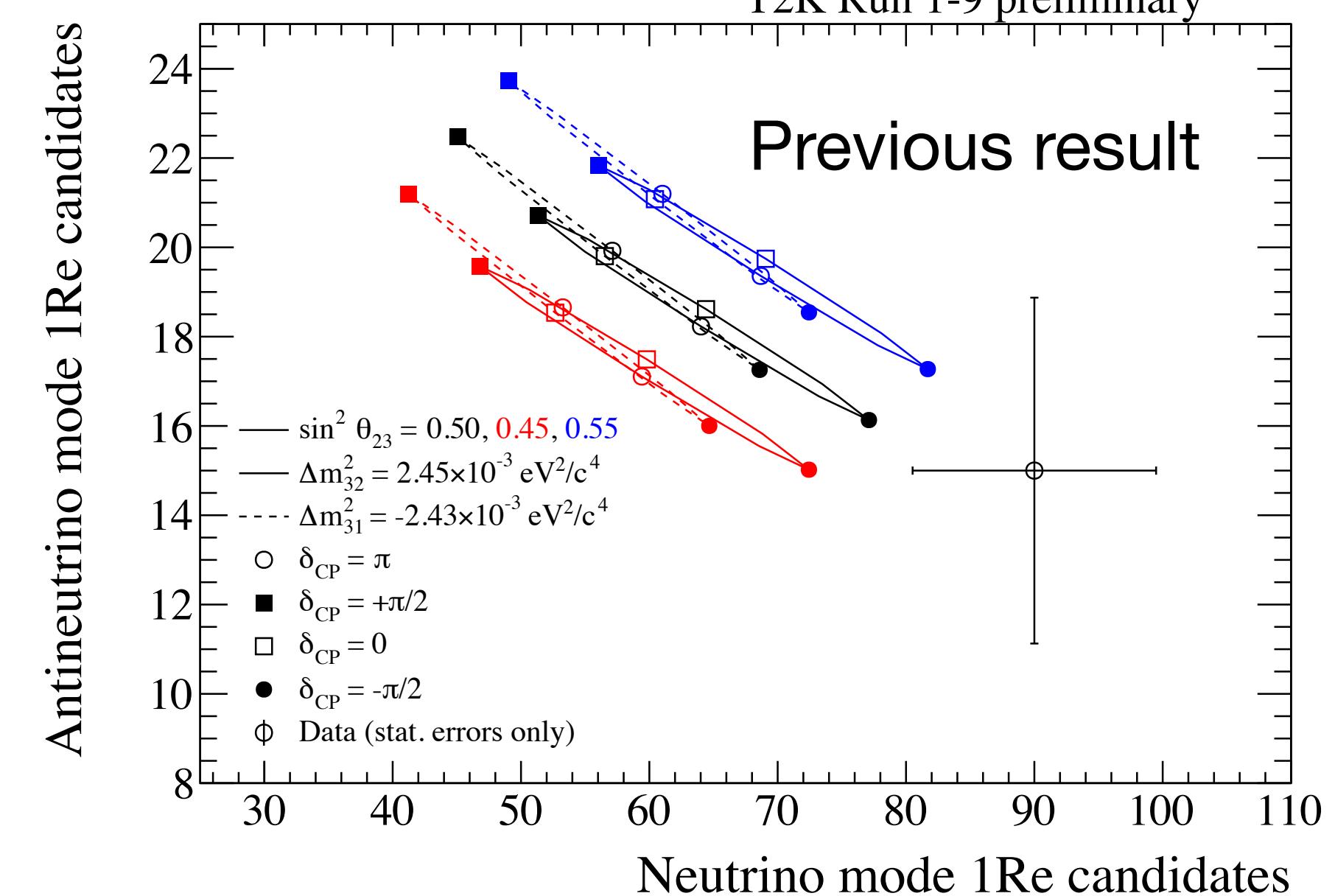
- Can sequentially add improvements to the OA2018 analysis (Nature publication), and examine the measurement sensitivity:
  - A: Nature result
  - B: Improved cross section and flux models mentioned earlier
  - C: Using the updated reactor constraint on  $\theta_{13}$  (from 2019 Particle Data Group)
  - D: New Super-K calibration, causing some events to drop in and out from samples
  - E: Addition of the new Run 10 T2K data
- **Greatest change in our measurement (slightly weaker constraint) is driven by the addition of new data**



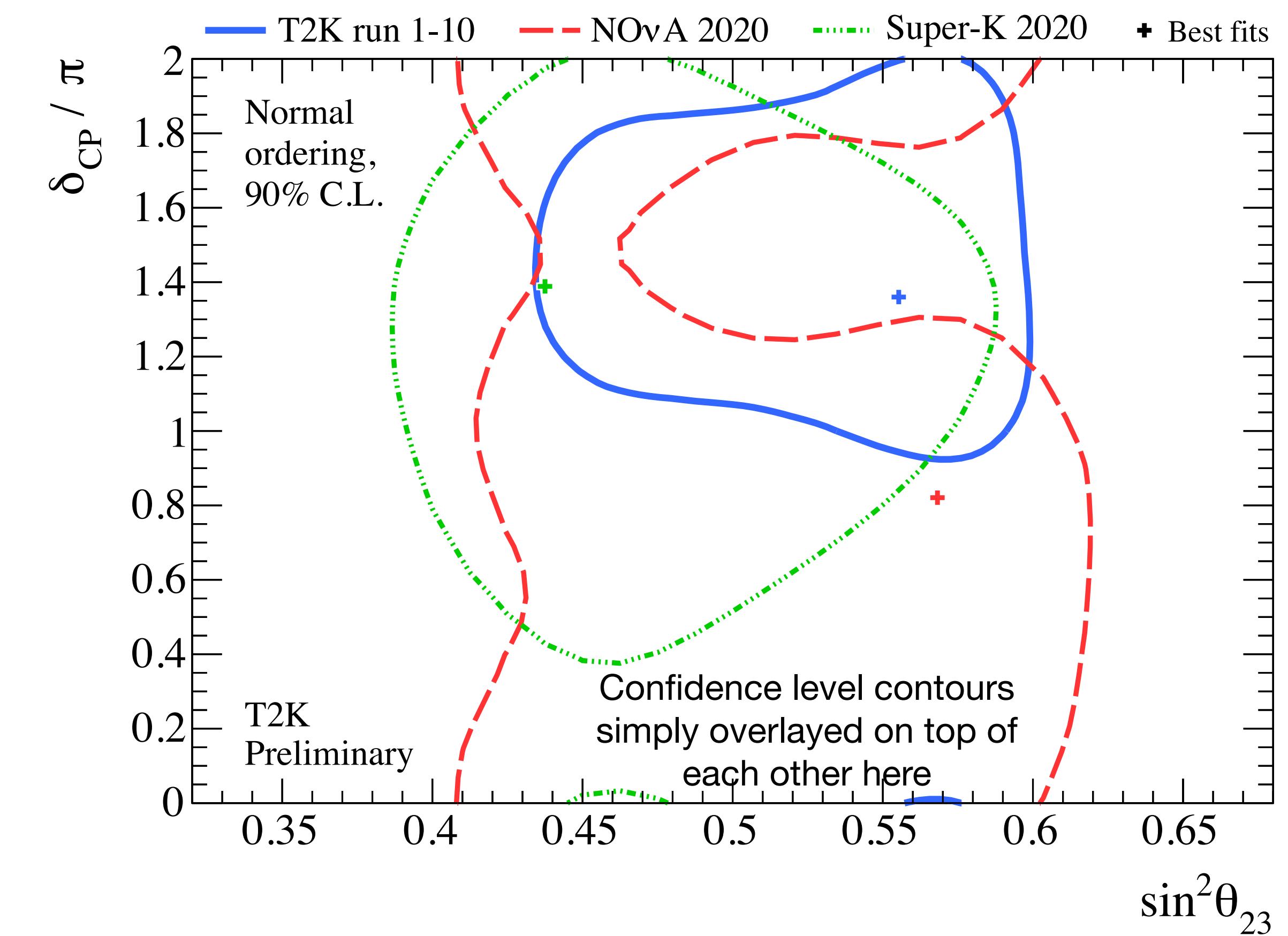
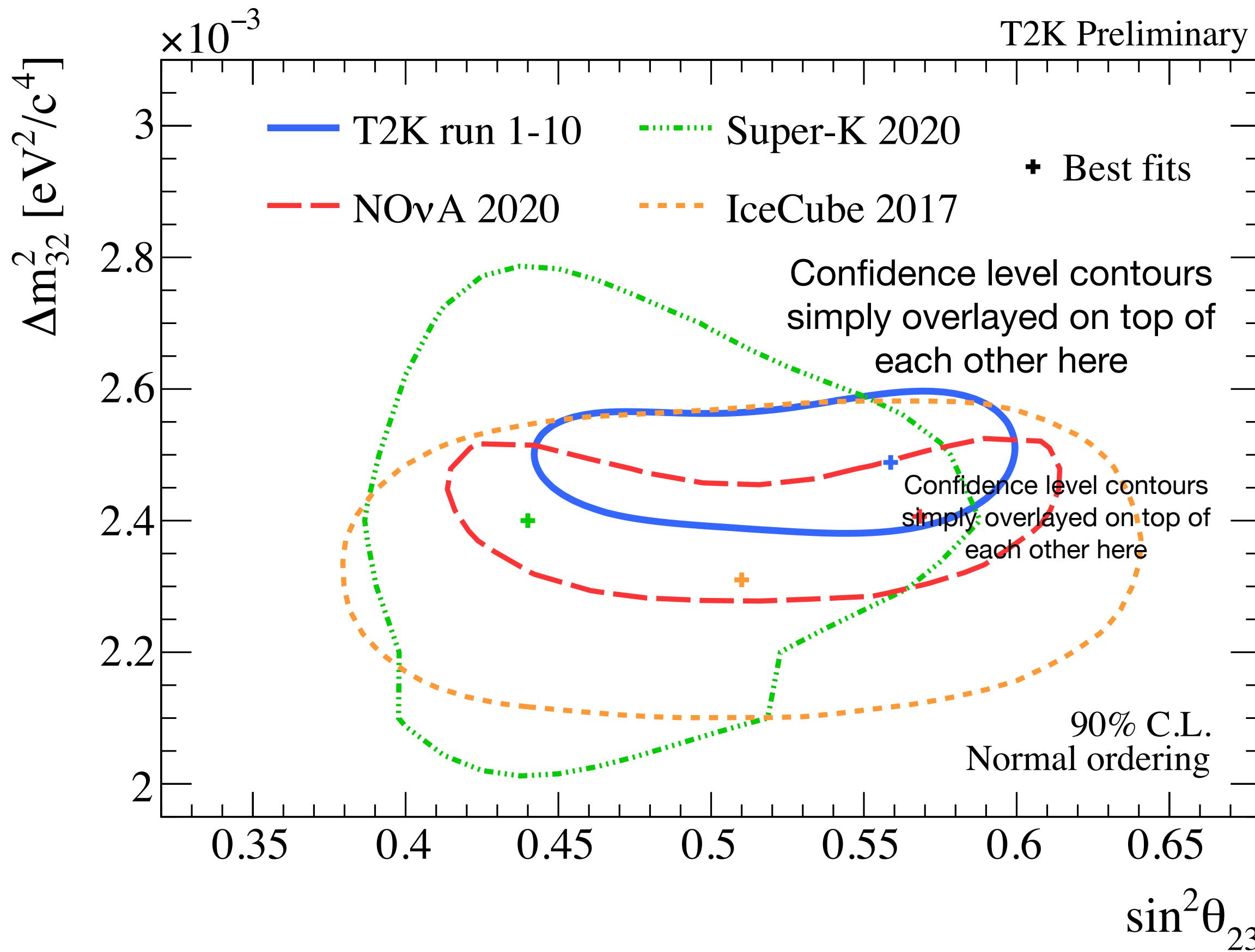
# Bi-Probability plots



- Change in electron neutrino event rate  $\sim o(45\%)$  between  $\delta_{CP} = \pi/2$  and  $\delta_{CP} = -\pi/2$
- Latest result is slightly closer to the PMNS prediction than last year



# Combined oscillation parameter knowledge: current landscape



NOvA 2020: results from [here](#)

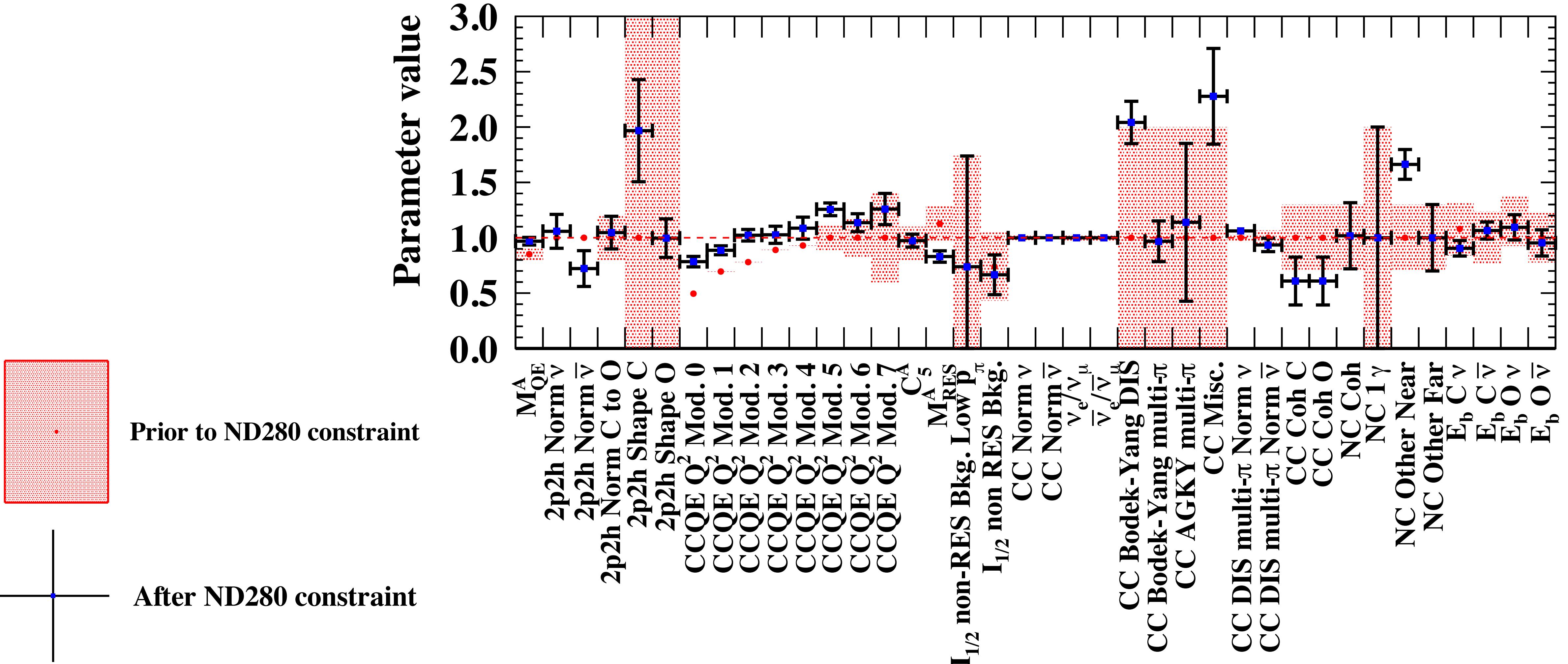
Super-K 2020: preliminary results from [here](#)

IceCube 2017: results from [here](#)

# Fits to near detector data: cross section model parameters

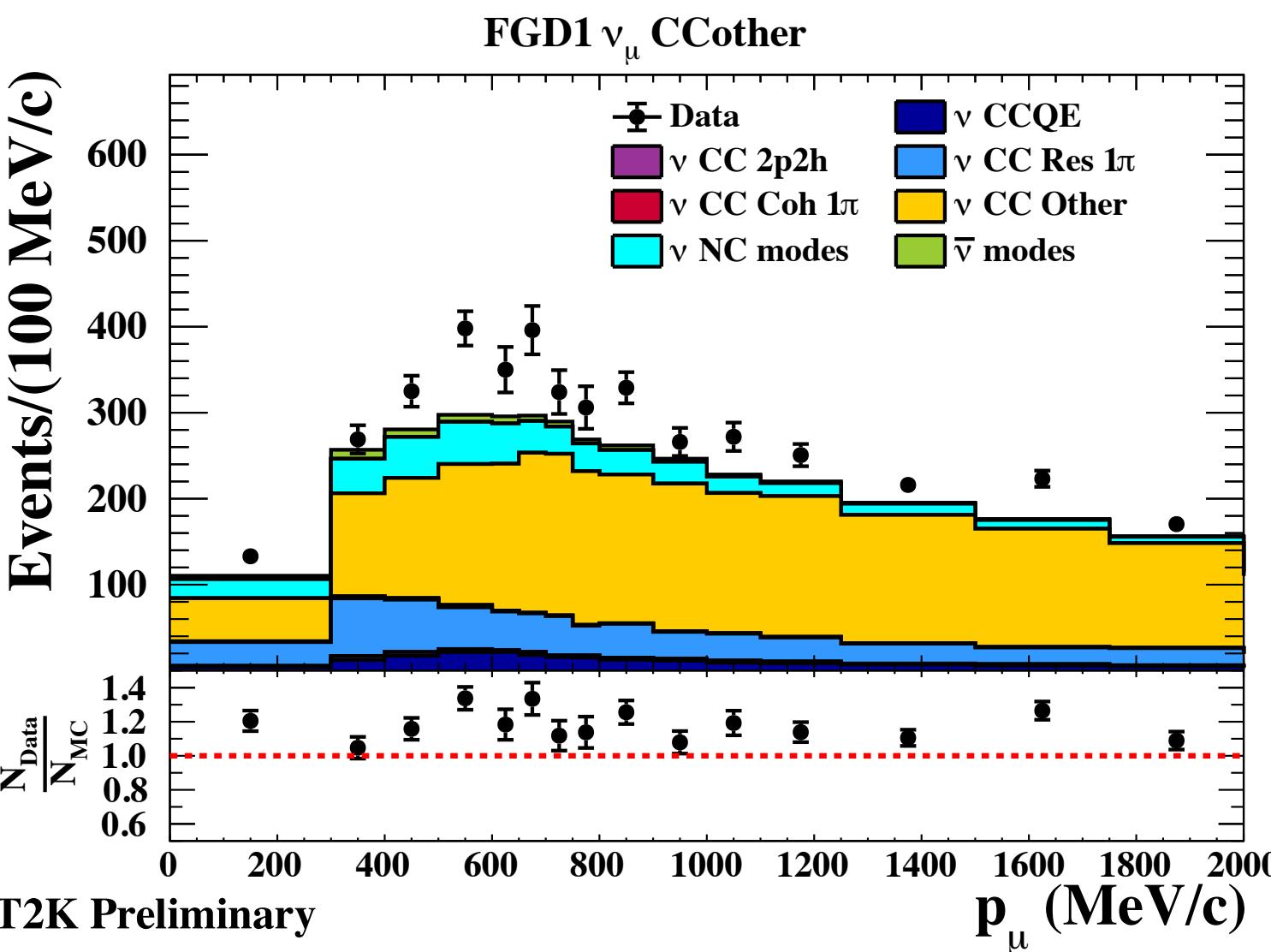
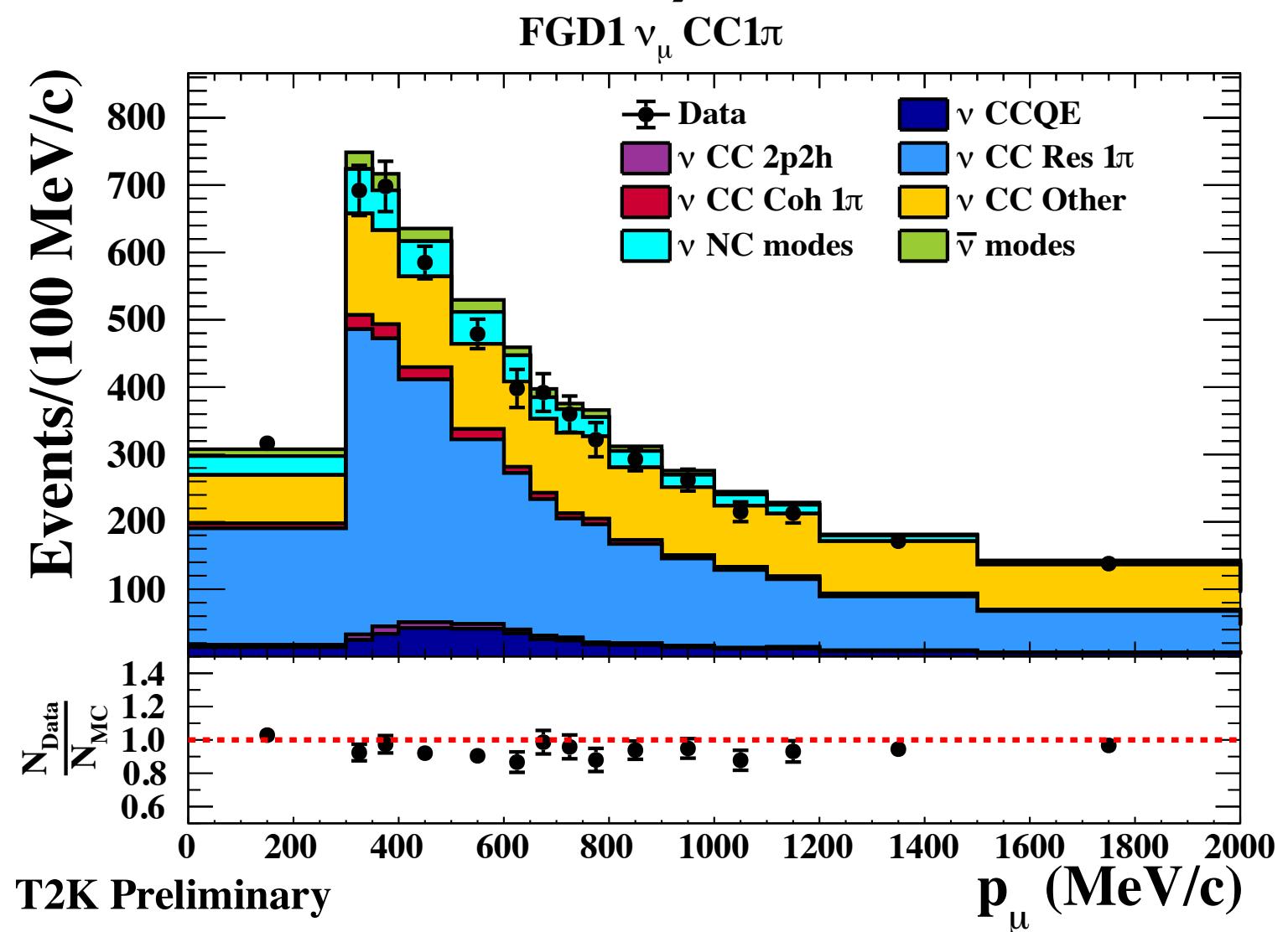
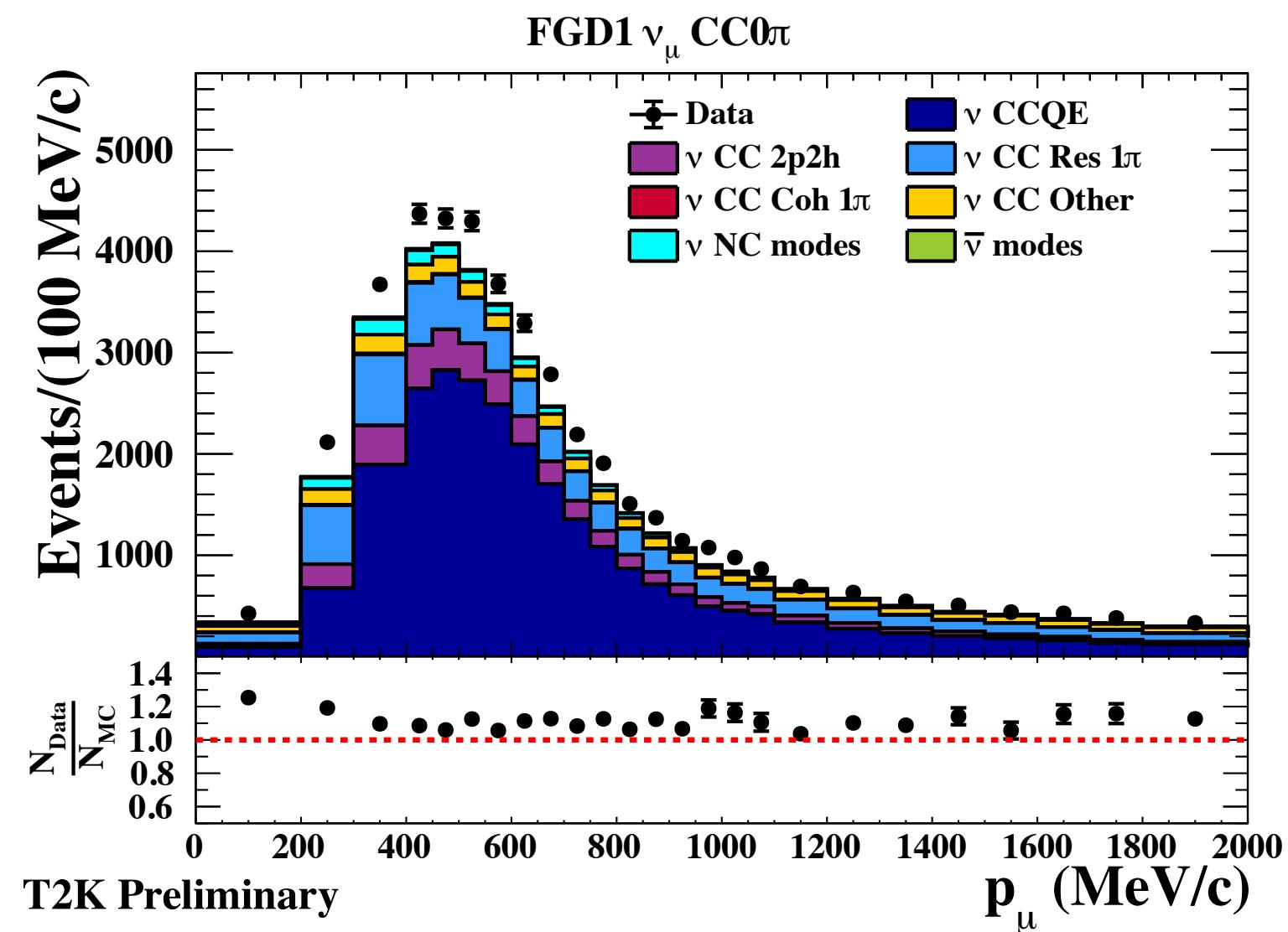
Cross-section

T2K Preliminary

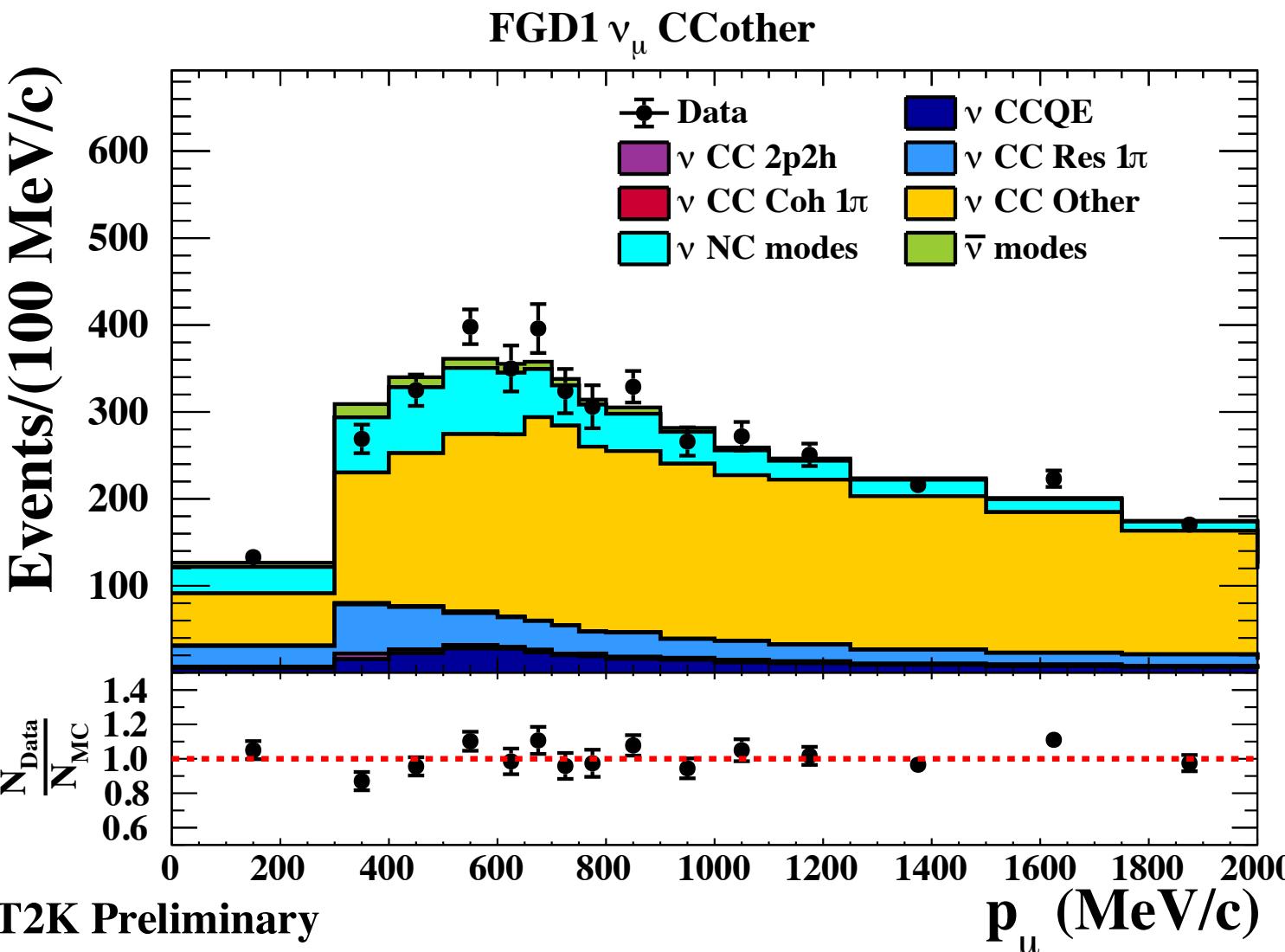
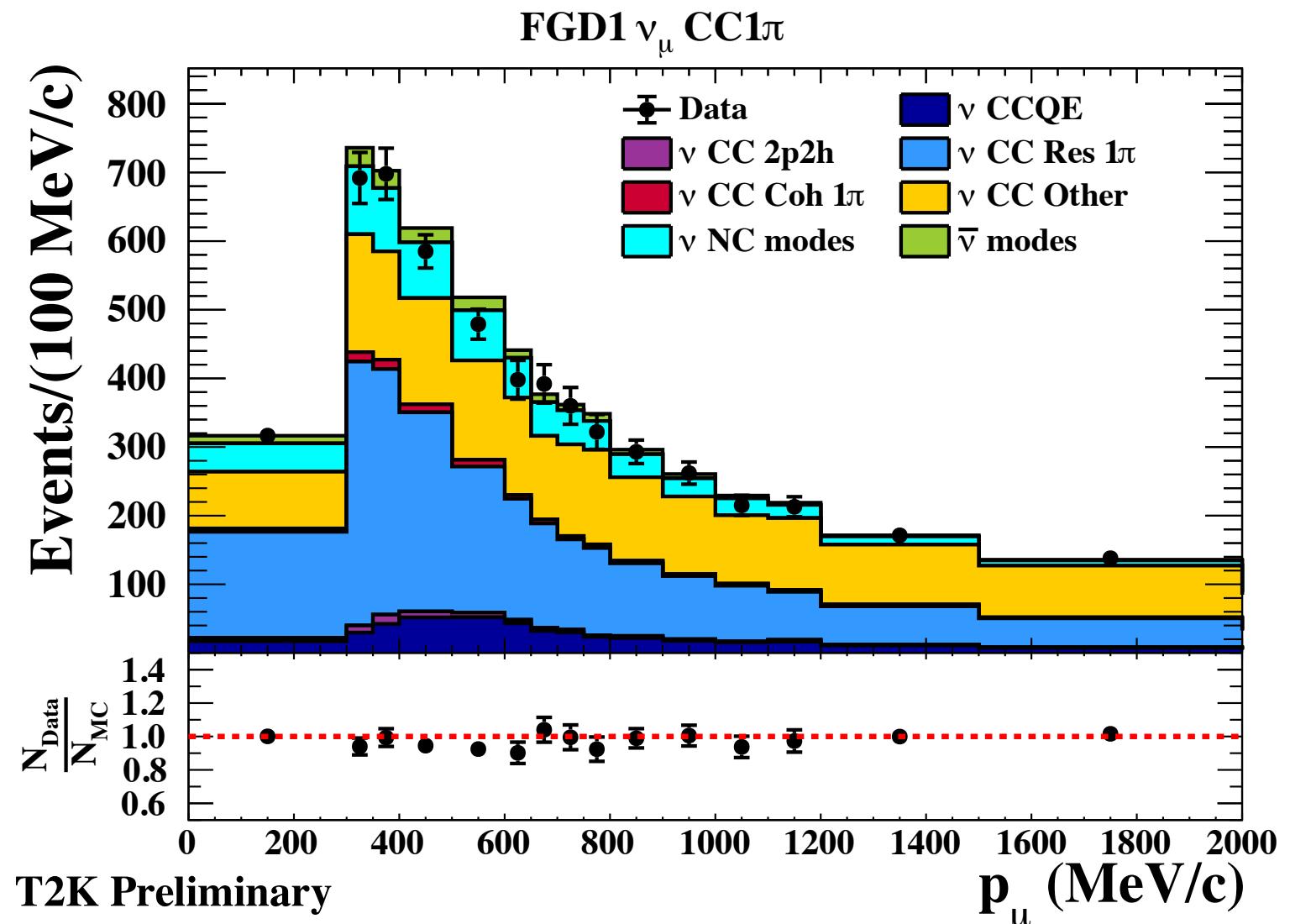
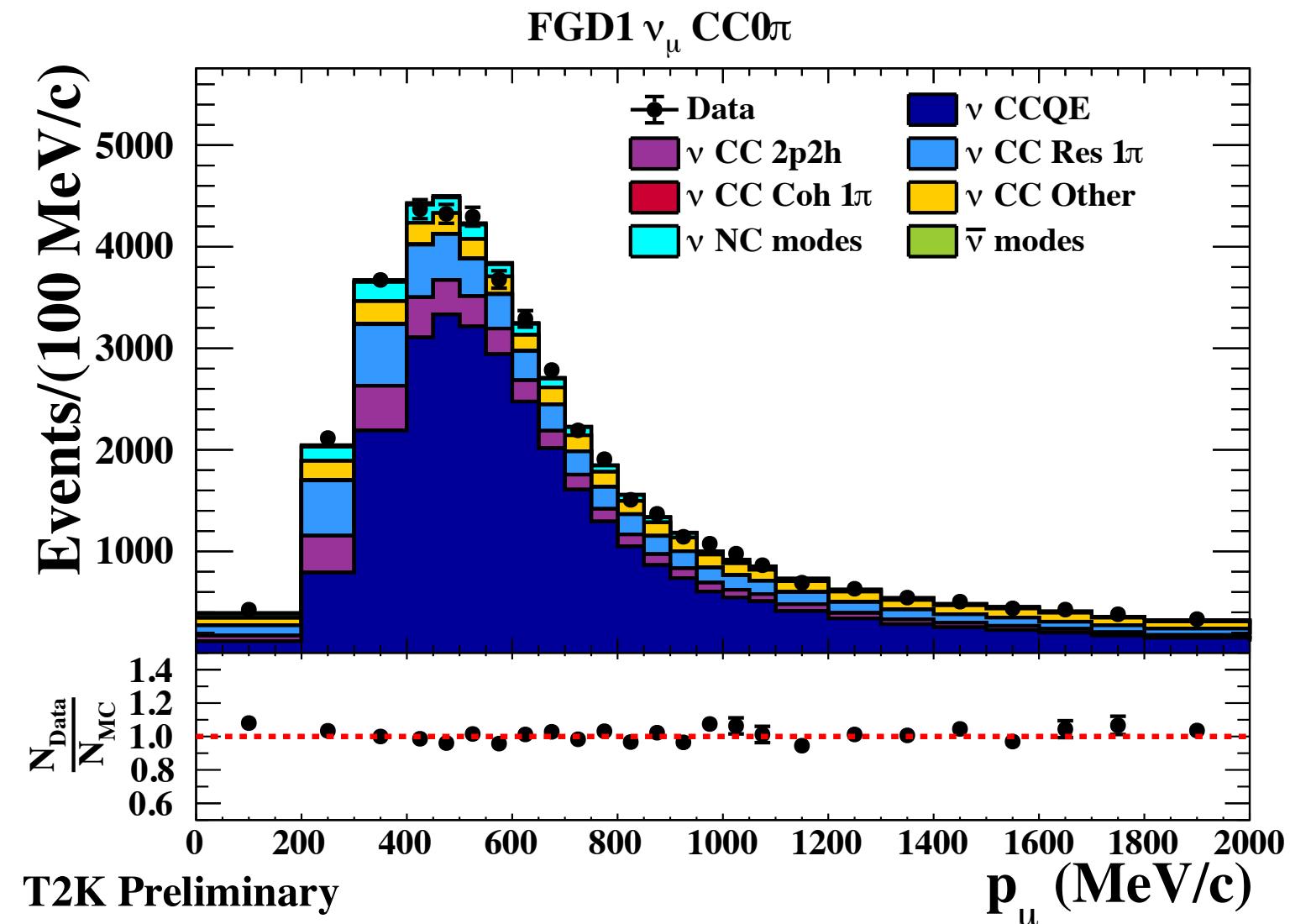


# Fits to near detector FGD1 $\nu_\mu$ signal data in $\nu$ -mode

*Before ND fit*

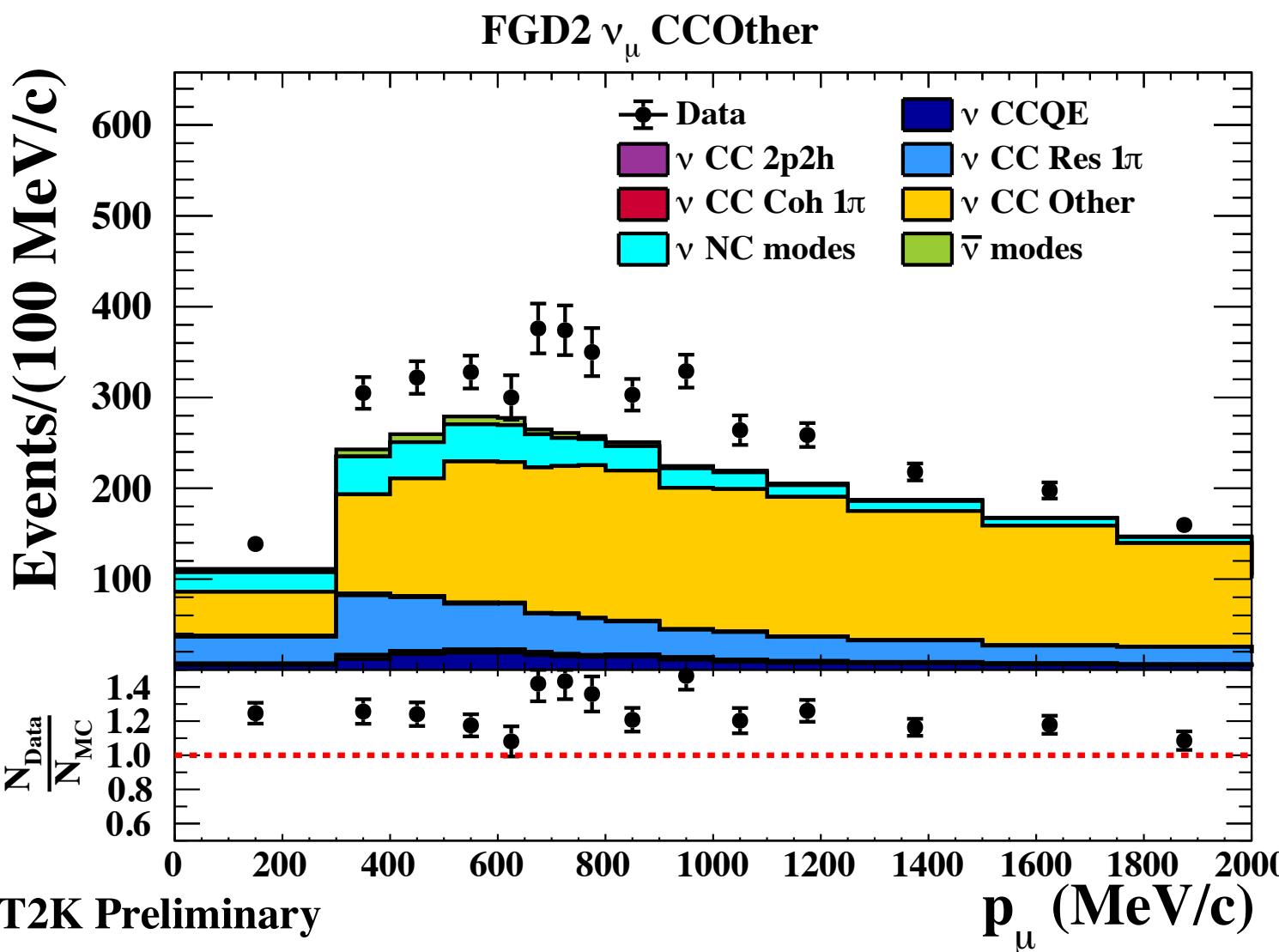
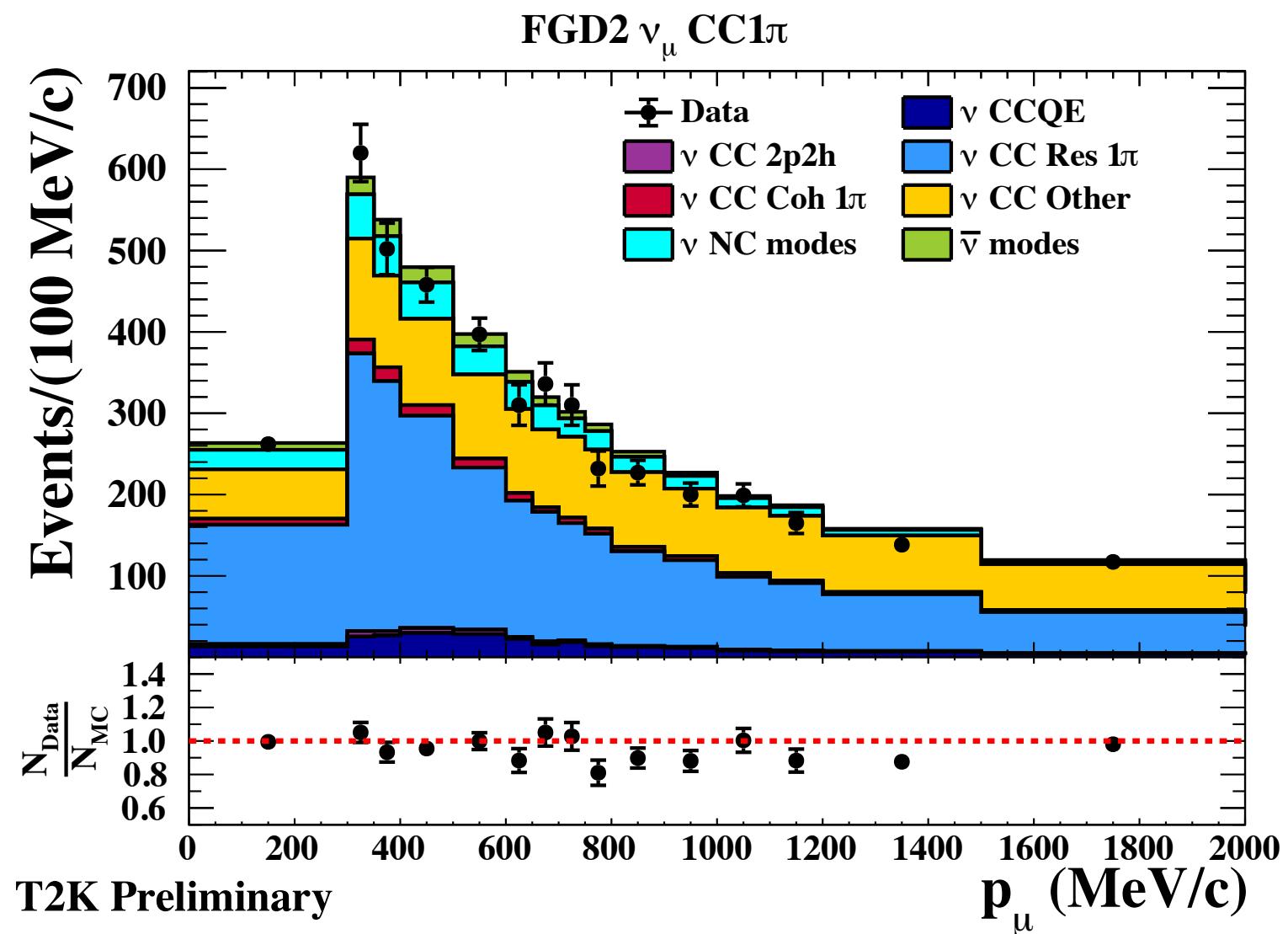
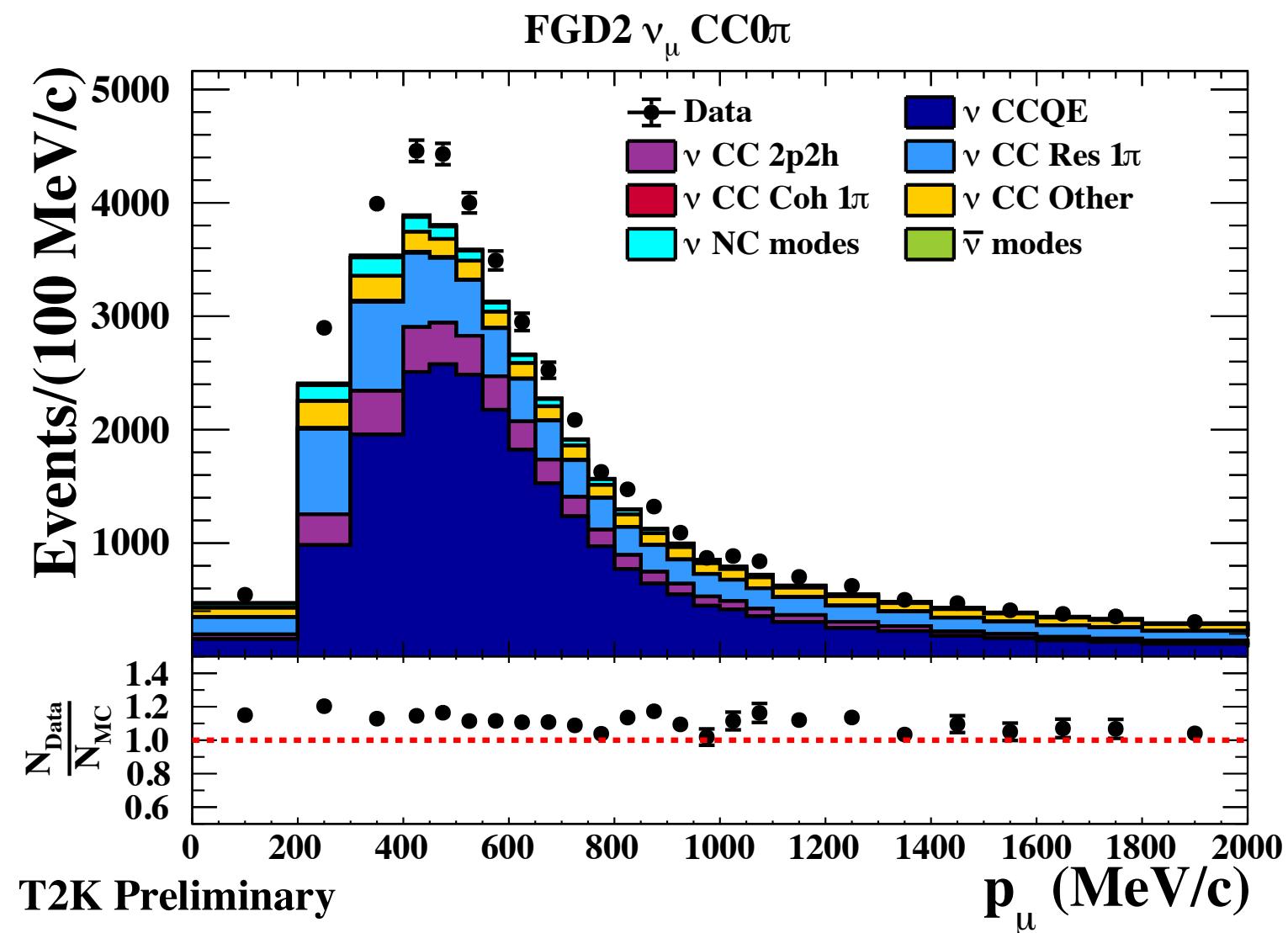


*After ND fit*

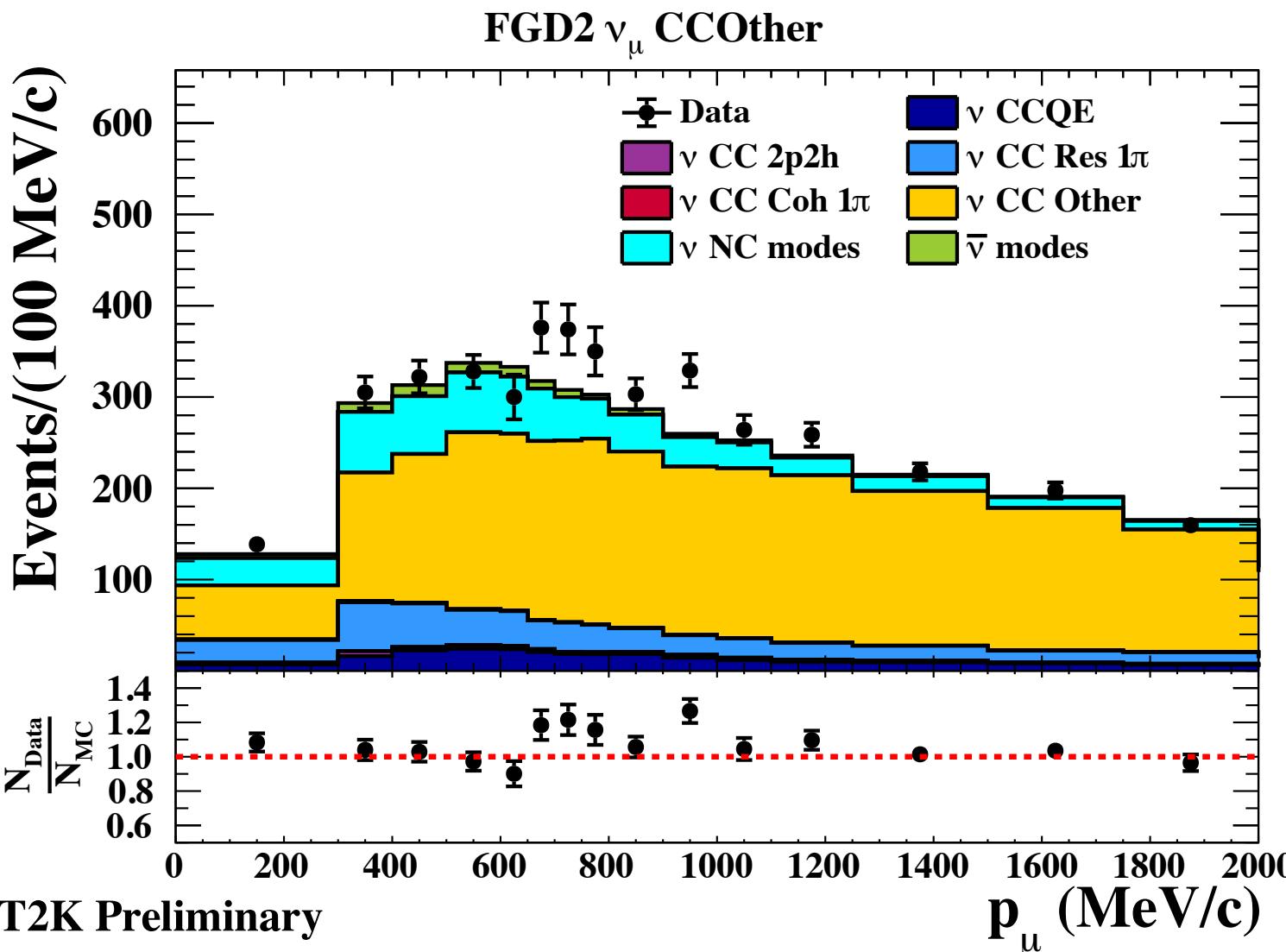
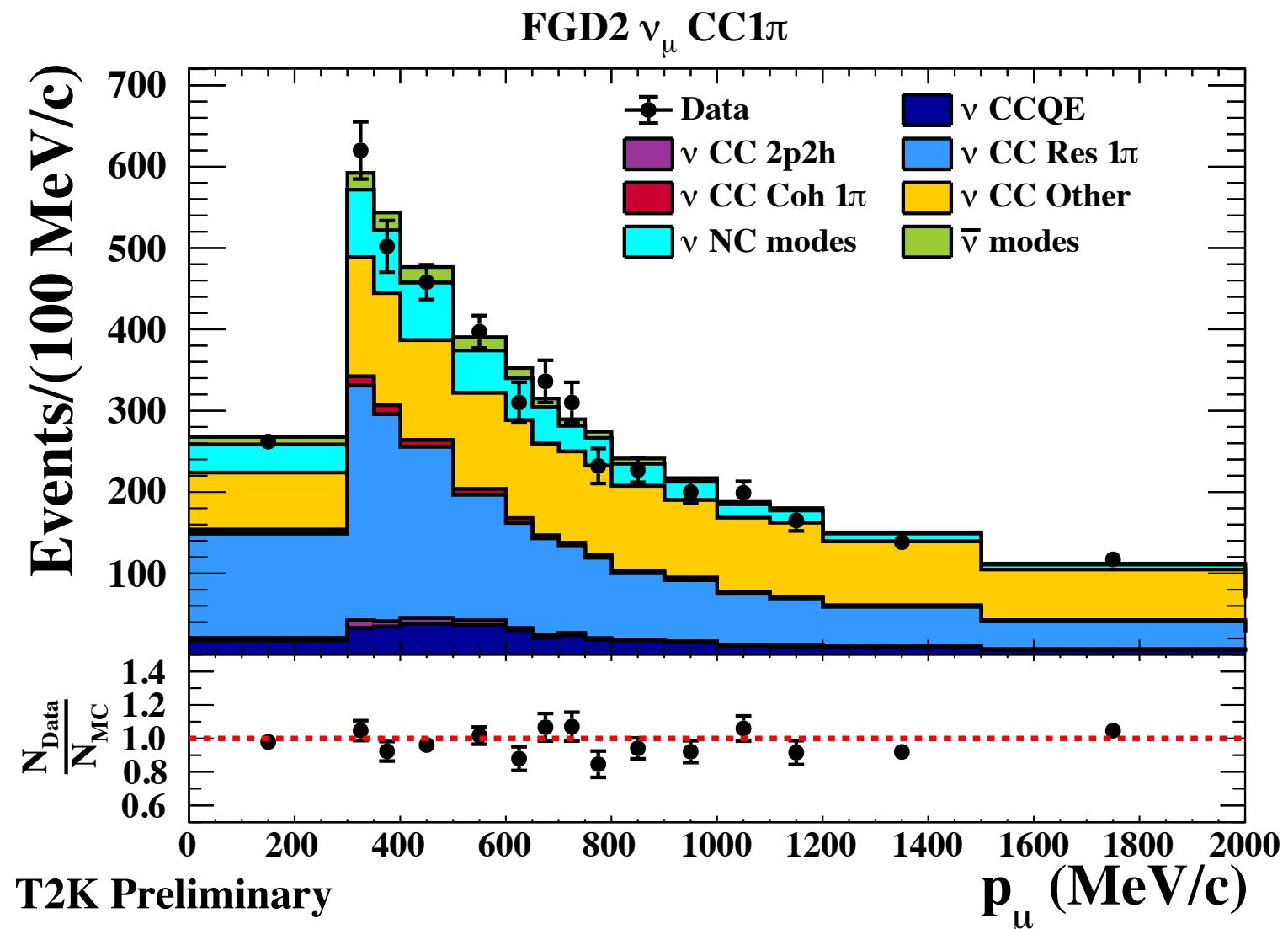
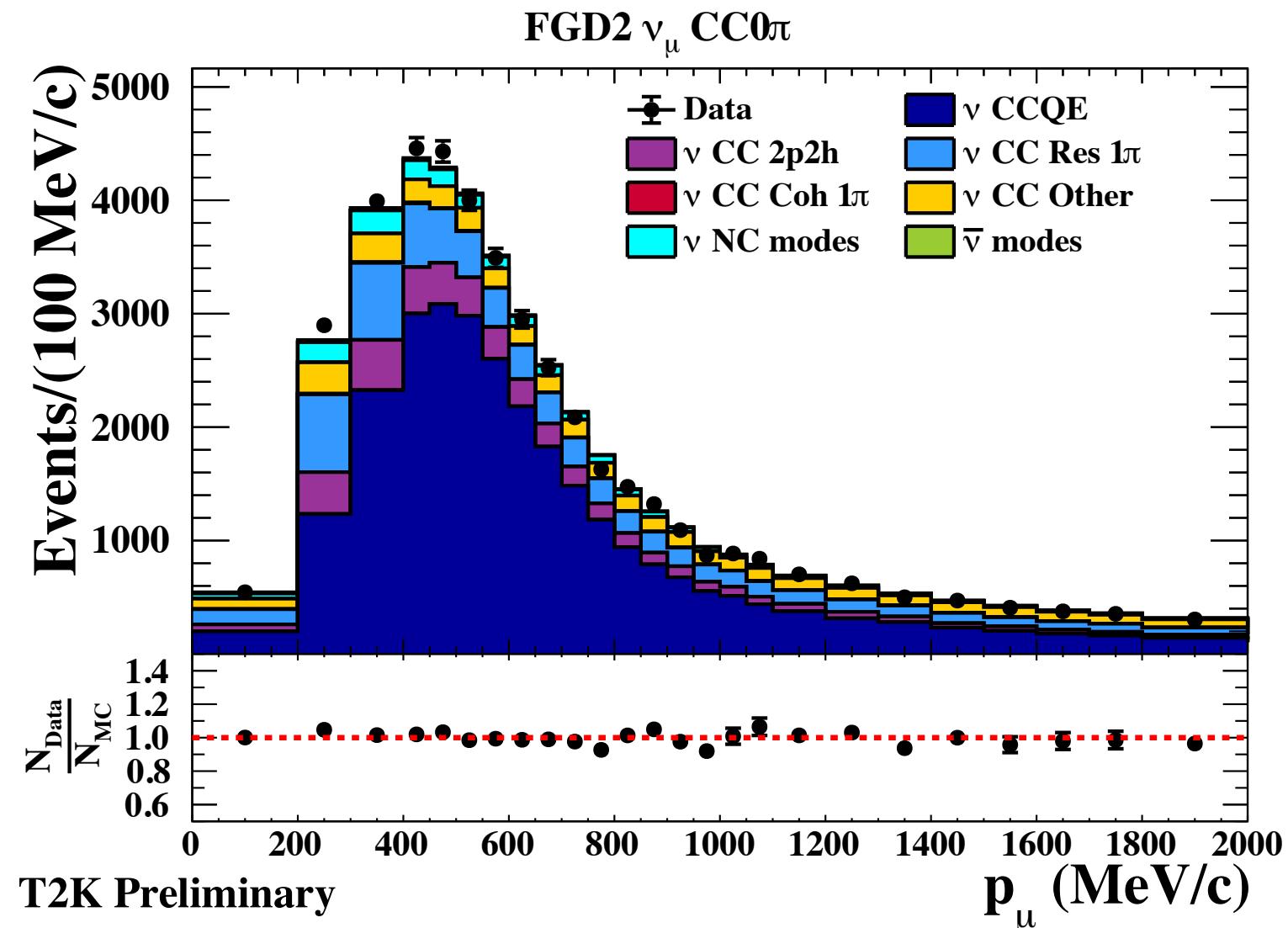


# Fits to near detector FGD2 $\nu_\mu$ signal data in $\nu$ -mode

*Before ND fit*

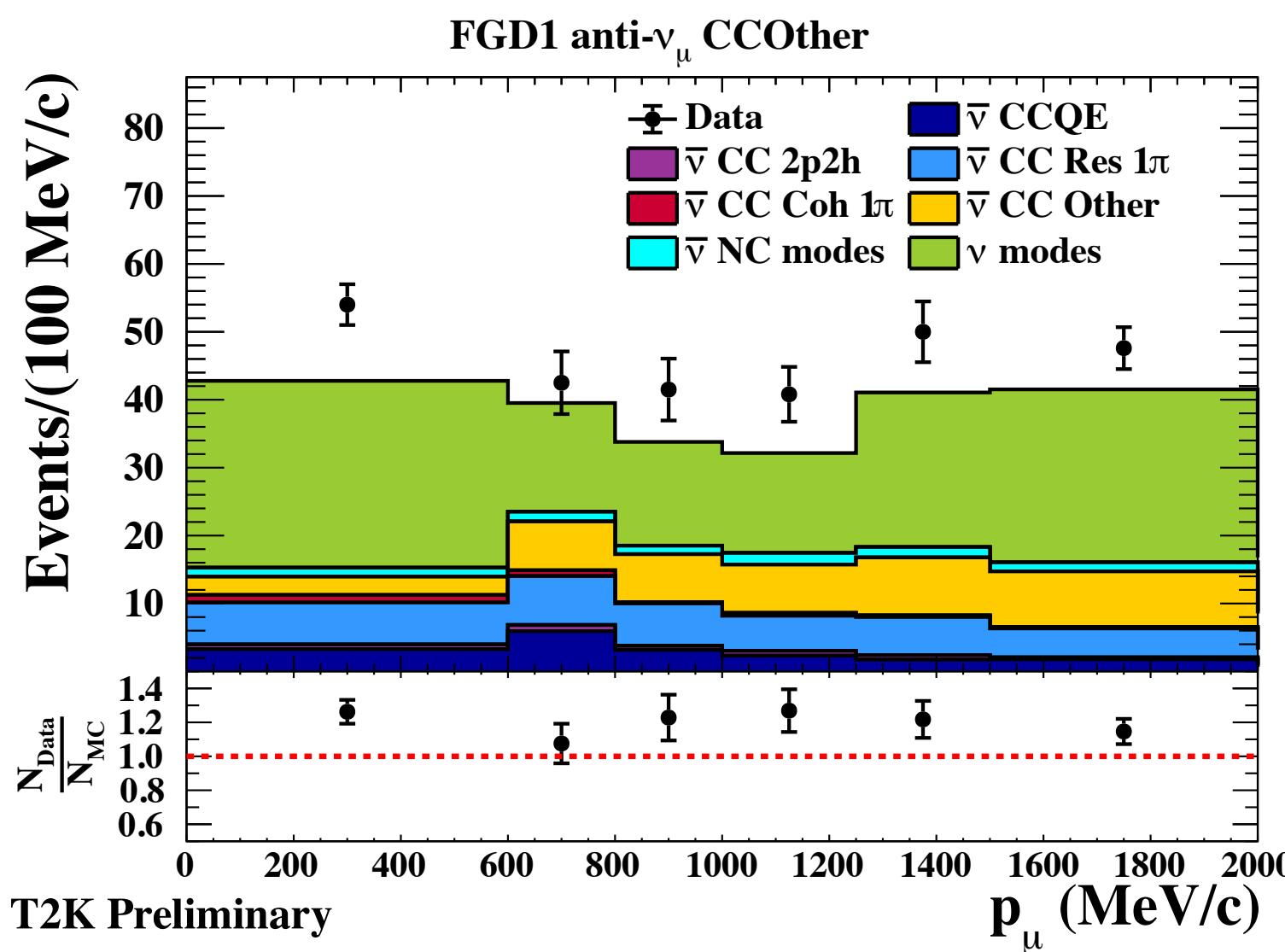
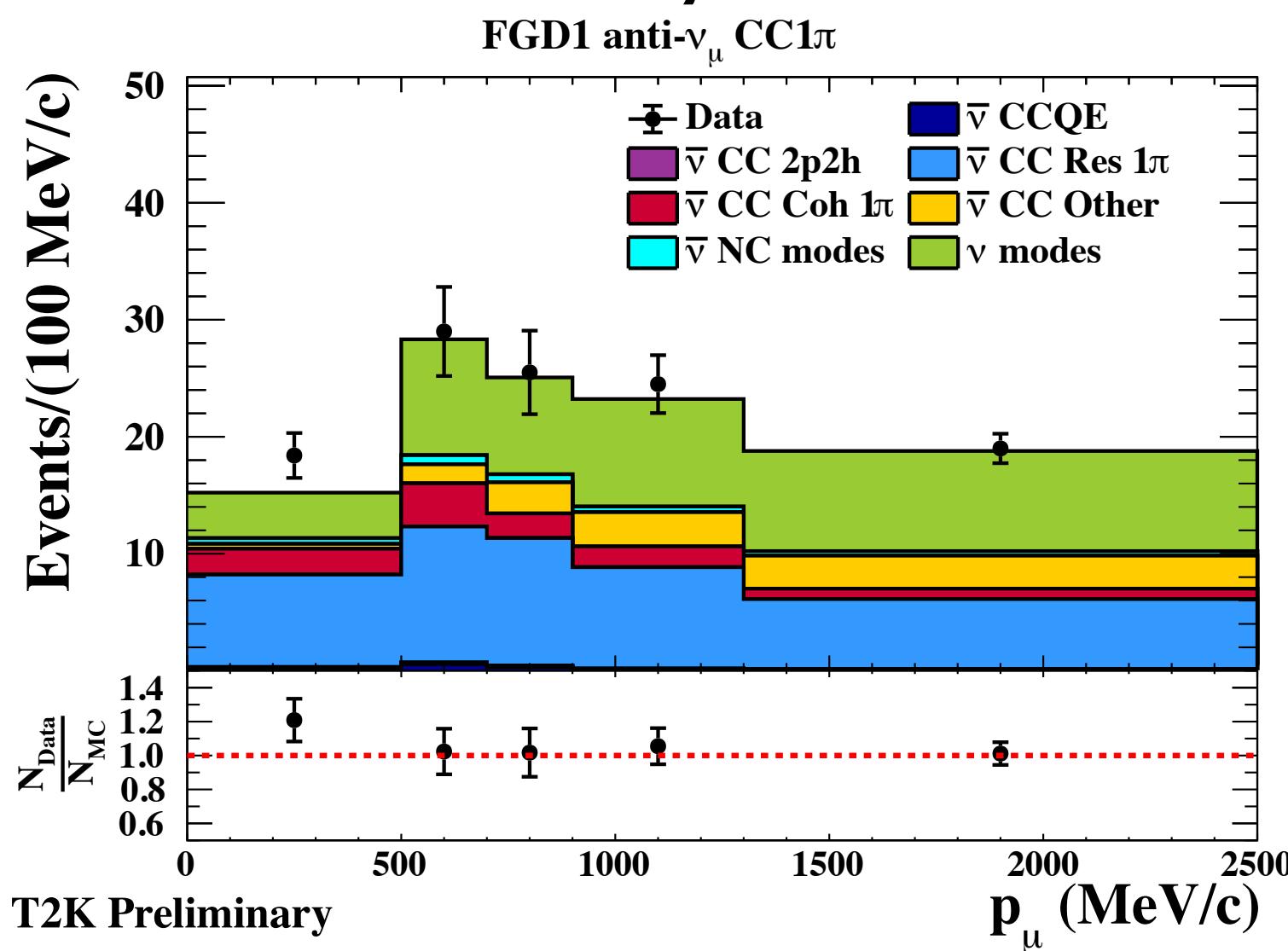
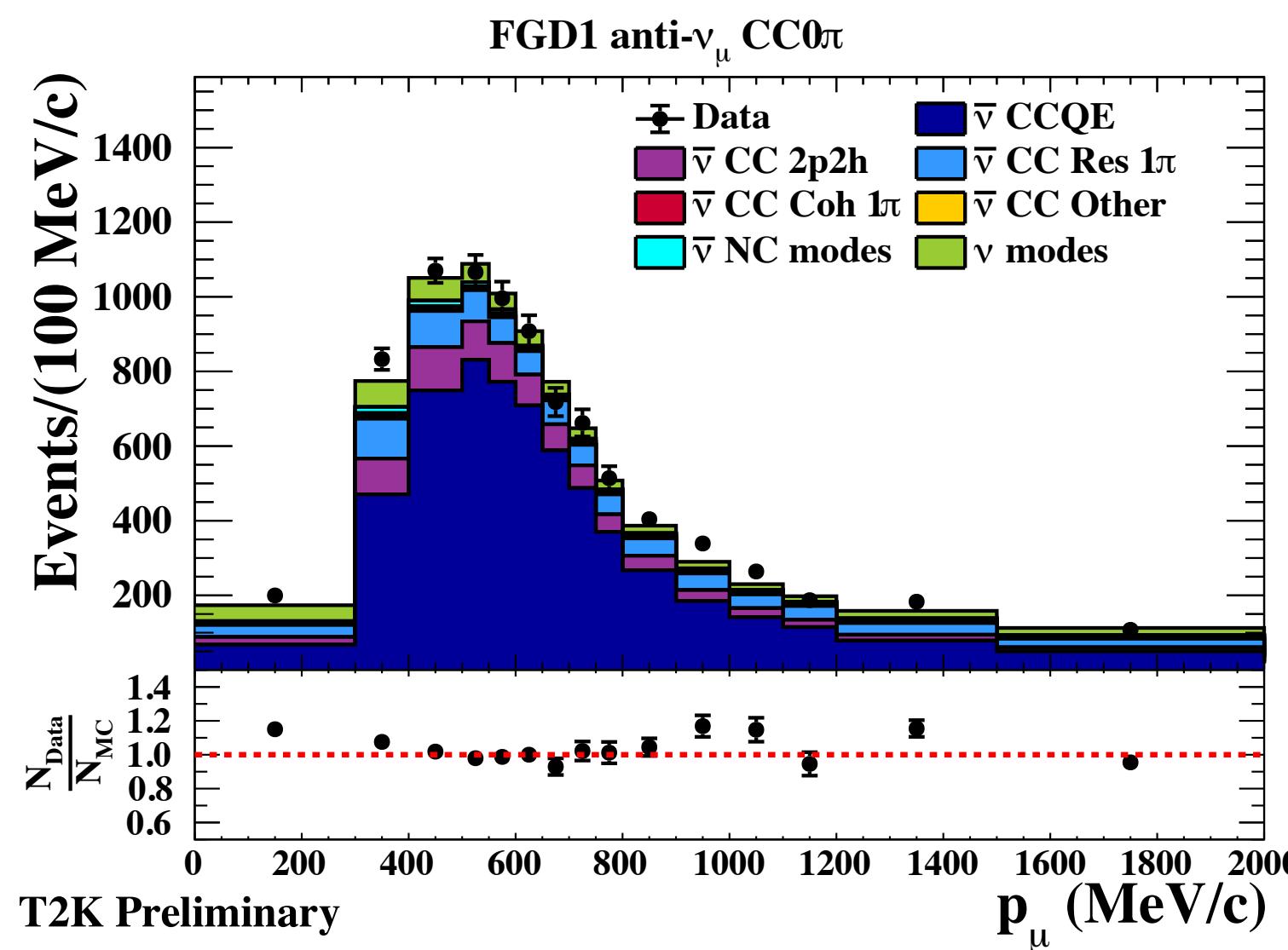


*After ND fit*

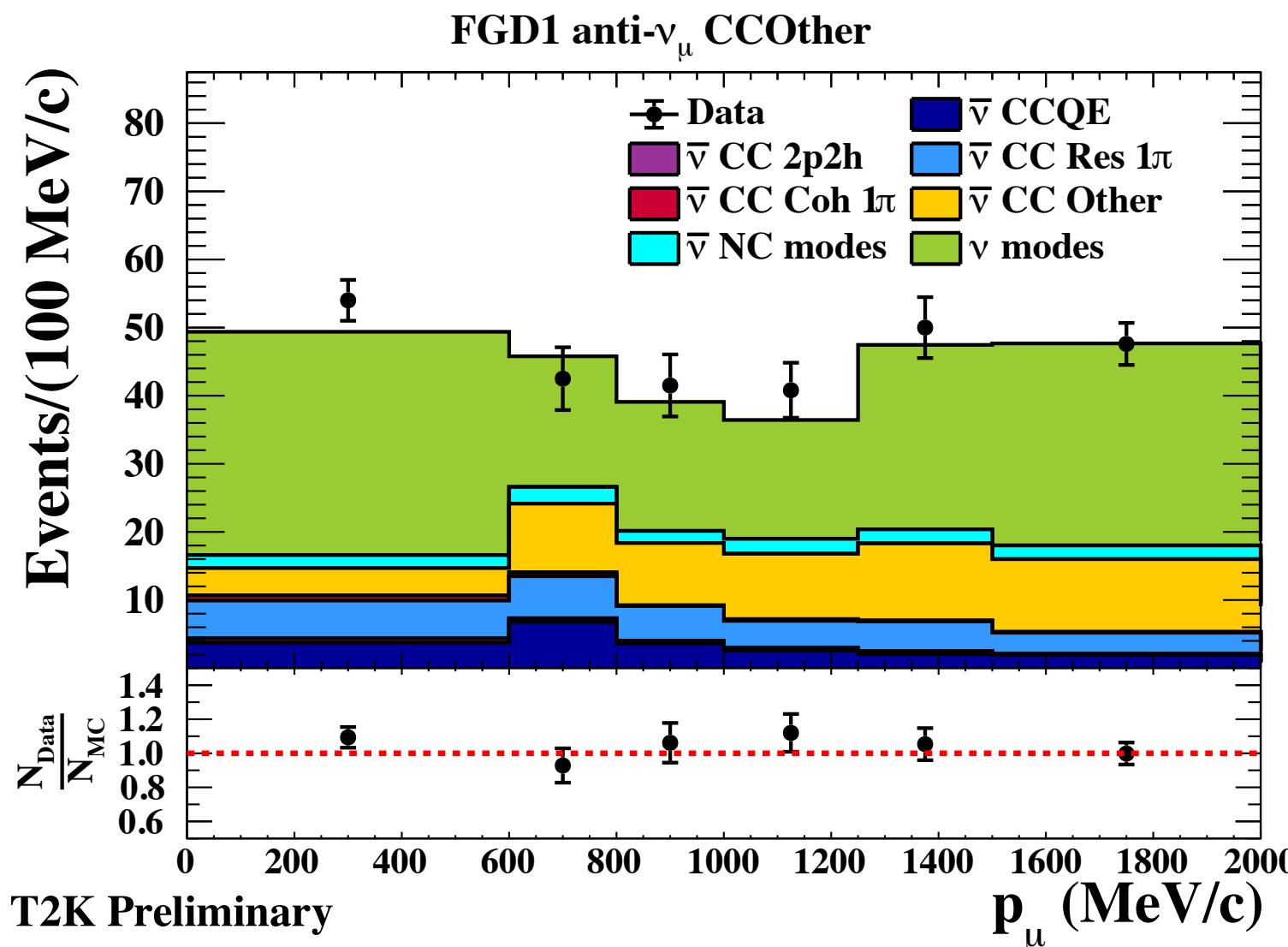
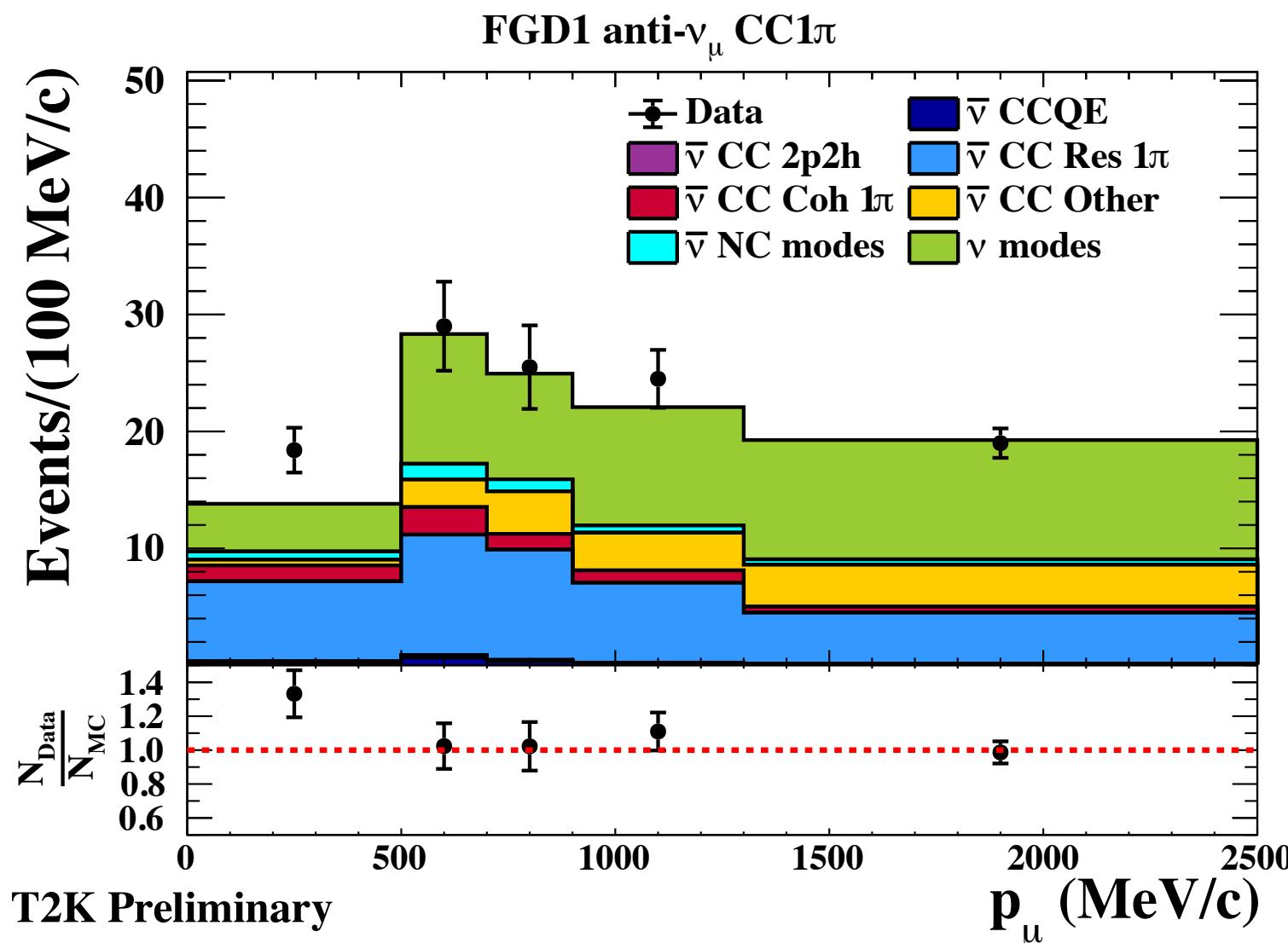
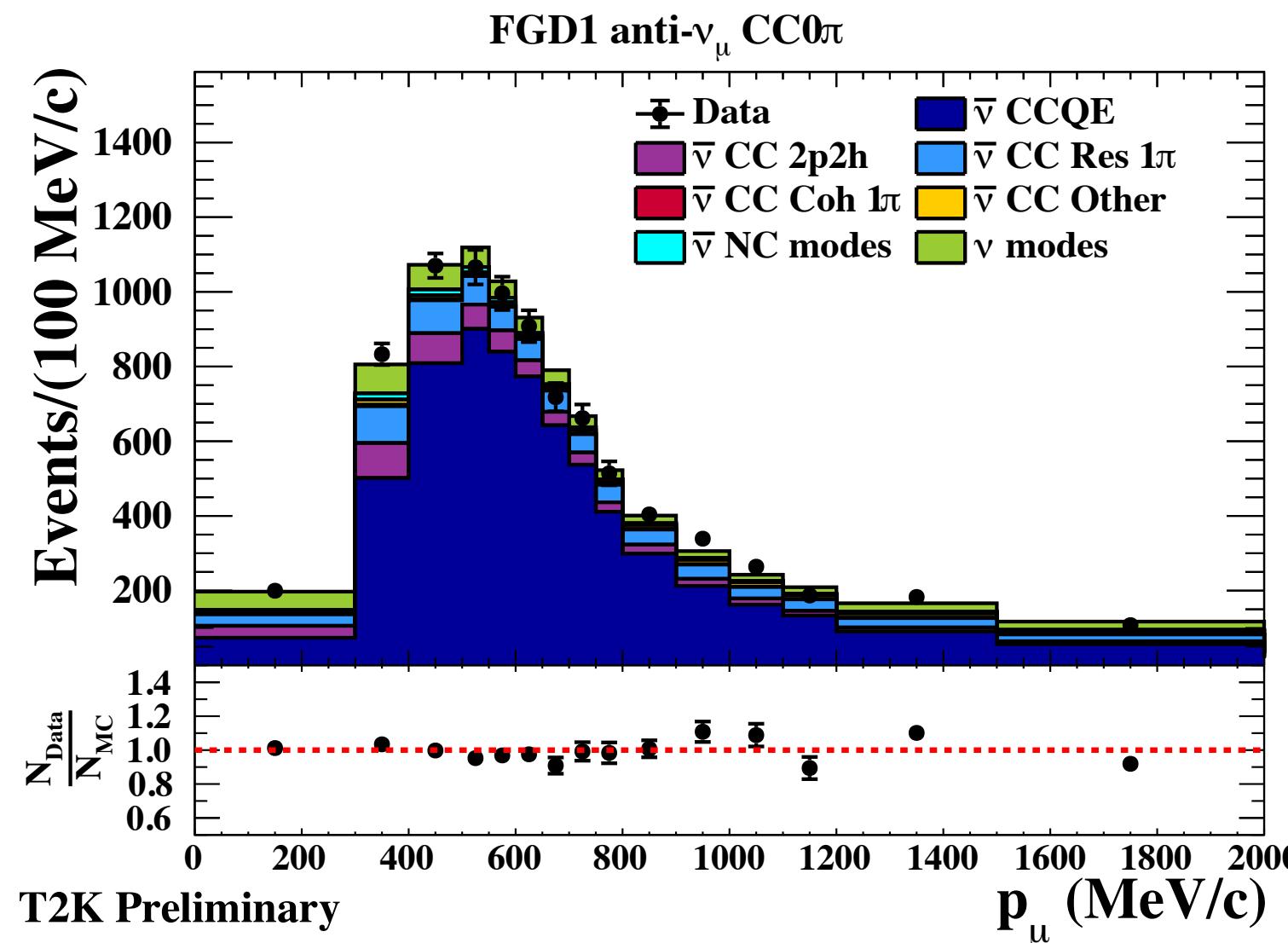


# Fits to near detector FGD1 $\bar{\nu}_\mu$ signal data in $\bar{\nu}$ -mode

Before ND fit

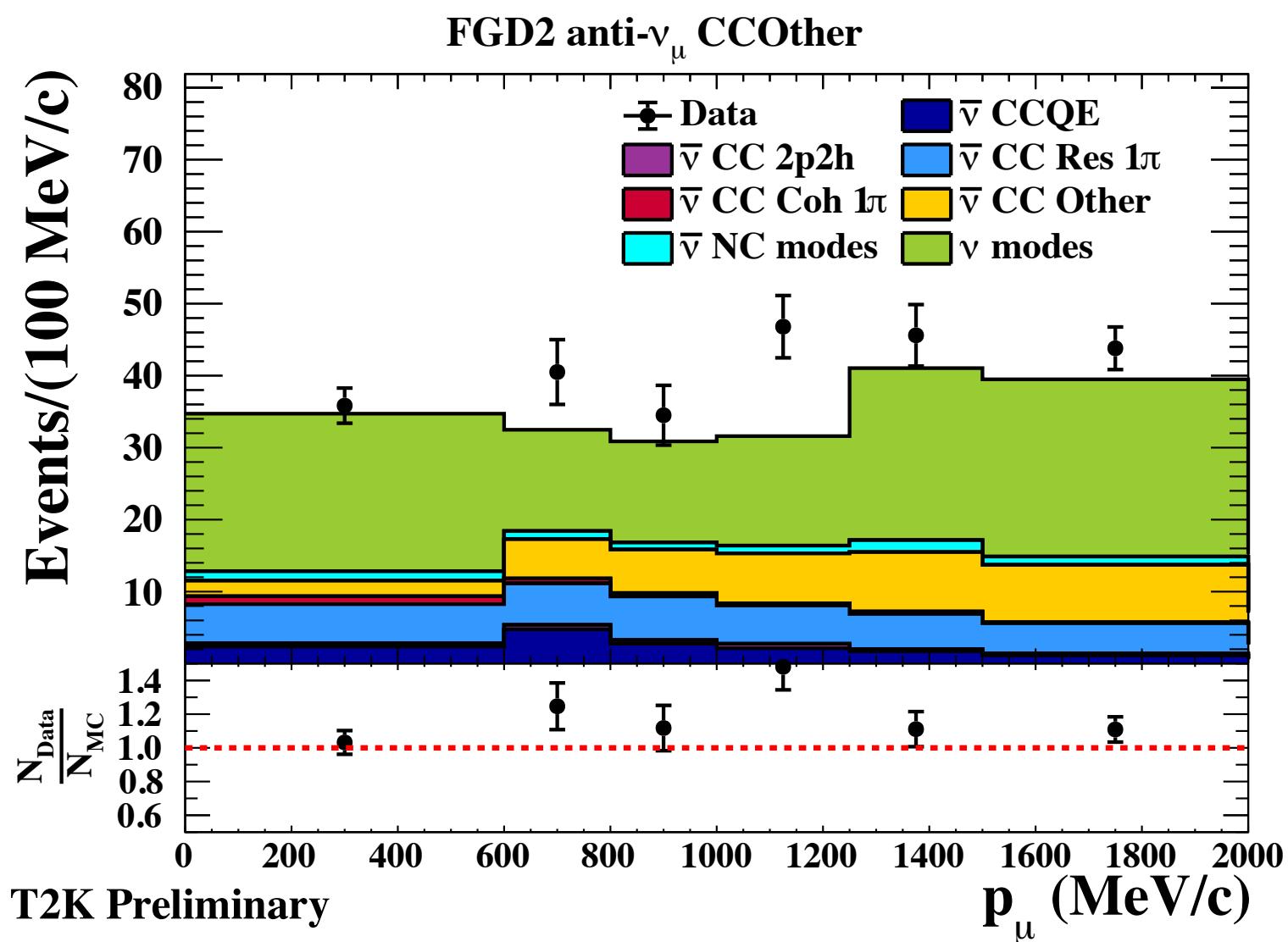
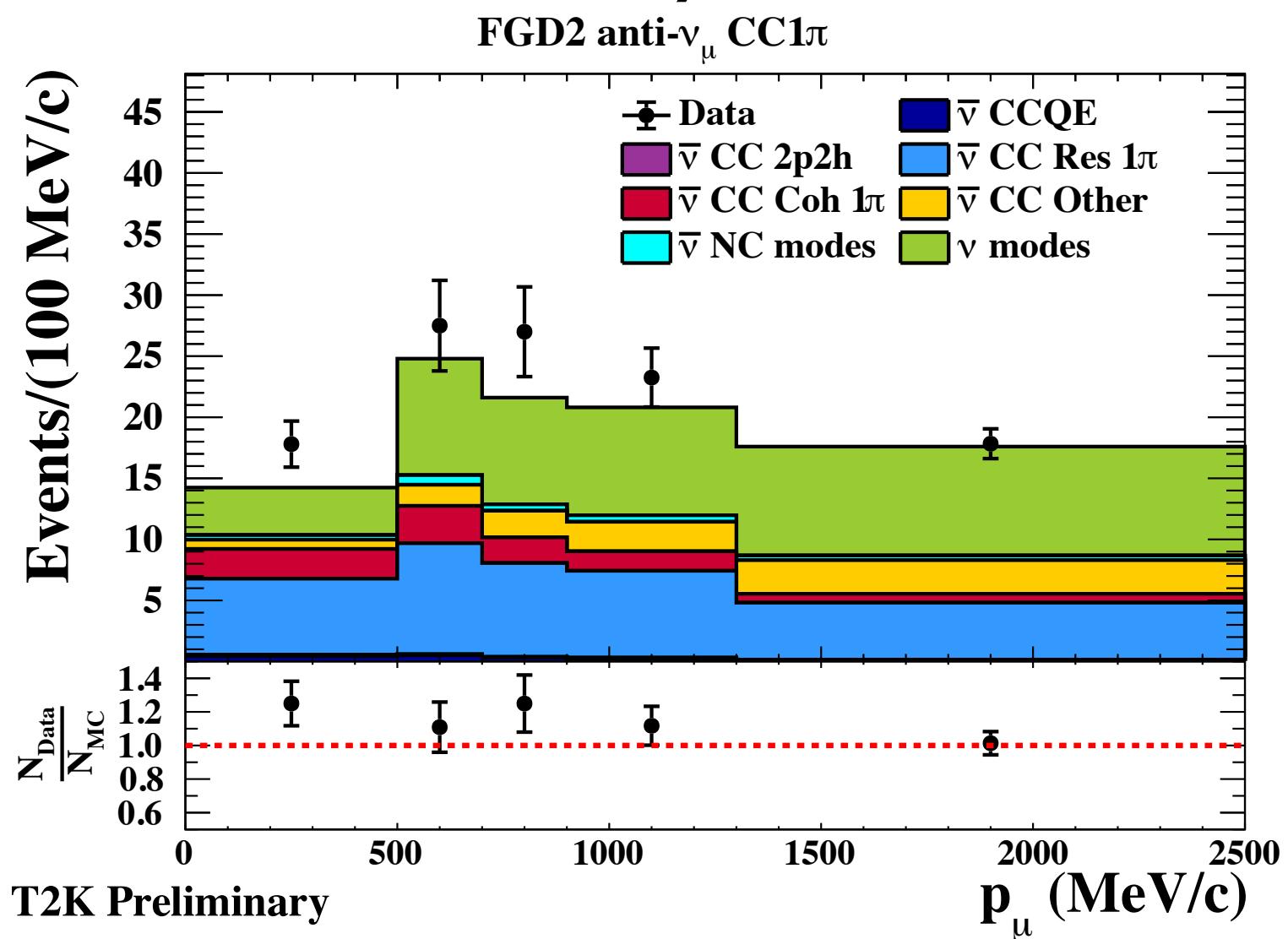
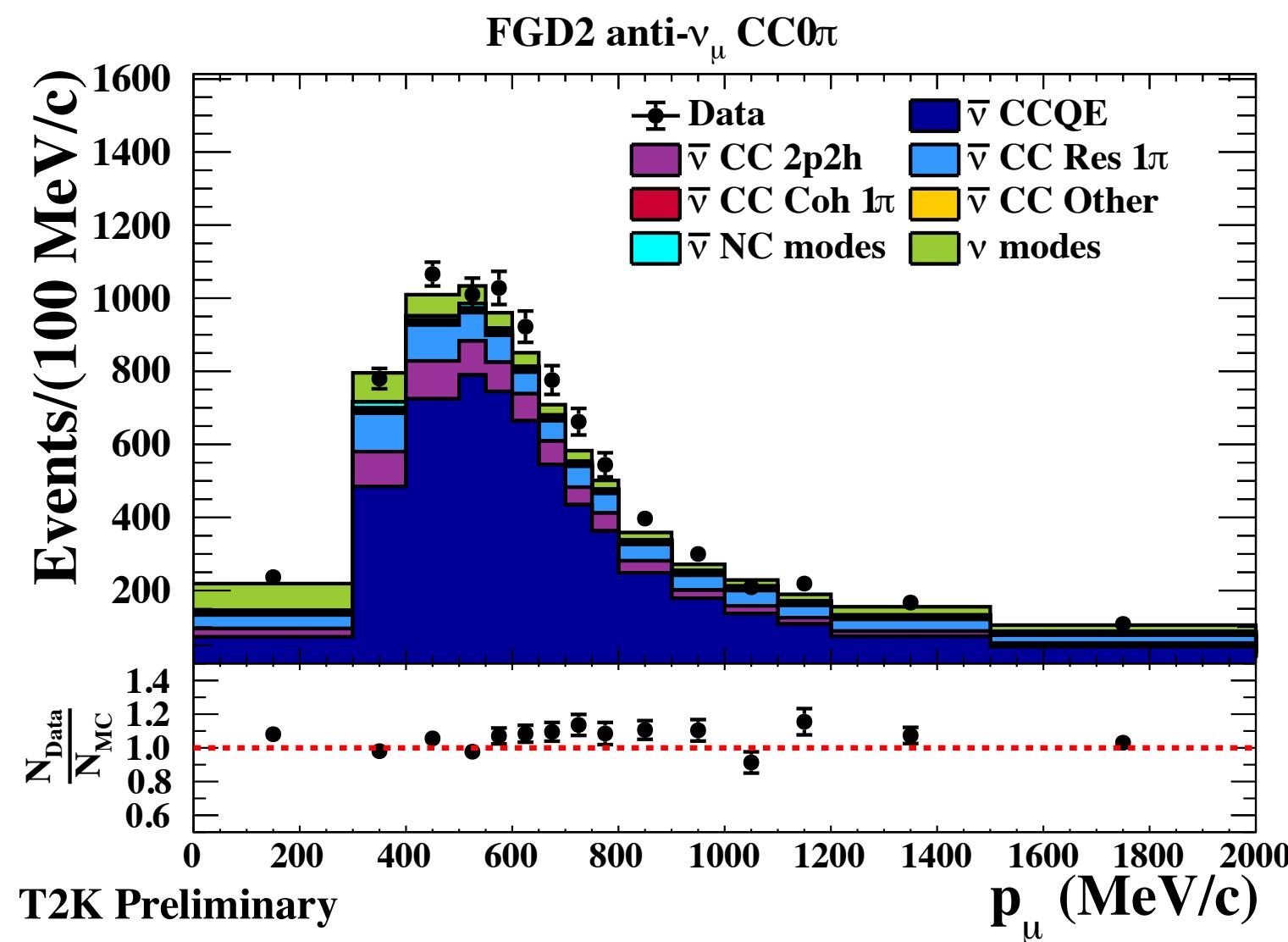


After ND fit

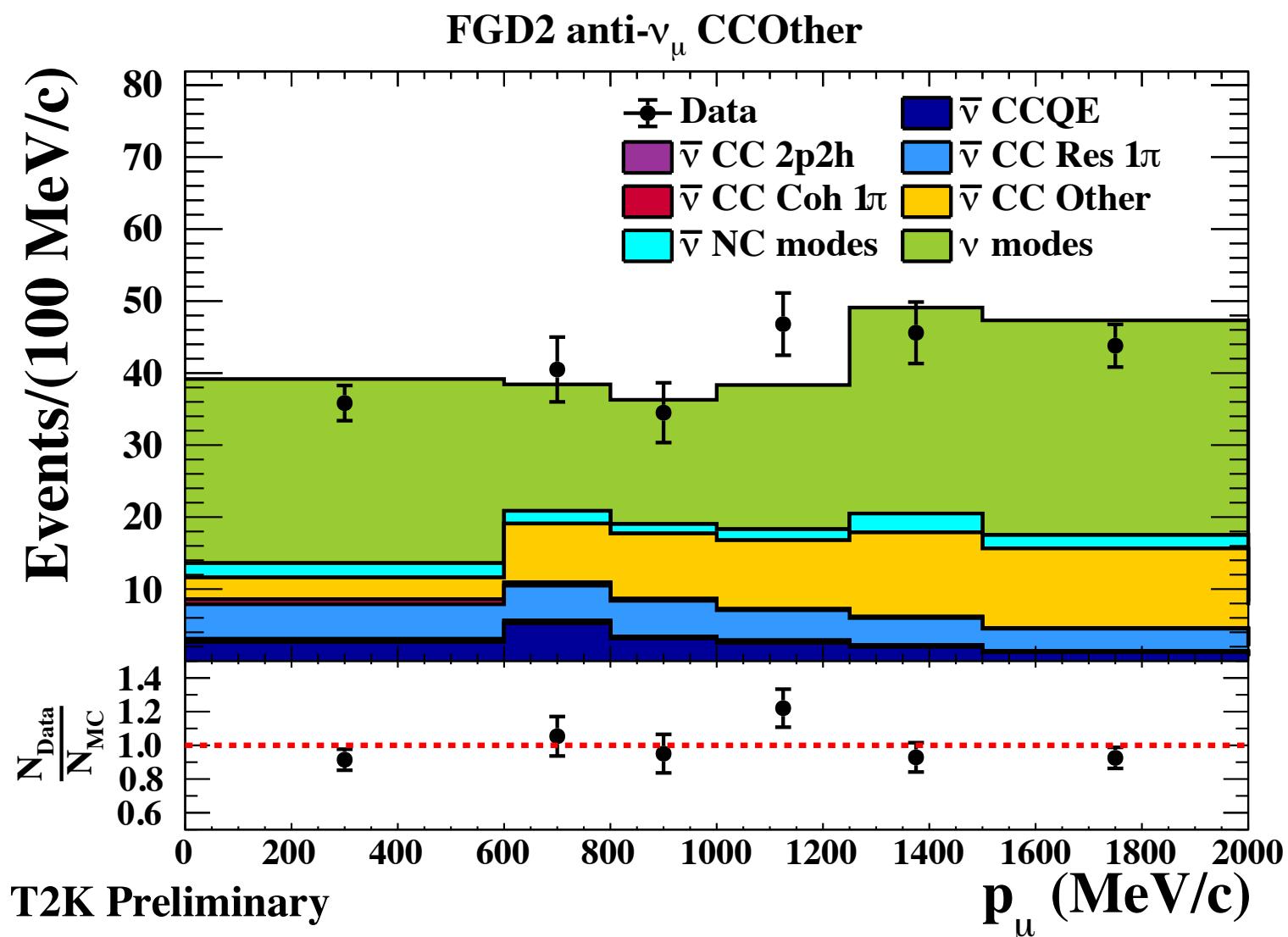
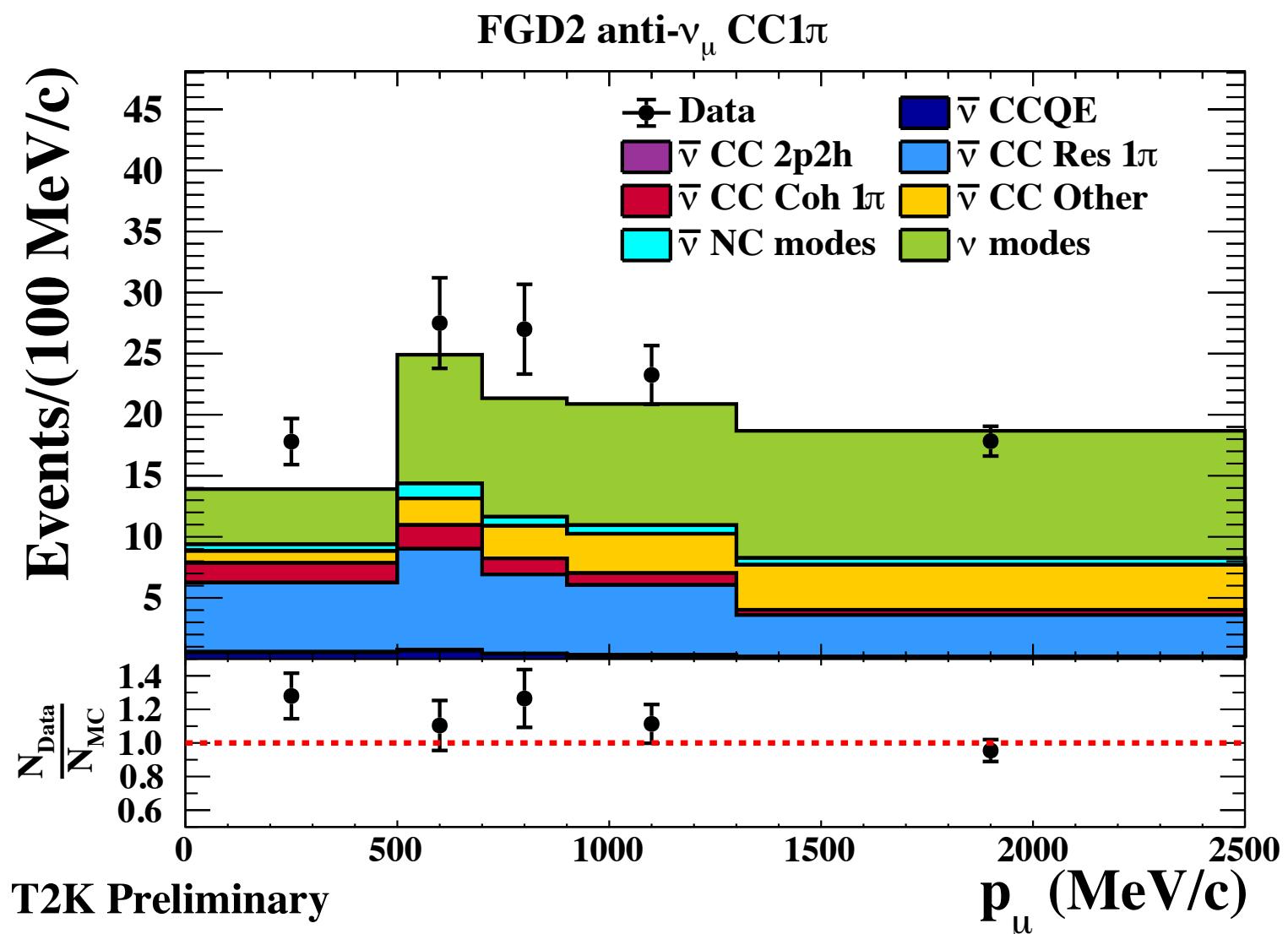
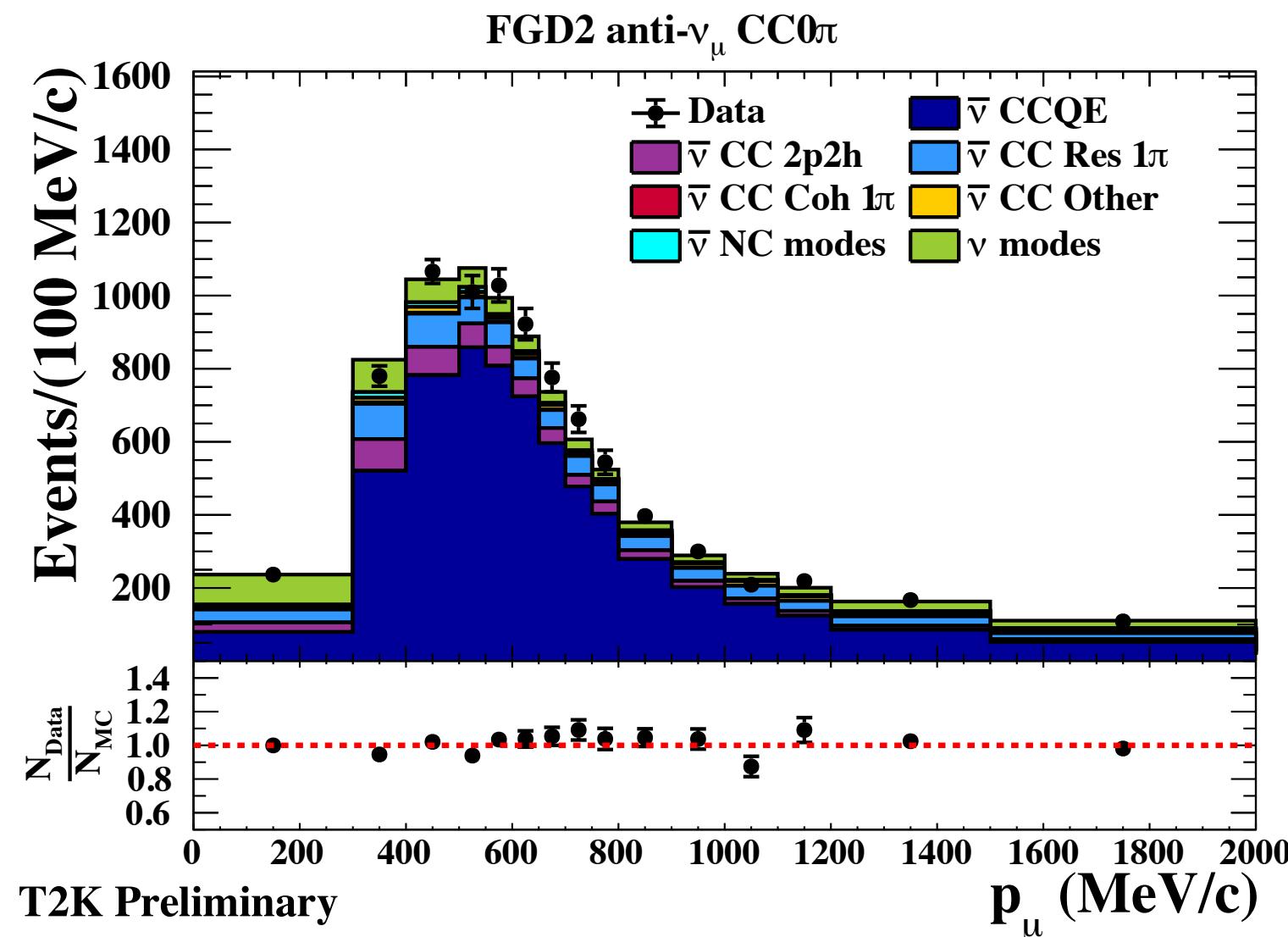


# Fits to near detector FGD2 $\bar{\nu}_\mu$ signal data in $\bar{\nu}$ -mode

*Before ND fit*

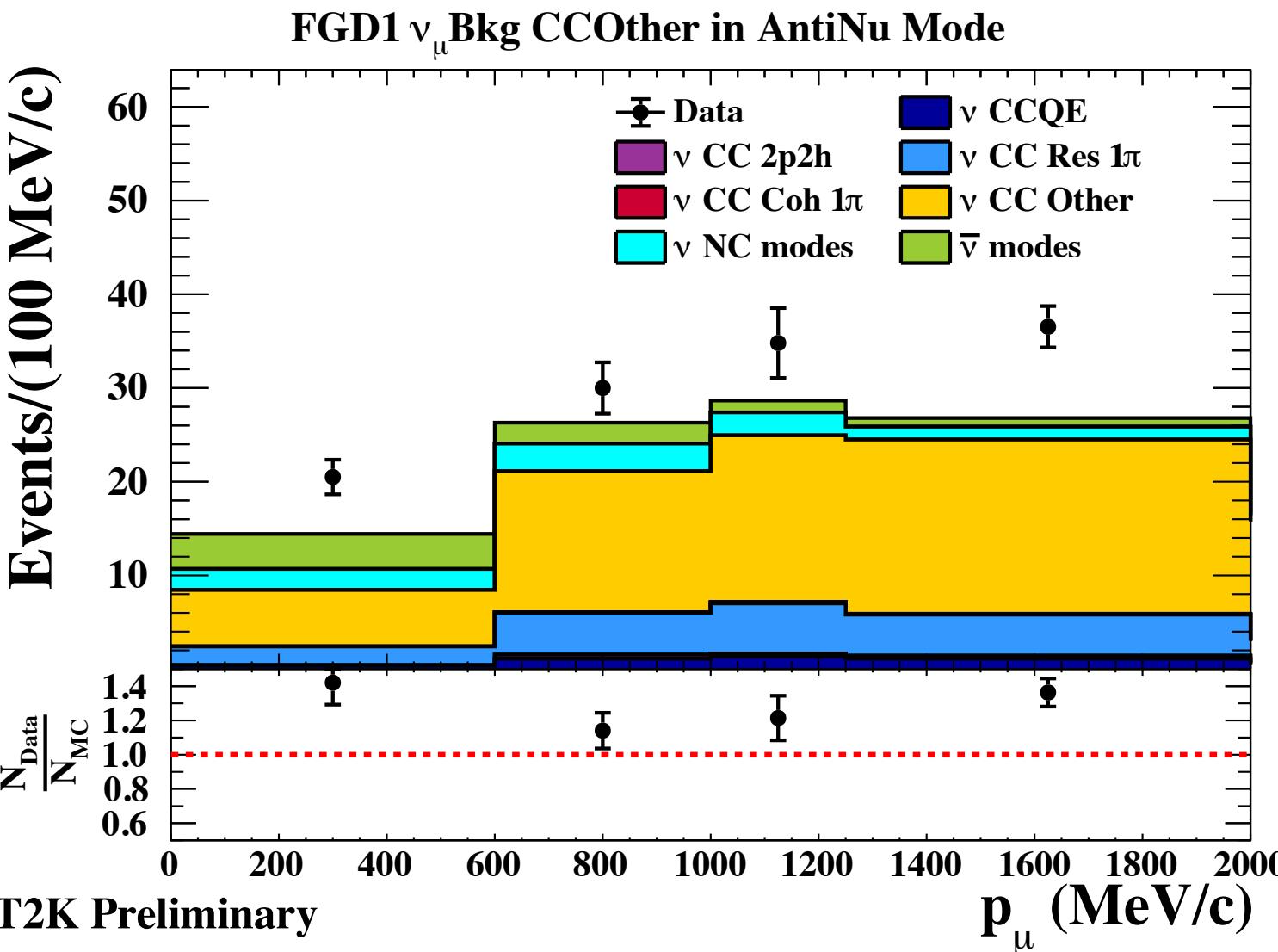
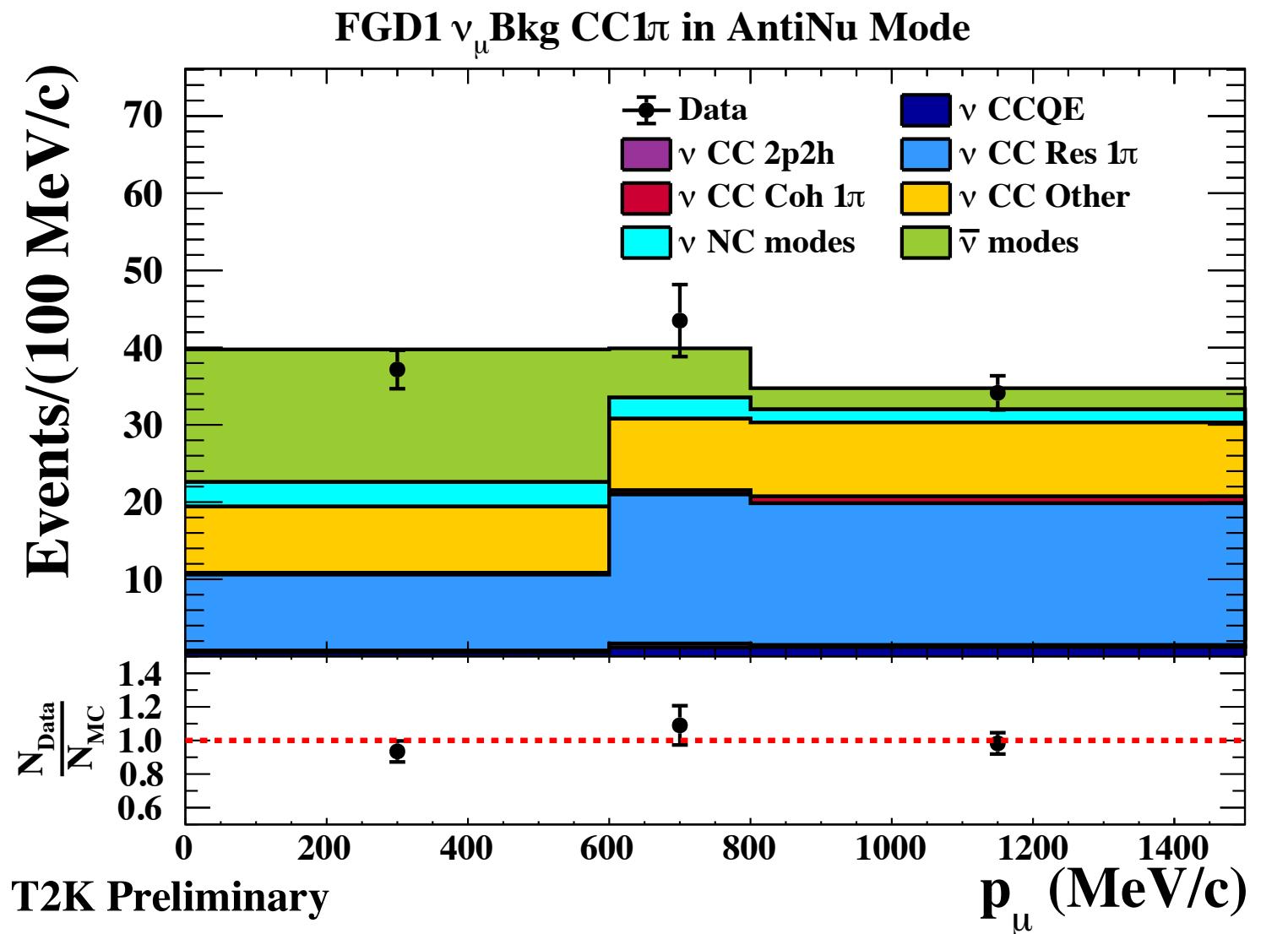
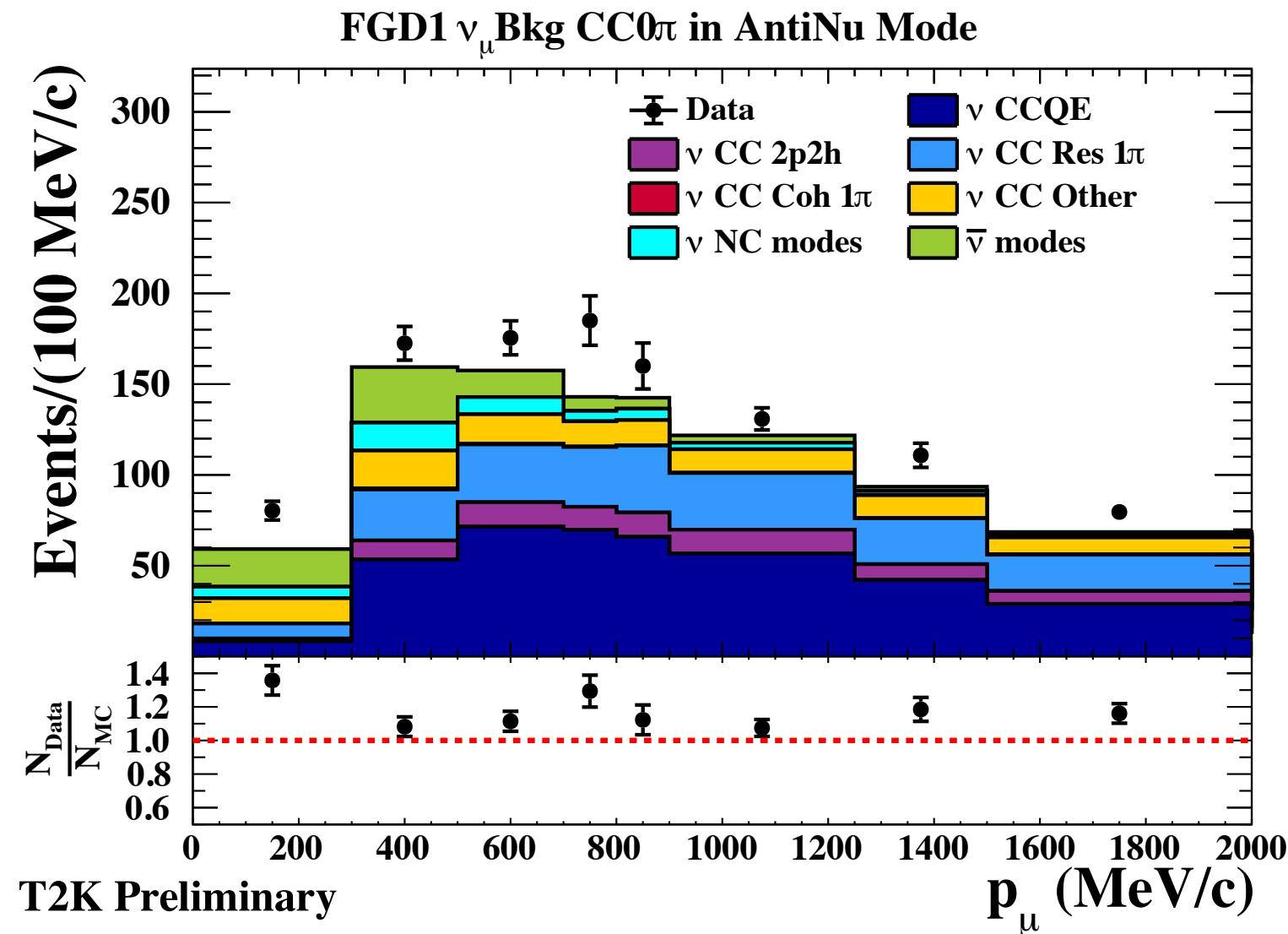


*After ND fit*

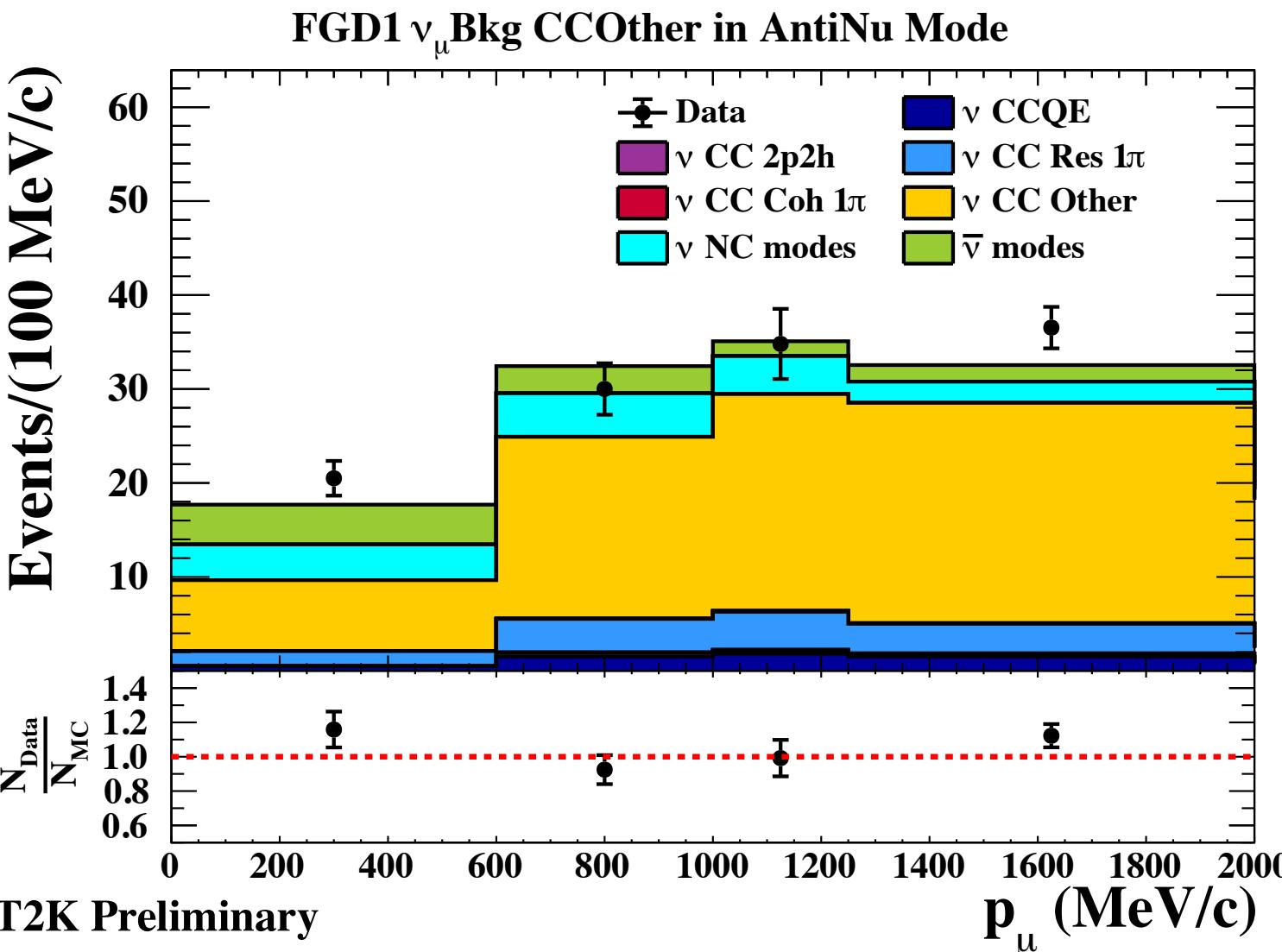
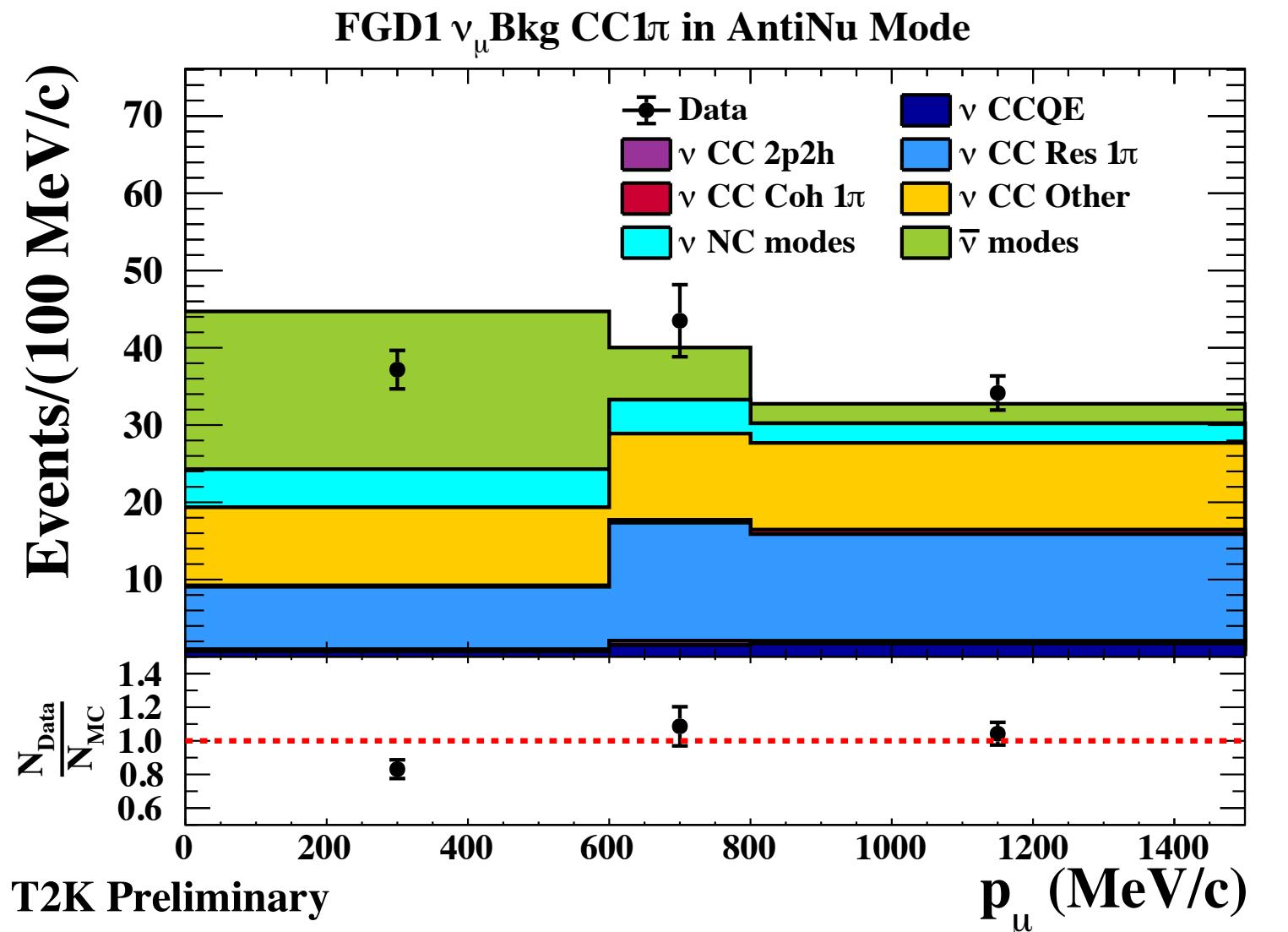
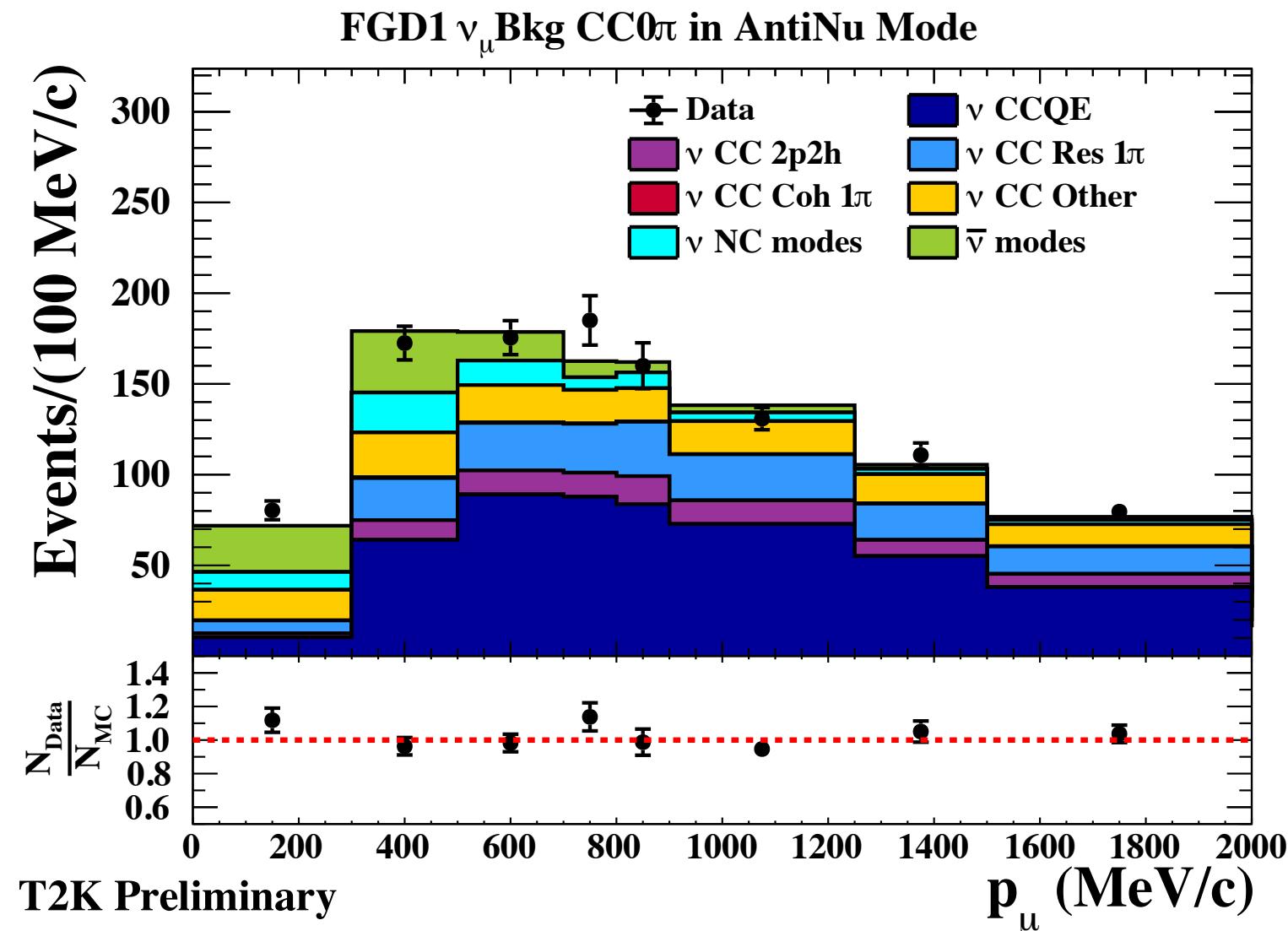


# Fits to near detector FGD1 $\nu_\mu$ wrong-sign bkg. in $\bar{\nu}$ -mode

Before ND fit

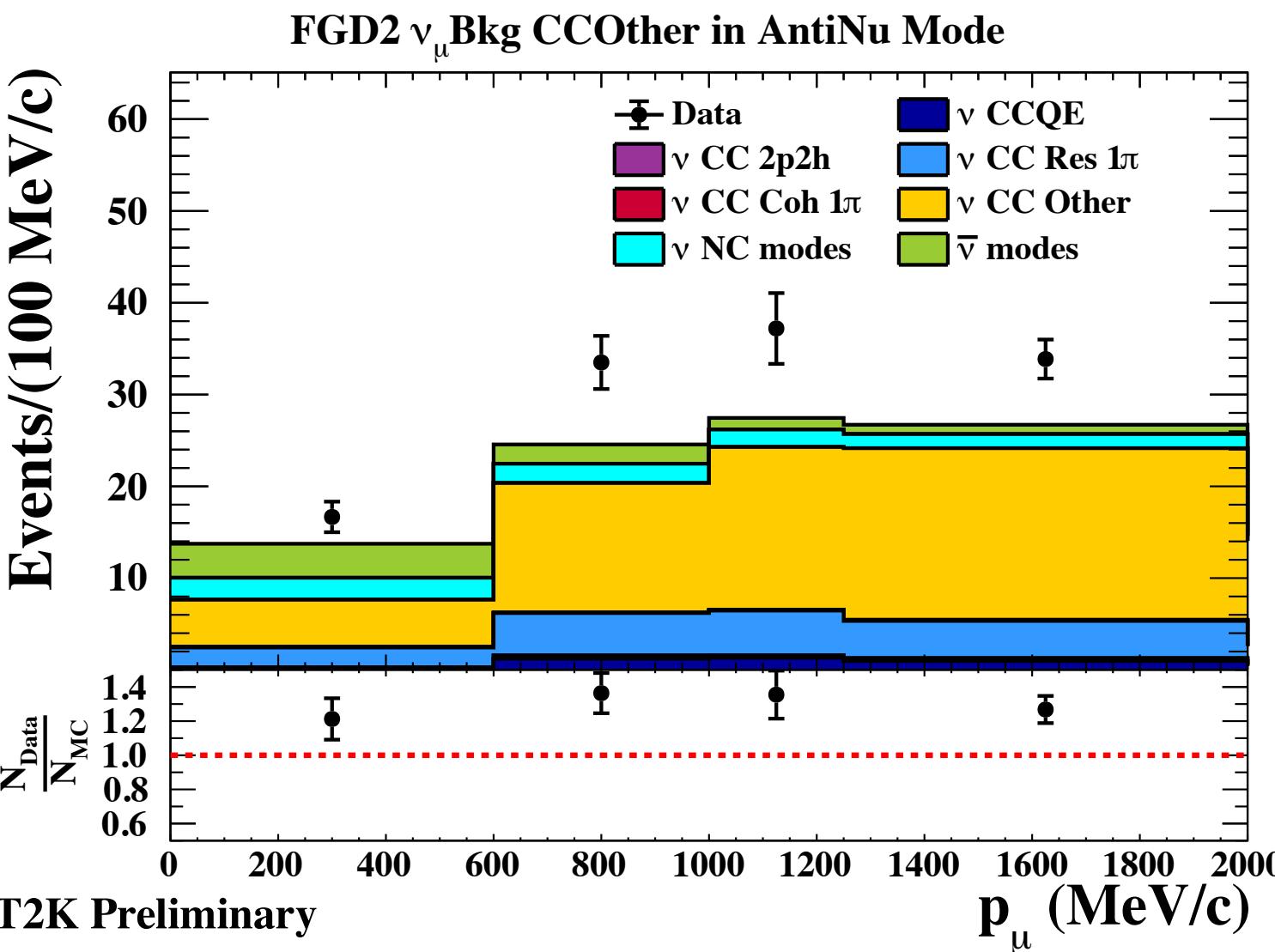
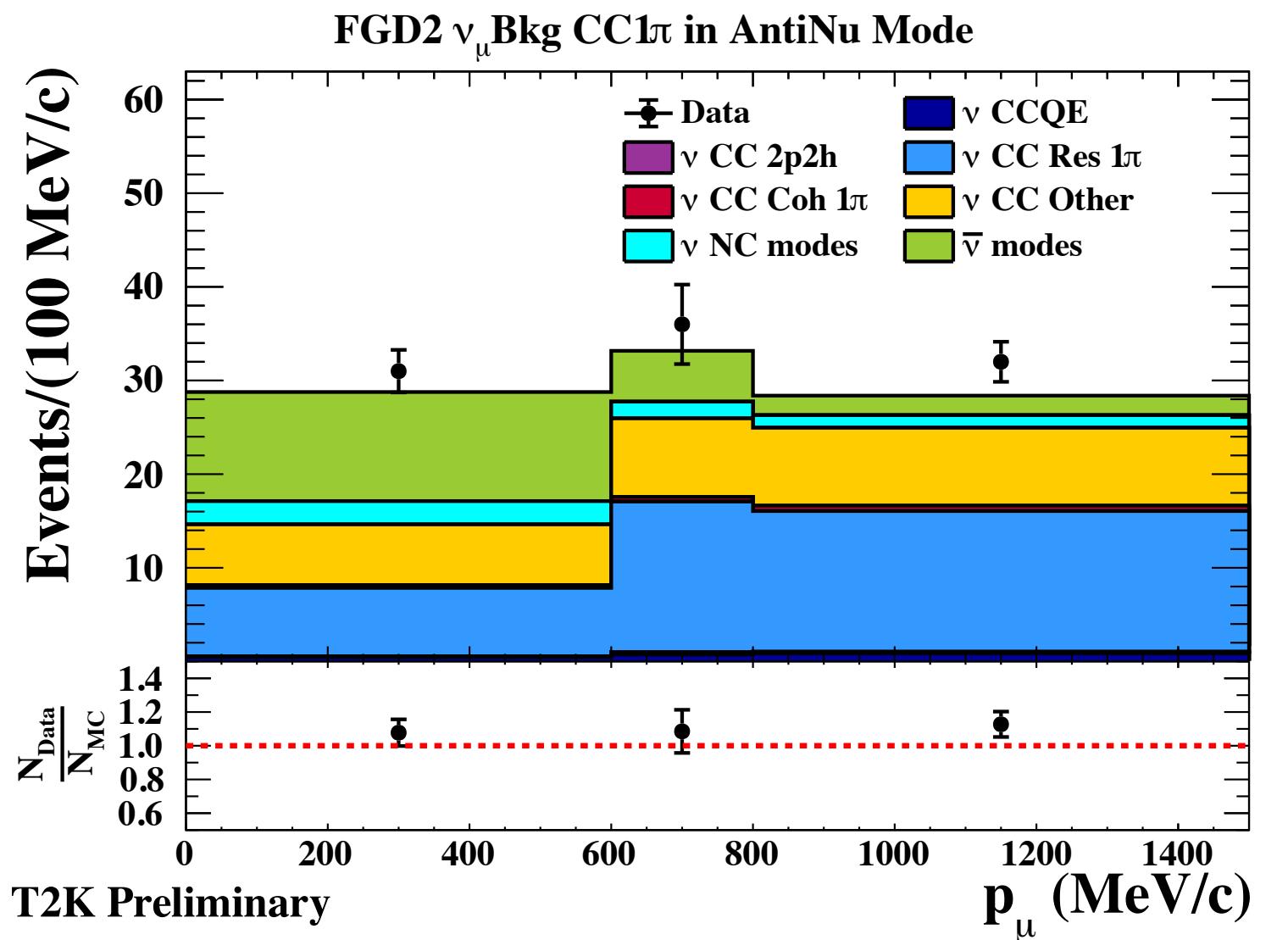
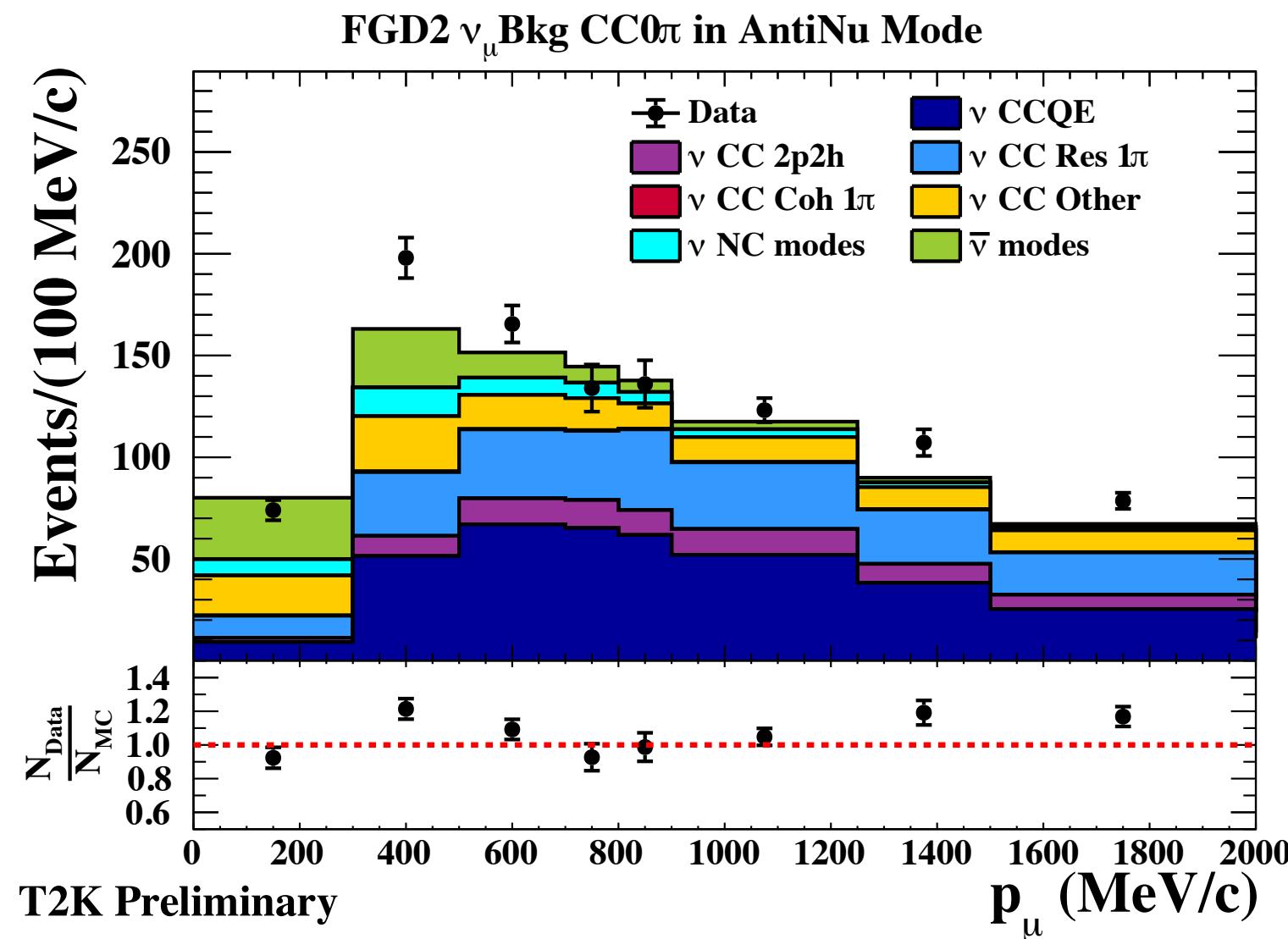


After ND fit

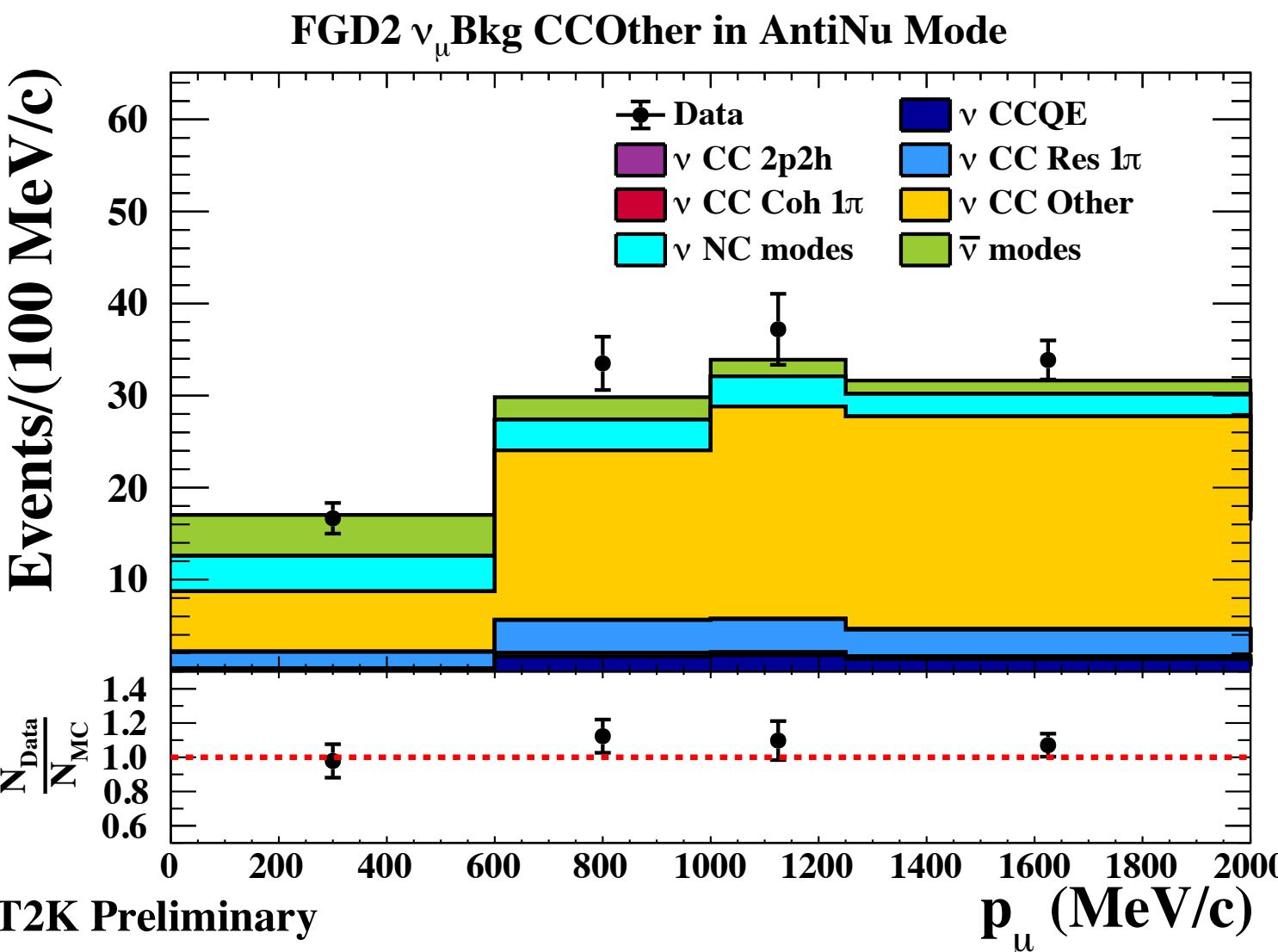
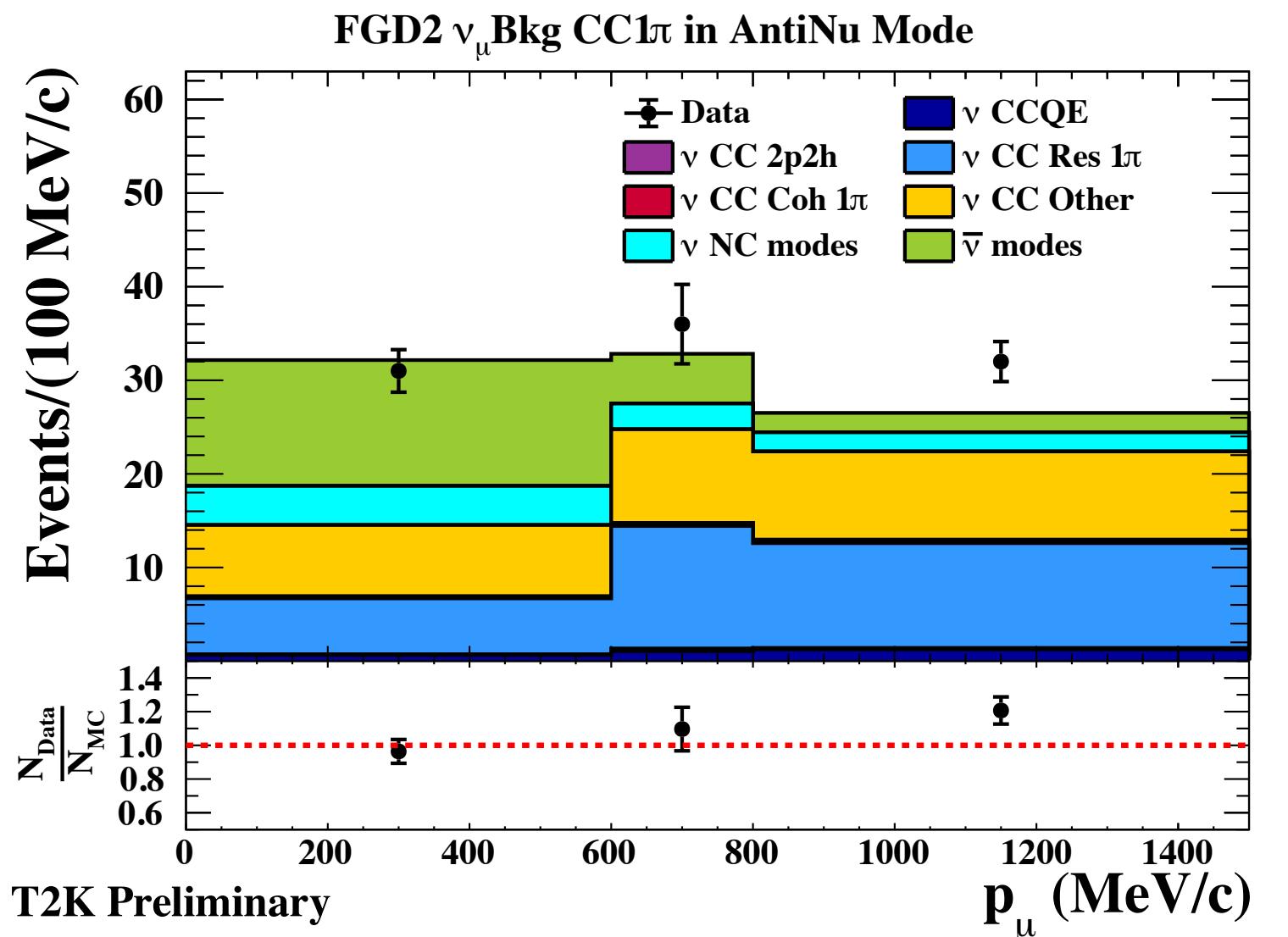
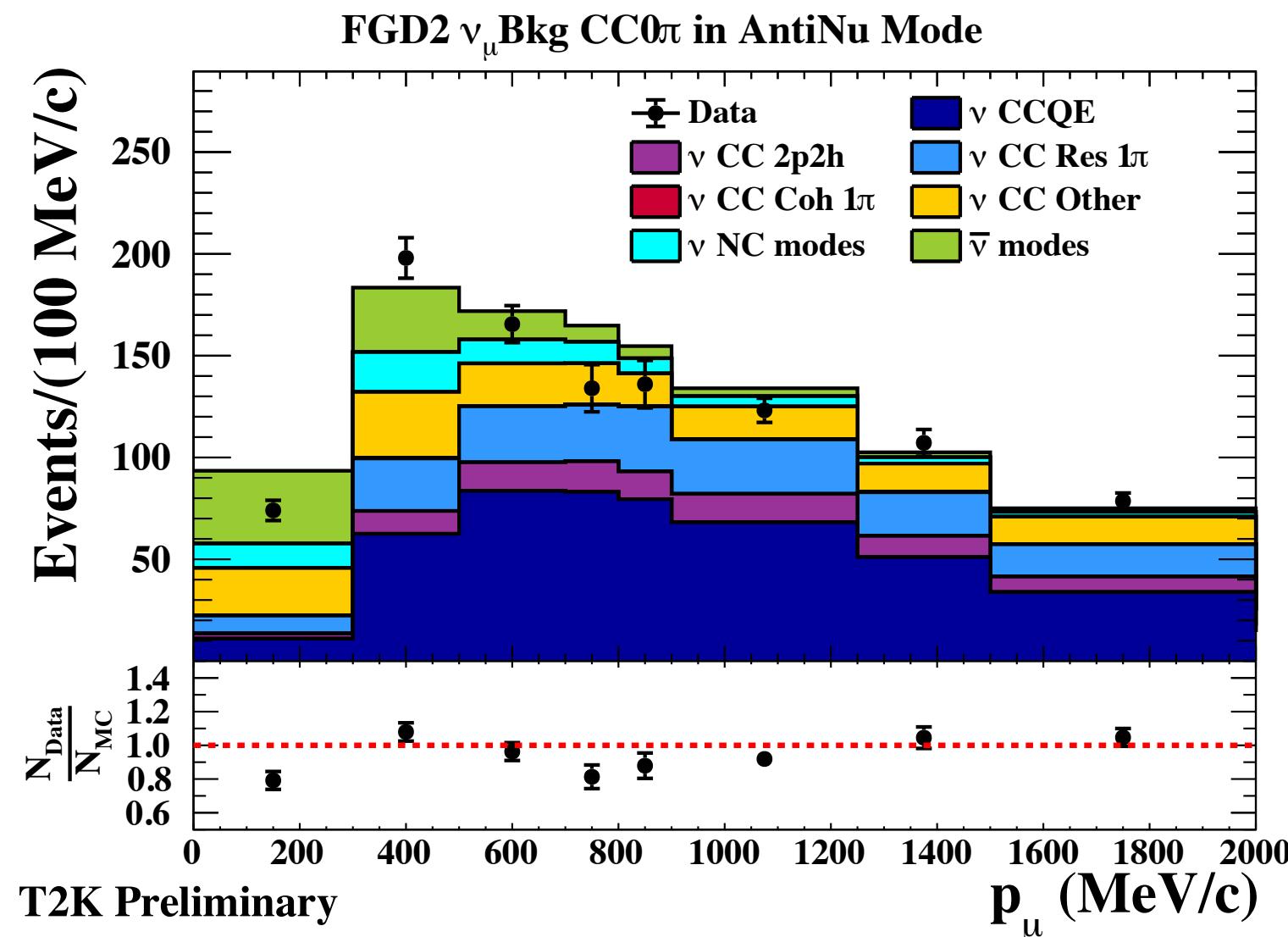


# Fits to near detector FGD2 $\nu_\mu$ wrong-sign bkg. in $\bar{\nu}$ -mode

*Before ND fit*

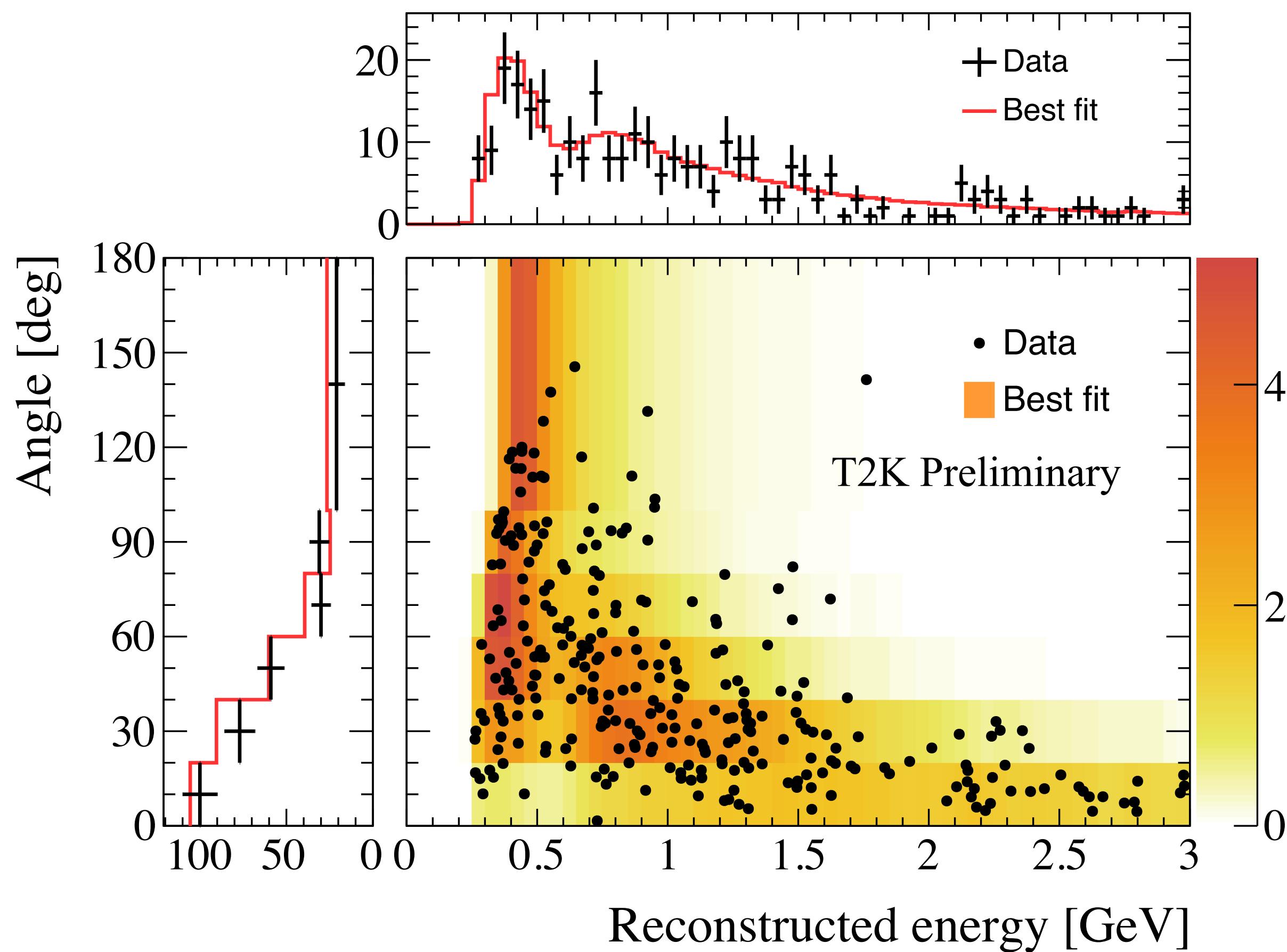


*After ND fit*

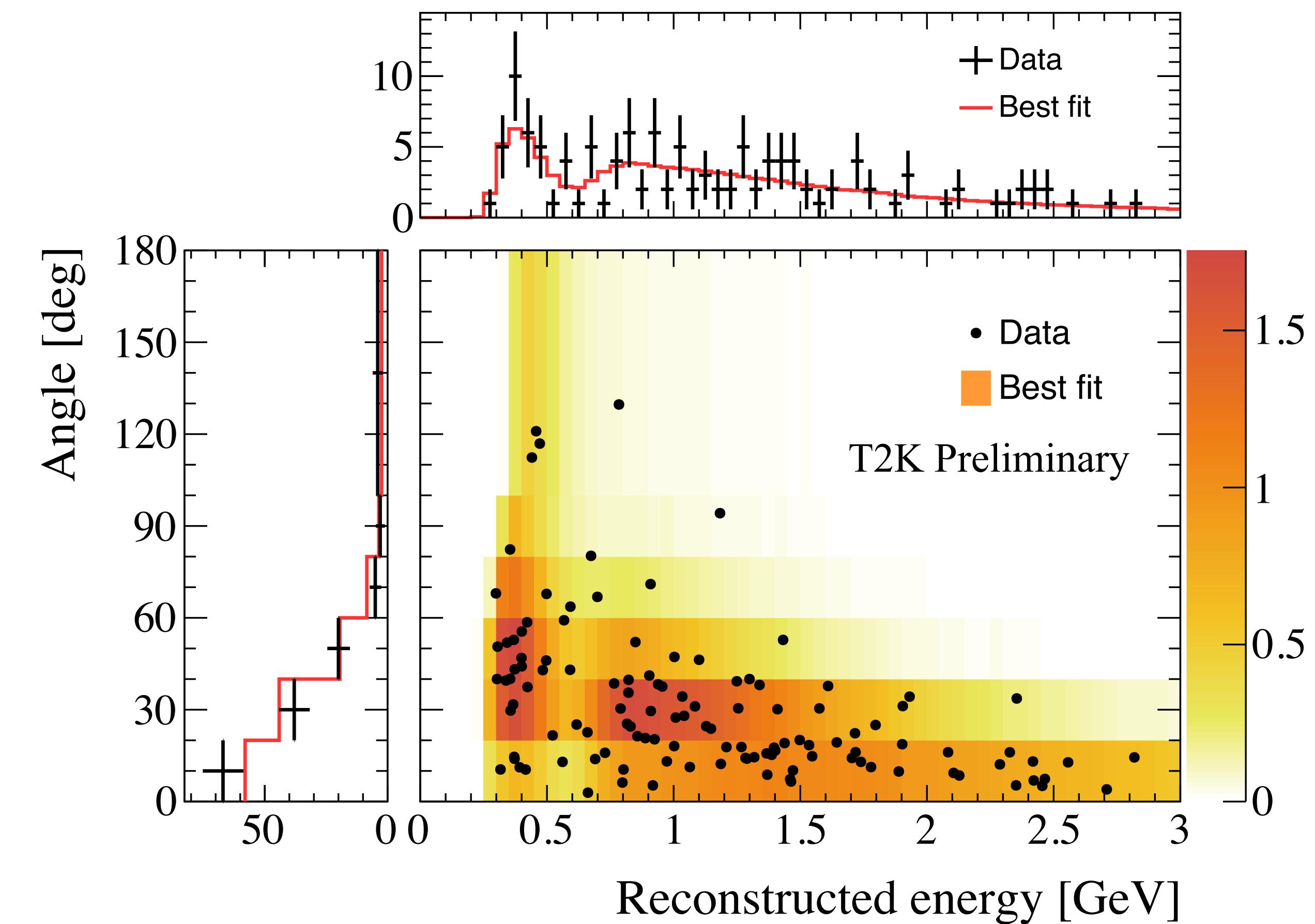


# Joint fit to near and far detector data: muon-like samples at Super-K

Neutrino mode  
1 muon-like ring

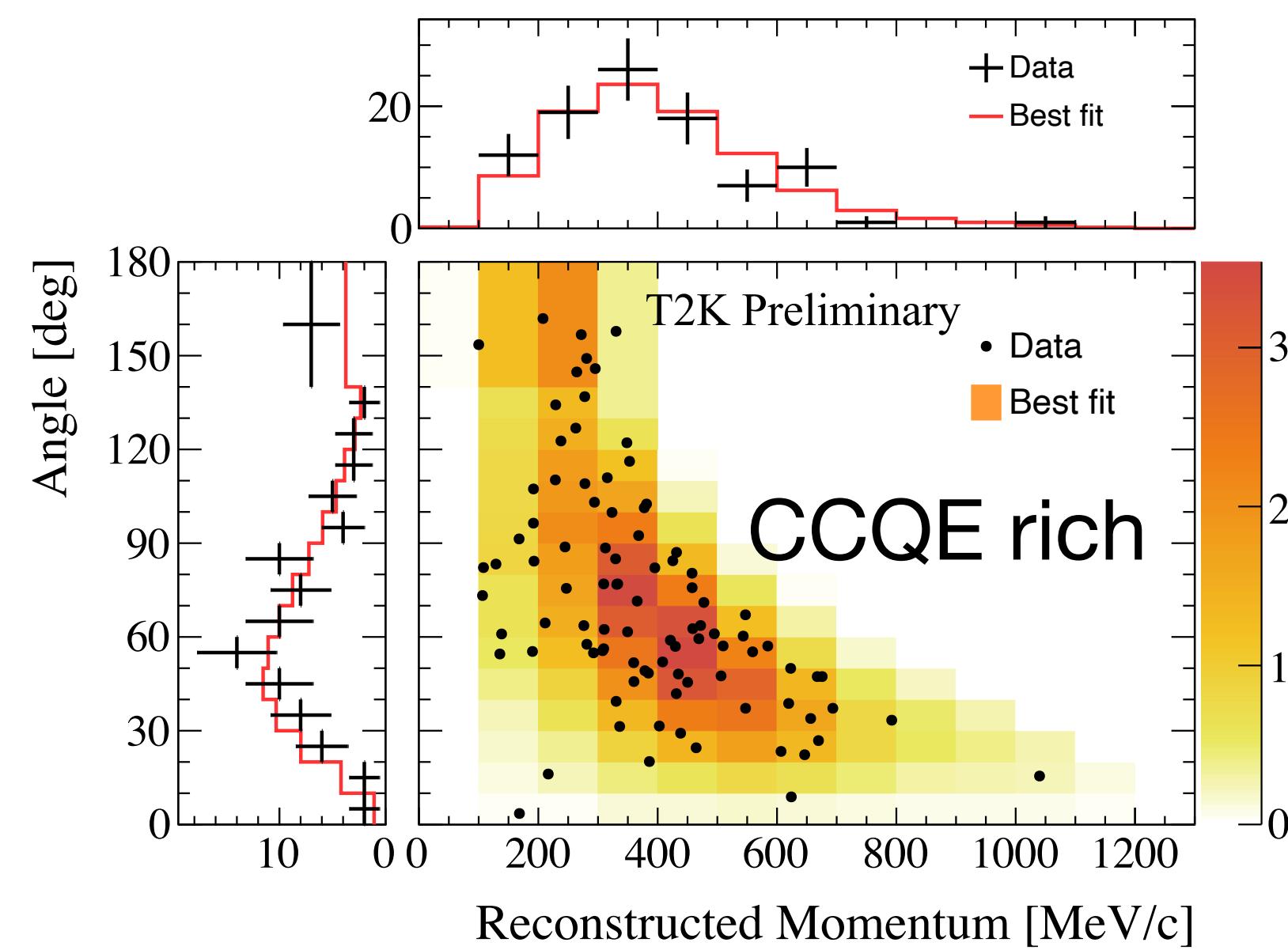


Anti-neutrino mode  
1 muon-like ring

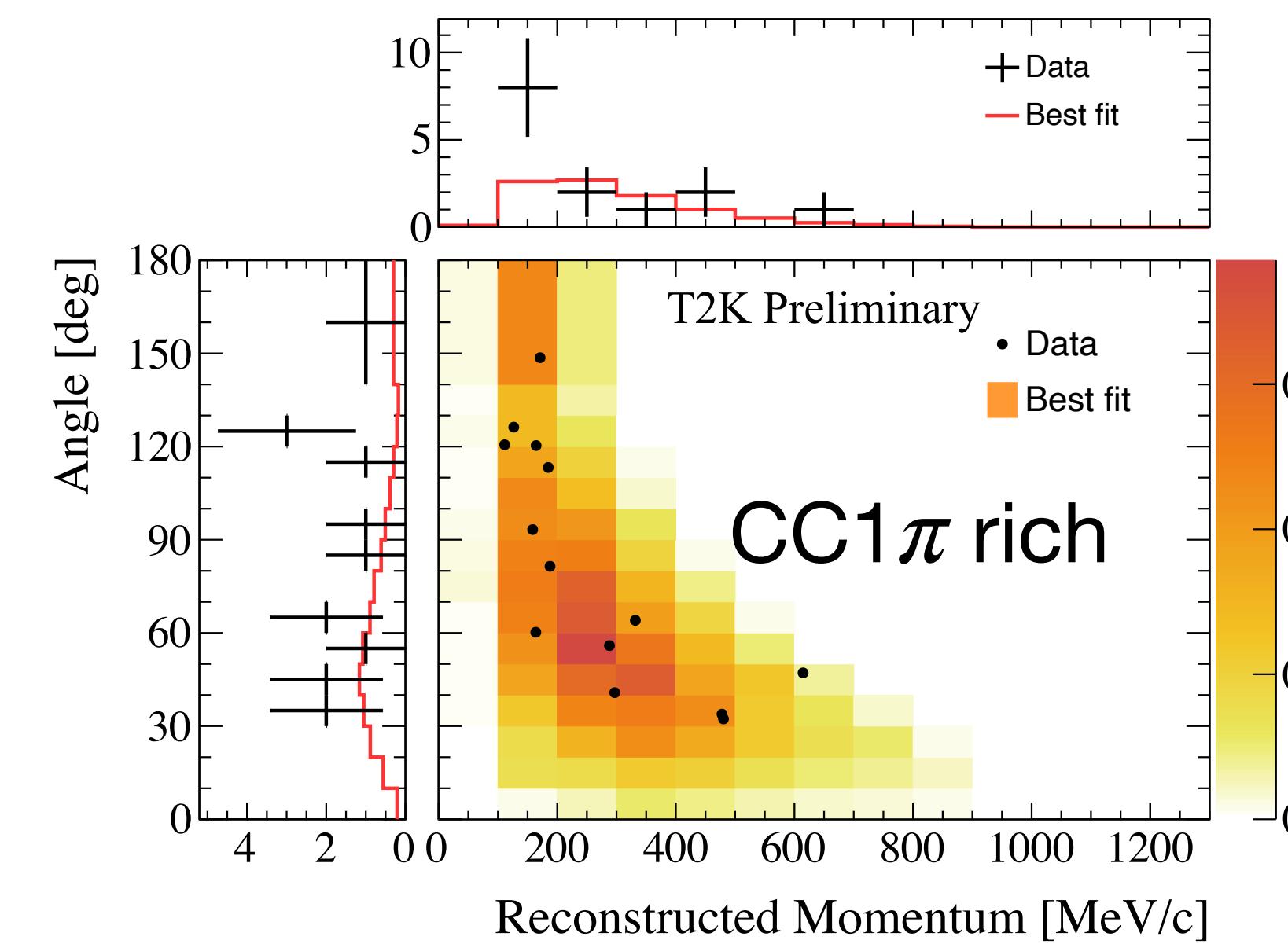


# Joint fit to near and far detector data: electron-like samples at Super-K

Neutrino mode  
1 electron-like ring



Neutrino mode  
1 electron-like ring + 1 decay electron



Anti-neutrino mode  
1 electron-like ring

